

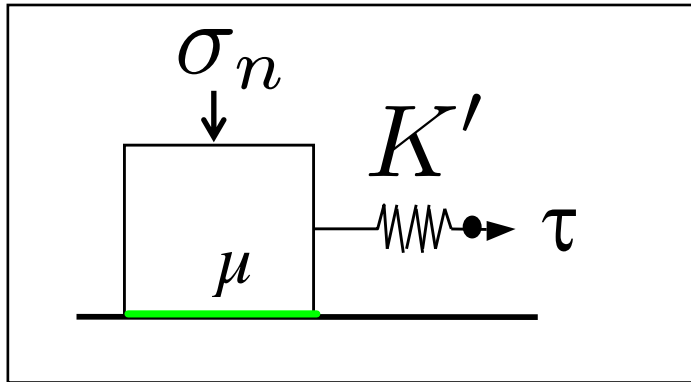
# Mechanics of Earthquakes and Faulting

Lecture 10, 2 Mar. 2021

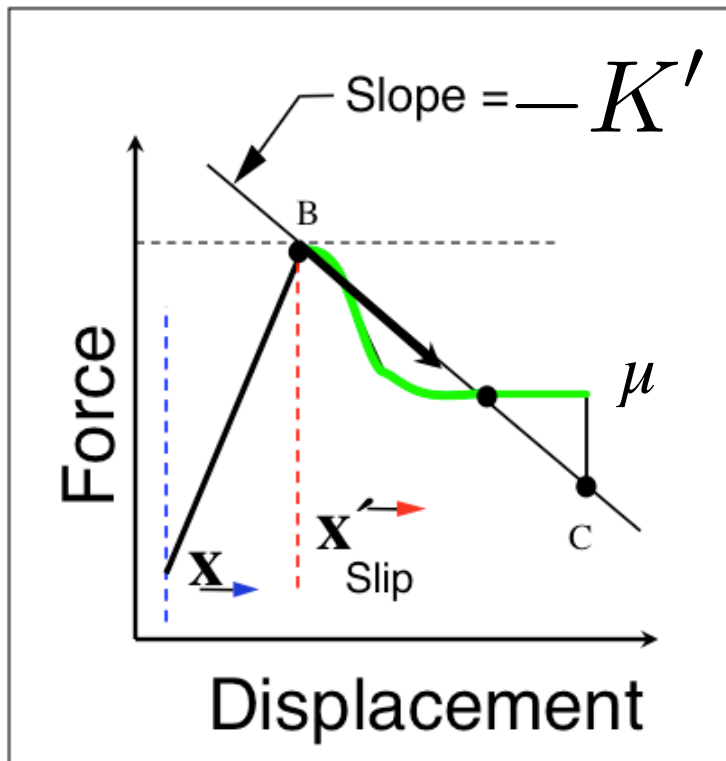
[www.geosc.psu.edu/Courses/Geosc508](http://www.geosc.psu.edu/Courses/Geosc508)

- Friction Constitutive Laws for Faulting
- Critical stiffness and stability transition from stable to unstable faulting
- SHS test to measure RSF parameters
- Earthquake Nucleation, Critical Fault Patch Size
- Stress Distribution for Propagating Rupture --(Crack Tip) Cohesive Zone
- Laboratory Observations of The transition from stable to unstable frictional sliding: Confirmation of the concept of a critical fault weakening rate with slip ( $k_c$ )

## Mechanics of Frictional Sliding: Stick-slip



$$\frac{\tau(\theta, v)}{\sigma_n} = \mu_o + a \ln \left( \frac{v}{v_o} \right) + b \ln \left( \frac{v_o \theta}{D_c} \right)$$

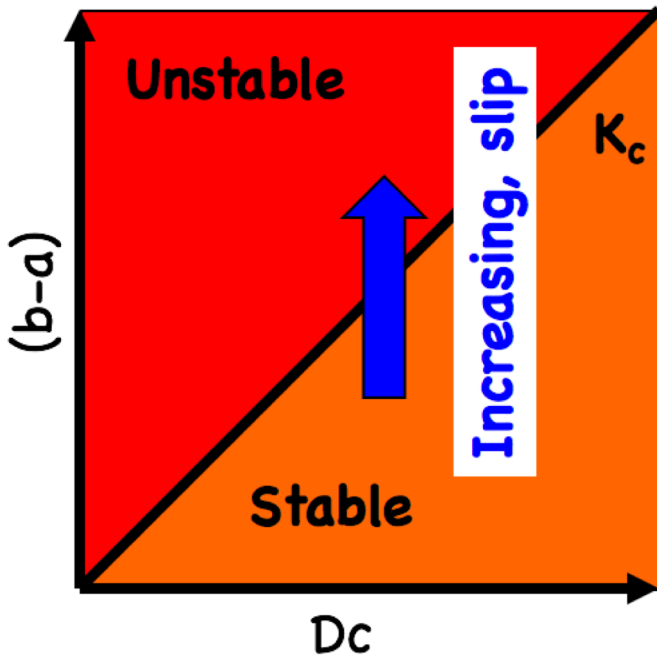


$$K_c = \frac{\sigma_n(b - a)}{D_c} \left[ 1 + \frac{mV_o^2}{\sigma_n a D_c} \right]$$

*Rice & Ruina, 1983; Gu et al., 1984; Roy and Marone, 1996*

$$K_c \approx \frac{\sigma_n(b - a)}{D_c}$$

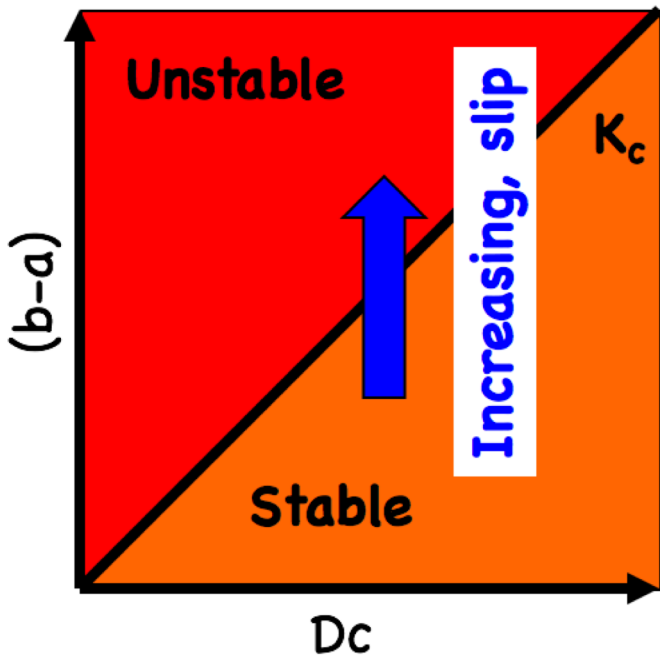
**Unstable if  $K < K_c$**



## Stability of sliding

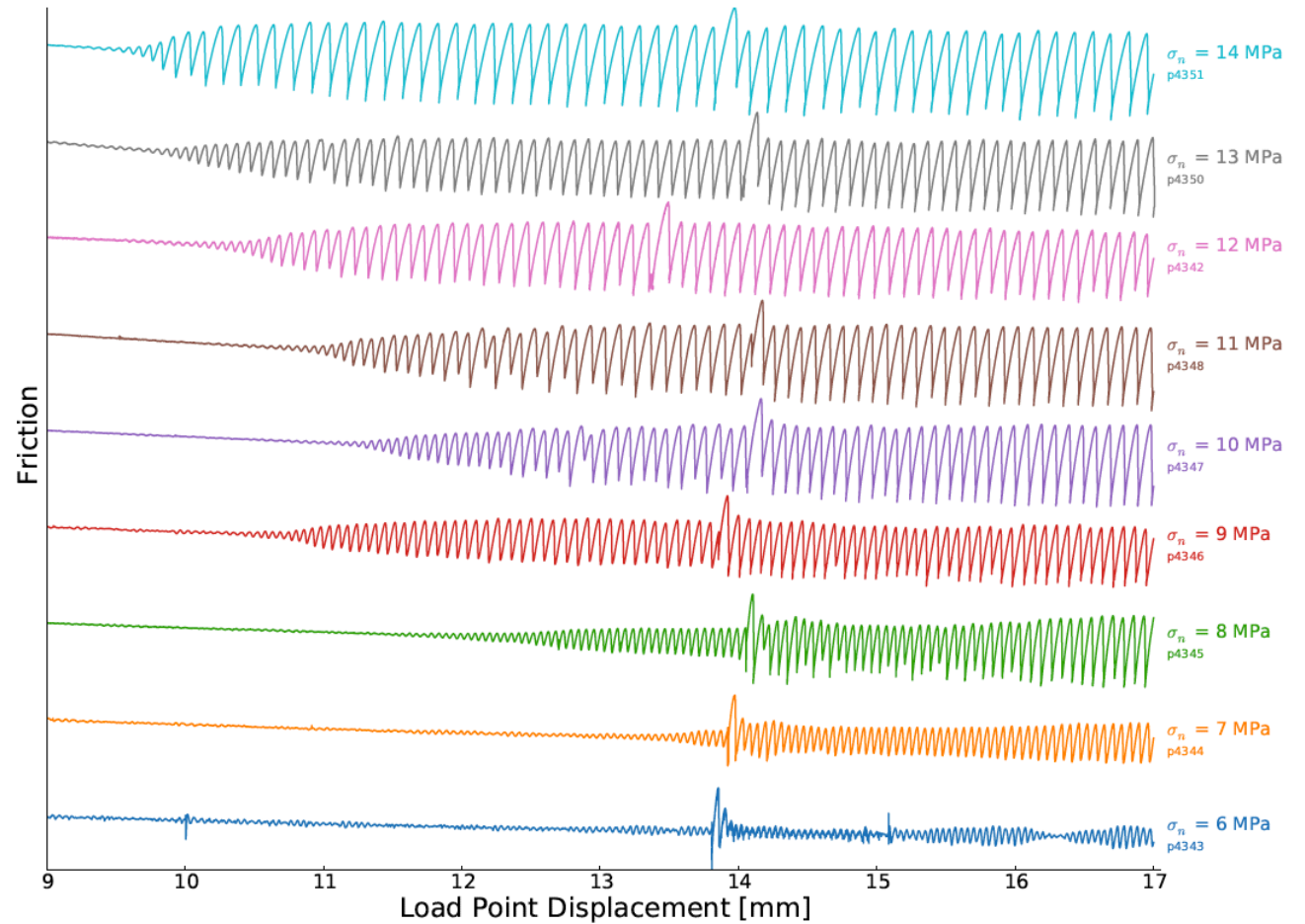
$$K_c = \frac{\sigma_n(b-a)}{D_c} \left[ 1 + \frac{mV_o^2}{\sigma_n a D_c} \right]$$

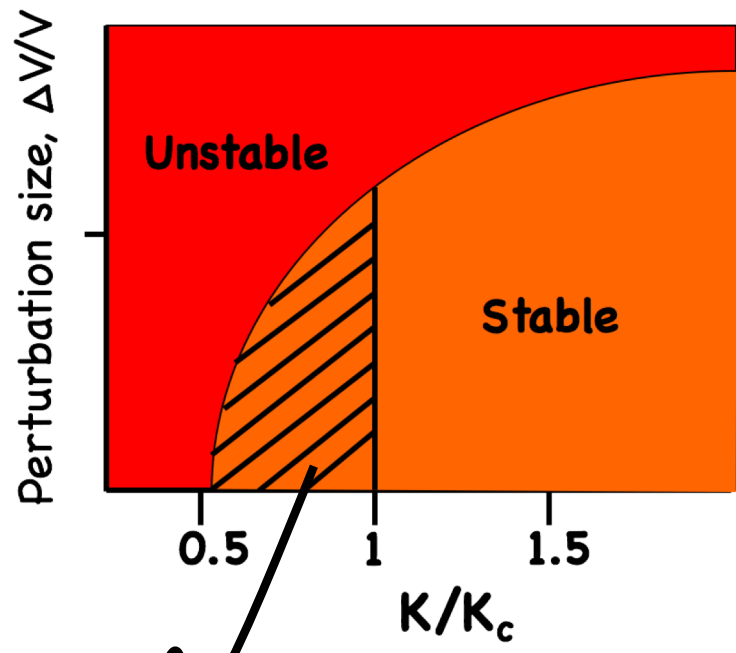
Stability line:  $K = K_c$



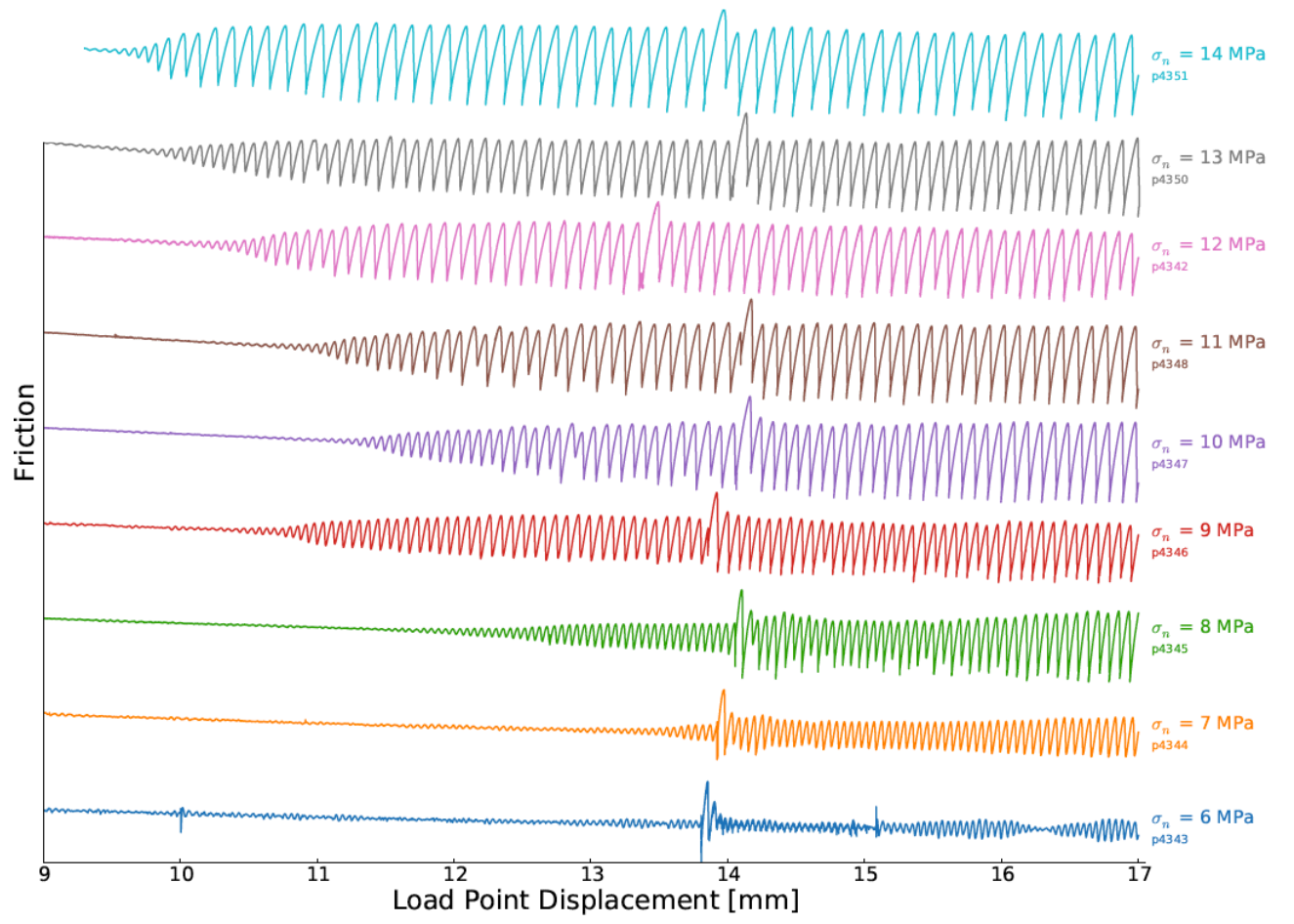
## Stability of sliding

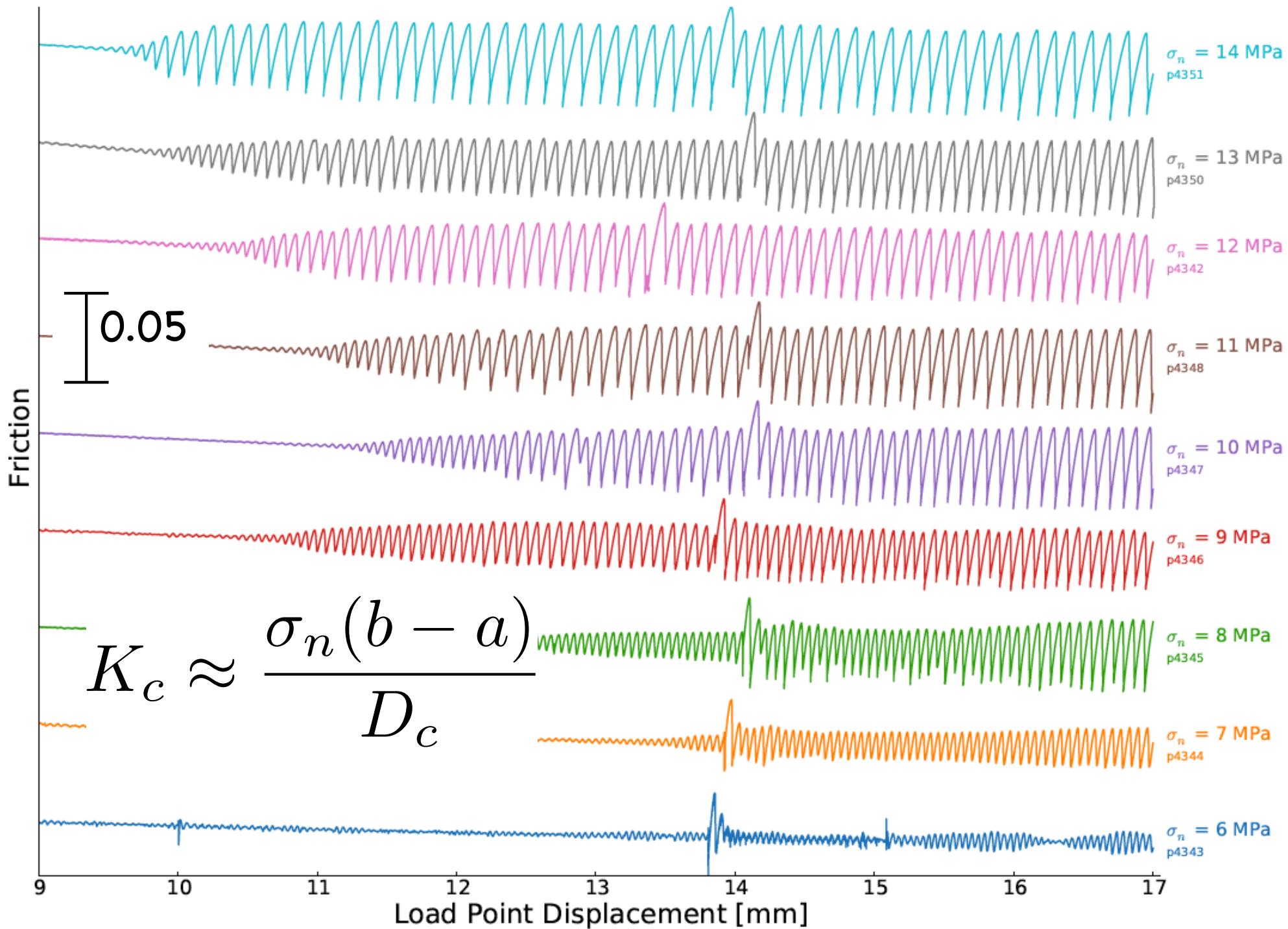
$$K_c = \frac{\sigma_n (b - a)}{D_c} \left[ 1 + \frac{m V_o^2}{\sigma_n a D_c} \right]$$

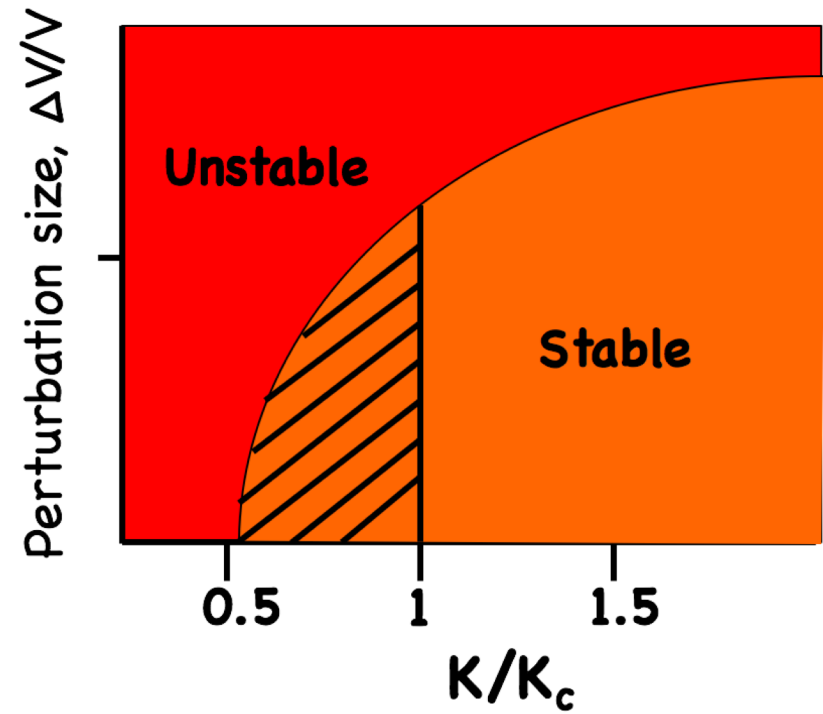
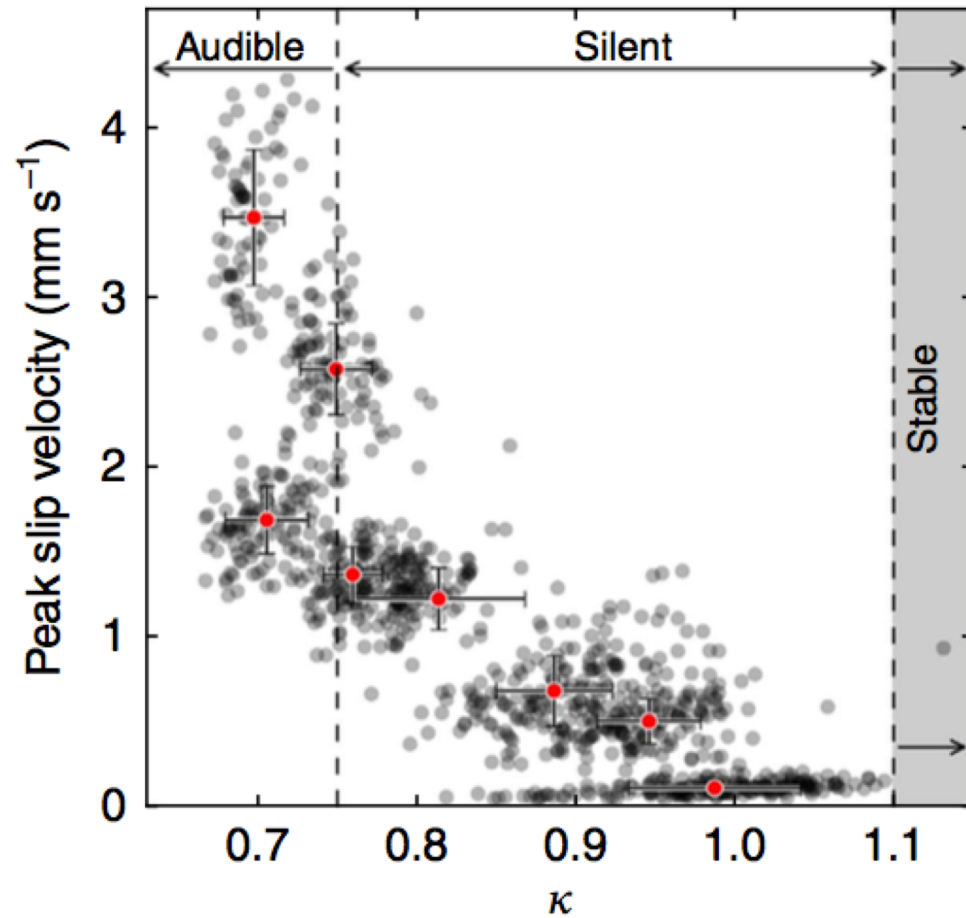




Quasidynamic

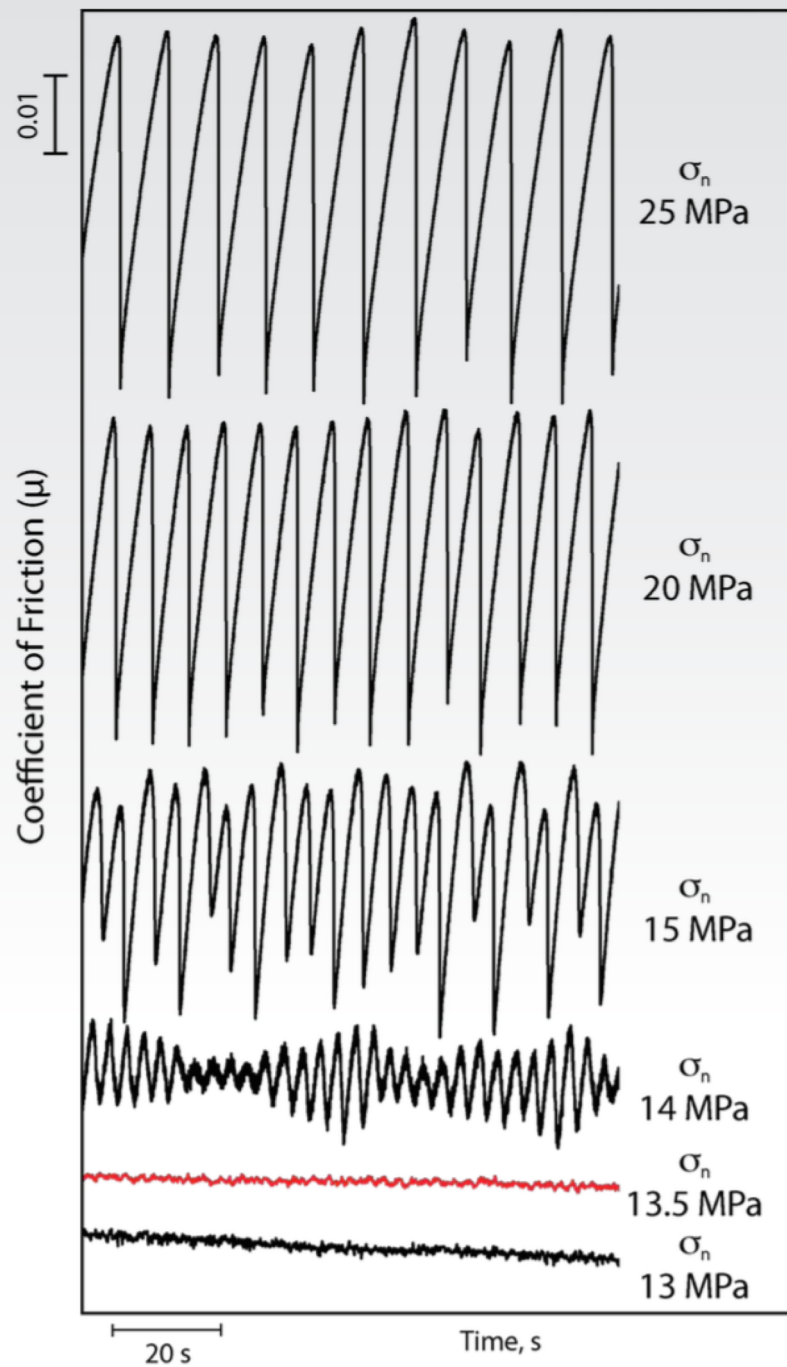






Unstable slip if

$$K < K_c \approx \frac{\sigma_n(b - a)}{D_c}$$



Complex Behavior  
(Period Doubling) Near  
The Stability Boundary



$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

reference value of base friction      reference velocity      critical slip distance

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

Dieterich, aging law

$$\frac{d\theta}{dt} = -\frac{V\theta}{D_c} \ln\left(\frac{V\theta}{D_c}\right)$$

Ruina, slip law

$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o\theta}{D_c}\right)$$

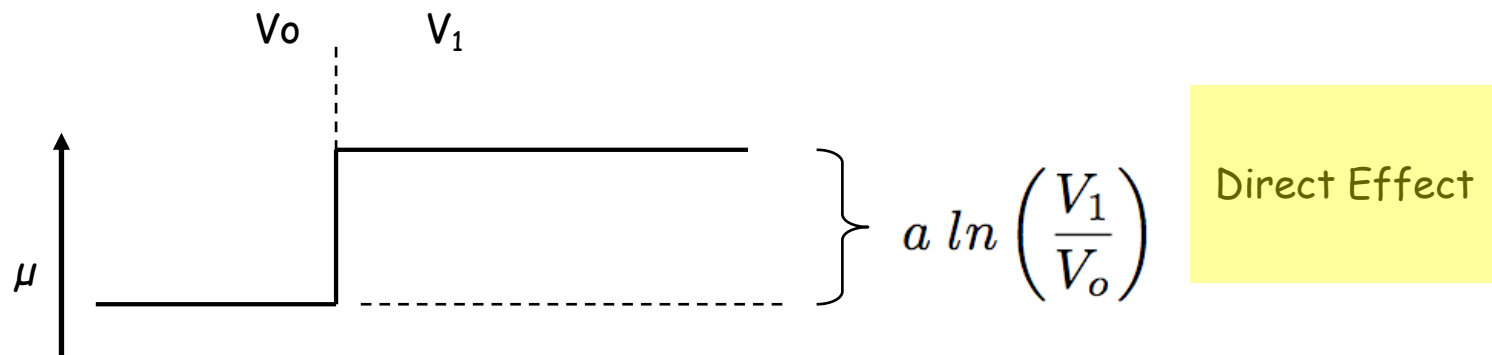
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Ruina, slip law

Velocity Step test



Friction vs.  
shear displacement

$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

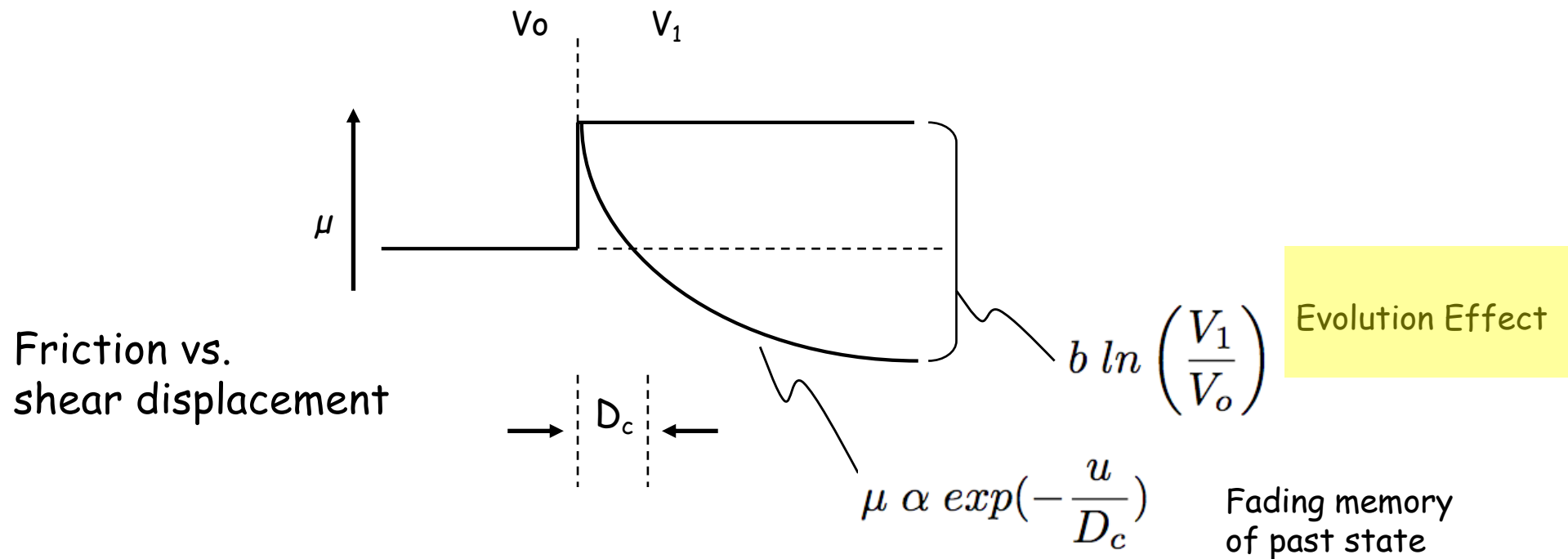
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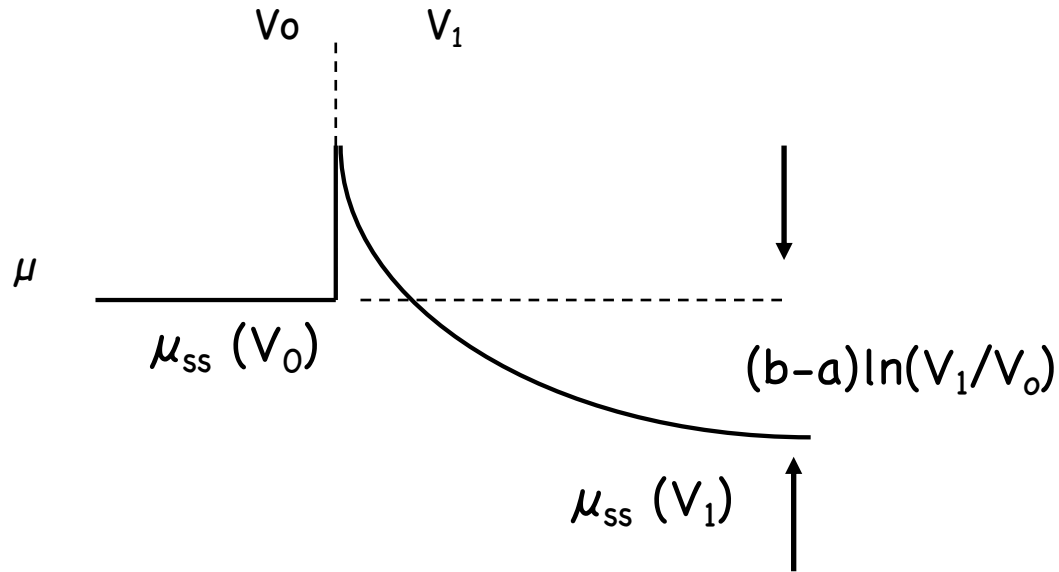
### Velocity Step test



$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o\theta}{D_c}\right)$$

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

Dieterich law



$$\frac{d\theta}{dt} = -\frac{V\theta}{D_c} \ln\left(\frac{V\theta}{D_c}\right)$$

Ruina law

Steady-state sliding:

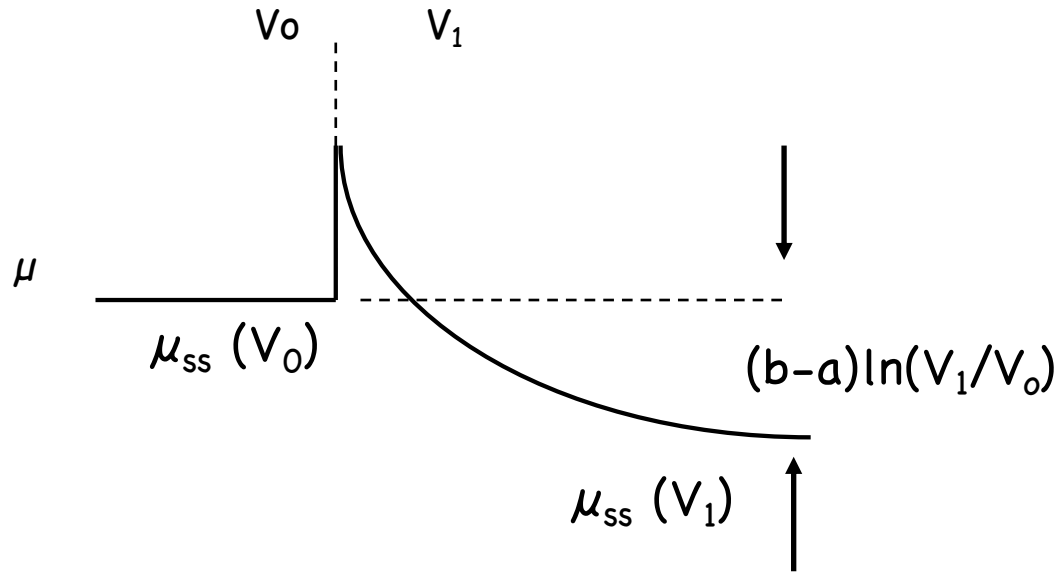
$$\frac{d\theta}{dt} = 0$$

Thus, for both laws  $V\theta/D_c = 1$ , which means:  $\theta_{ss} = \frac{D_c}{V_1}$

$$(1) \quad \mu(\theta, V) = \mu_o + a \ln \left( \frac{V}{V_o} \right) + b \ln \left( \frac{V_o \theta}{D_c} \right)$$

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

Dieterich law



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Ruina law

Steady-state sliding:

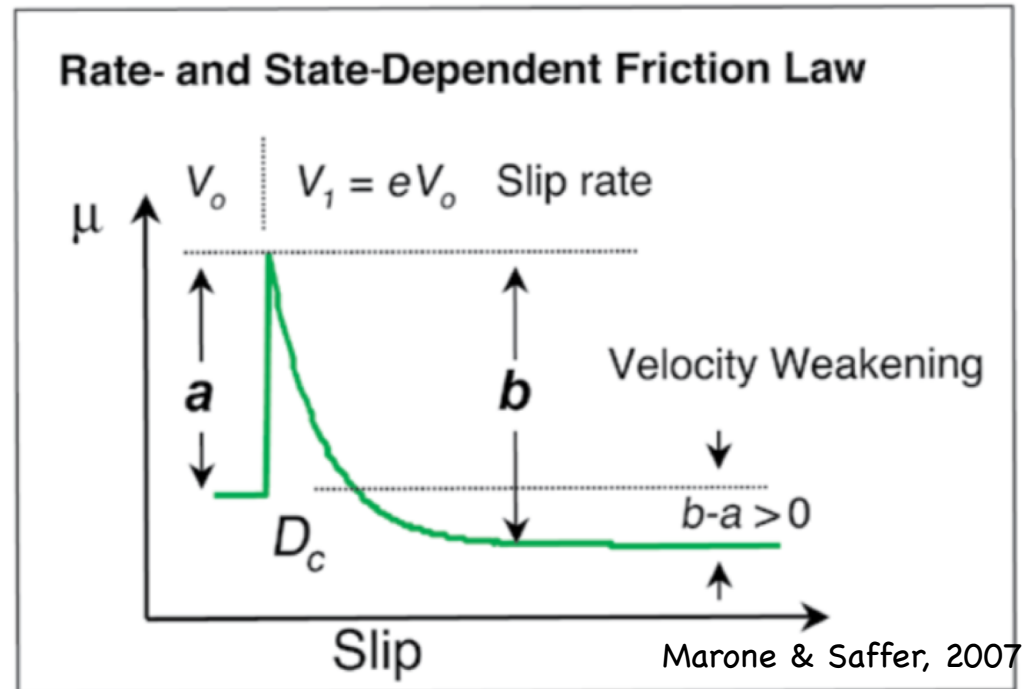
$$\theta_{ss} = \frac{D_c}{V_1}$$

then (1) becomes:  $\mu_1 - \mu_o = (a - b) \ln \left( \frac{V_1}{V_o} \right)$

or  $(a - b) = \frac{\Delta\mu}{\Delta \ln V}$

$$(1) \quad \mu(\theta, V) = \mu_o + a \ln \left( \frac{V}{V_o} \right) + b \ln \left( \frac{V_o \theta}{D_c} \right)$$

(b)

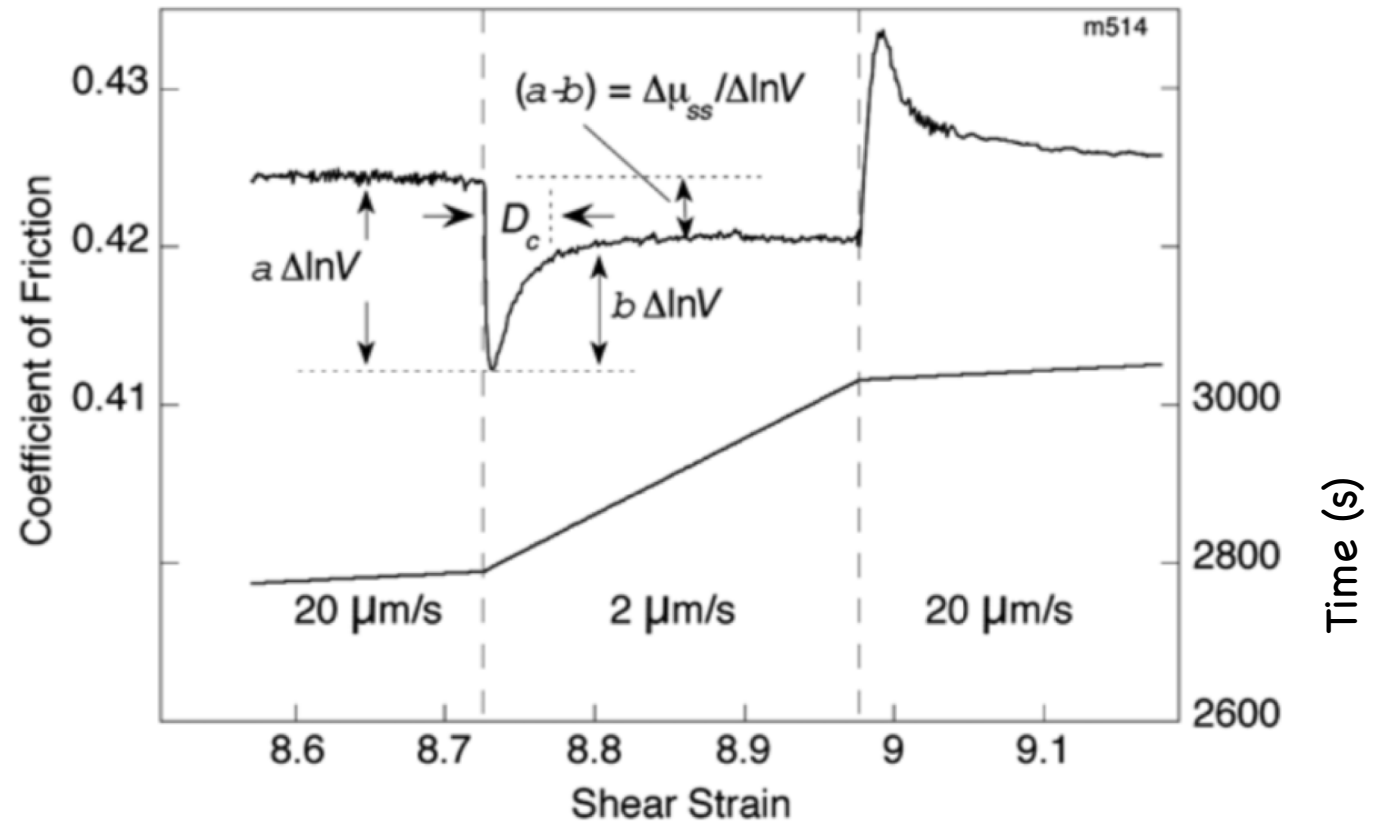


Steady-state sliding:

$$\theta_{ss} = \frac{D_c}{V_1}$$

then (1) becomes: 
$$\mu_1 - \mu_o = (a - b) \ln \left( \frac{V_1}{V_o} \right)$$

or 
$$(a - b) = \frac{\Delta \mu}{\Delta \ln V}$$



Marone & Saffer, 2007

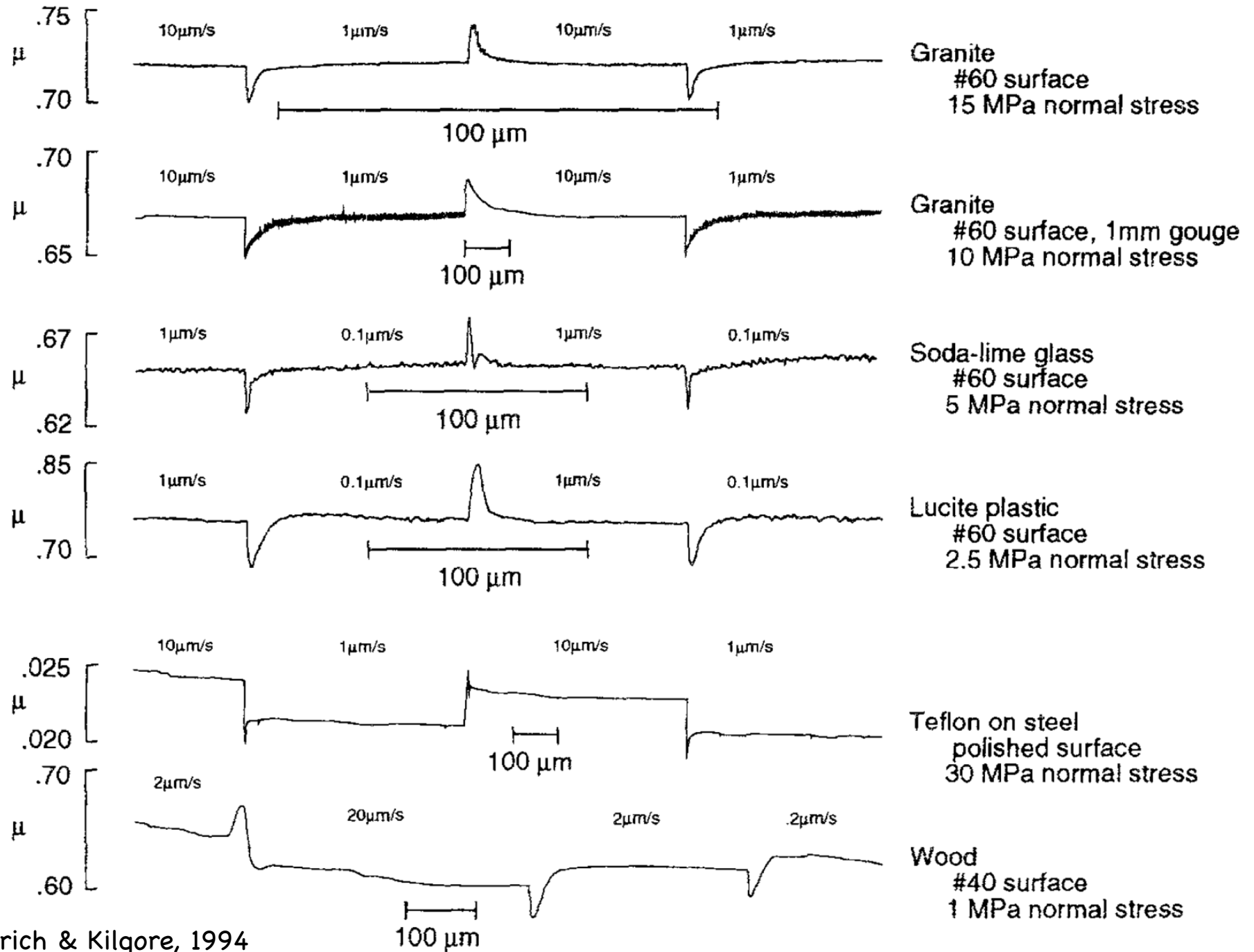
Steady-state sliding:

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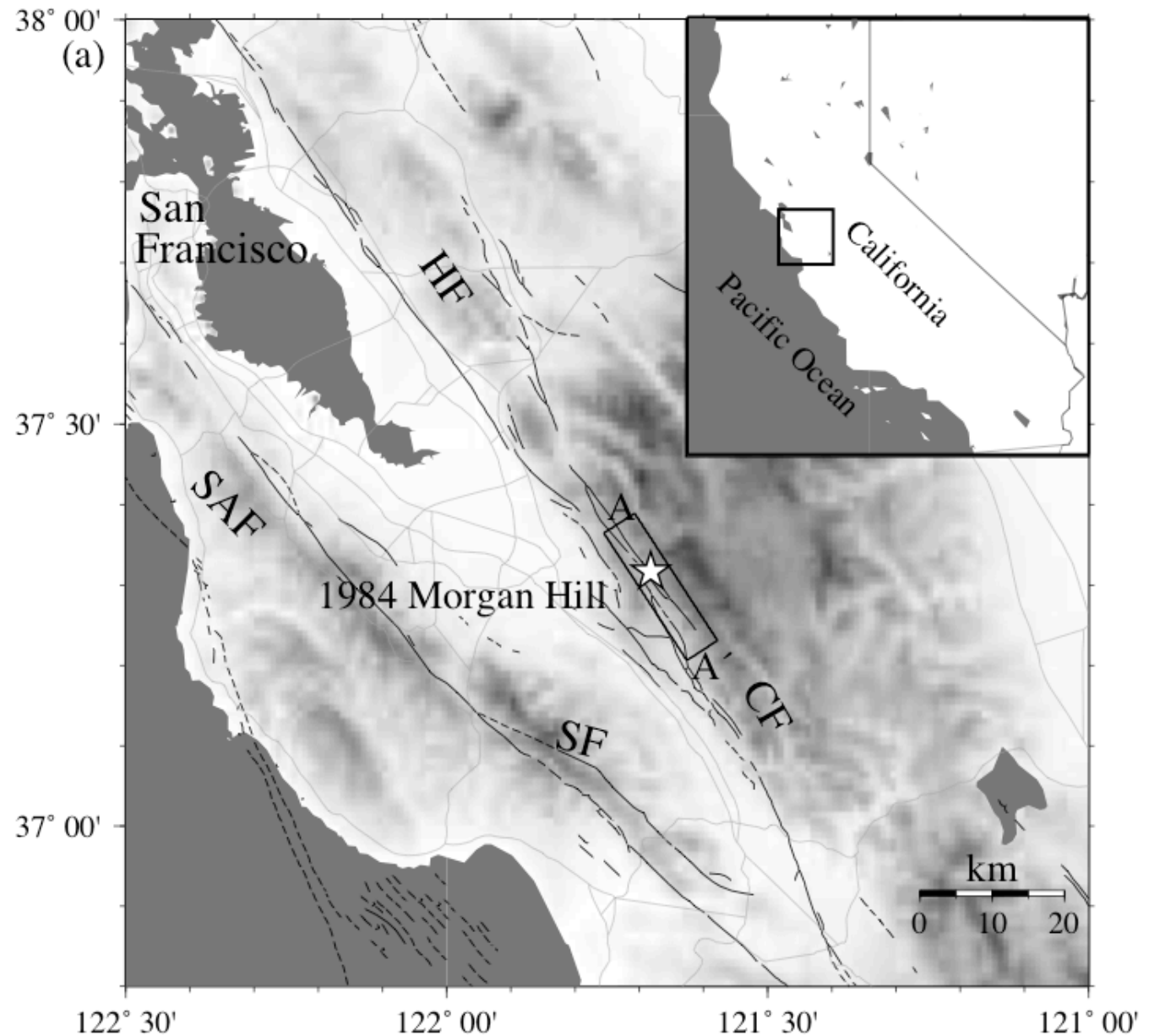
# Universality of rate and state friction



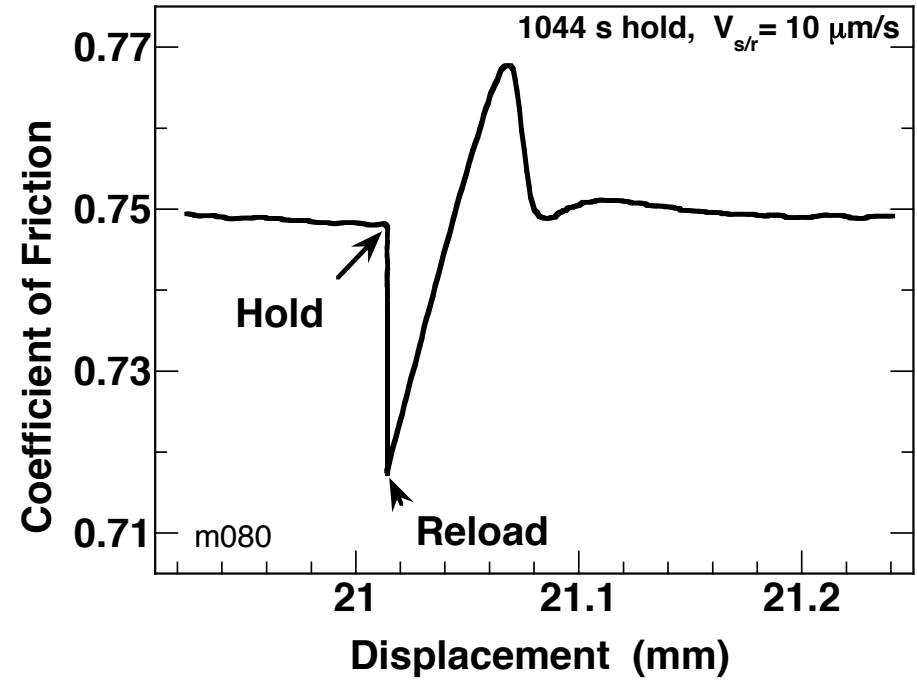
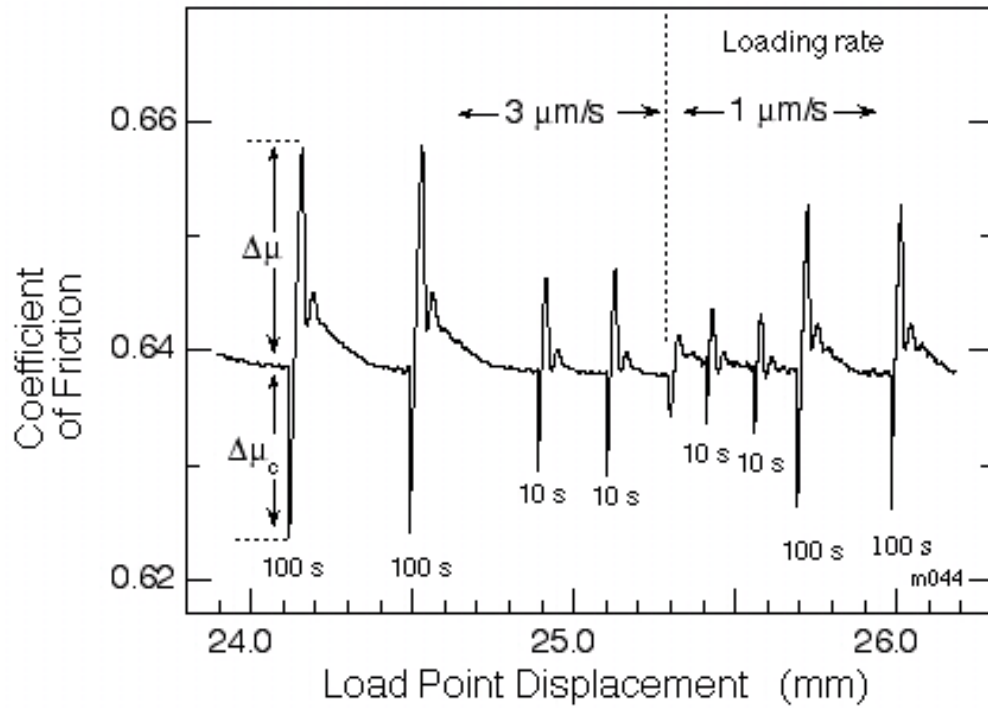


# Fault Healing and the Seismic Cycle: Repeating Earthquakes

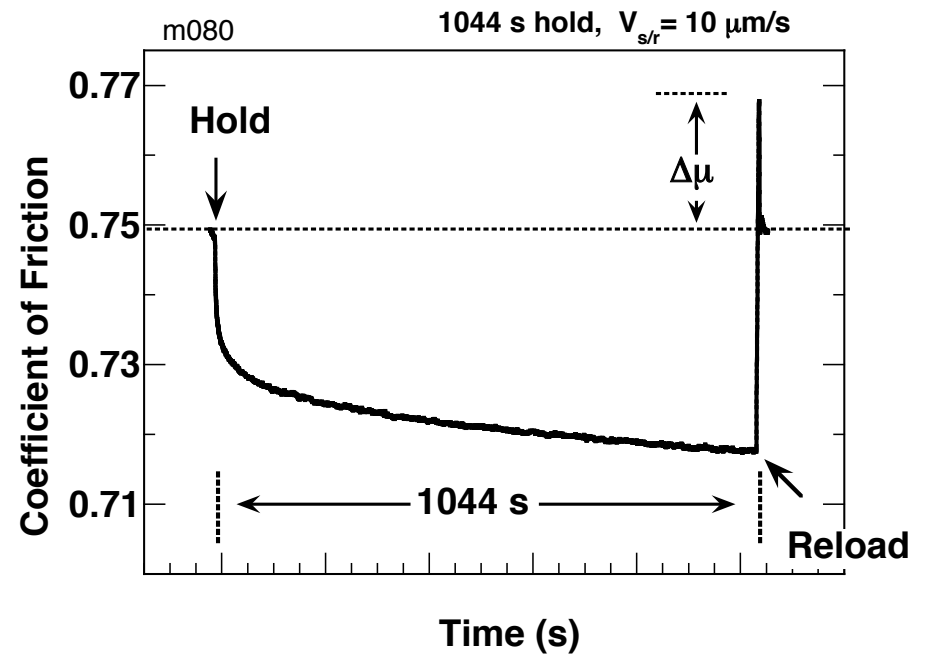
How do faults  
regain strength  
between  
earthquakes?

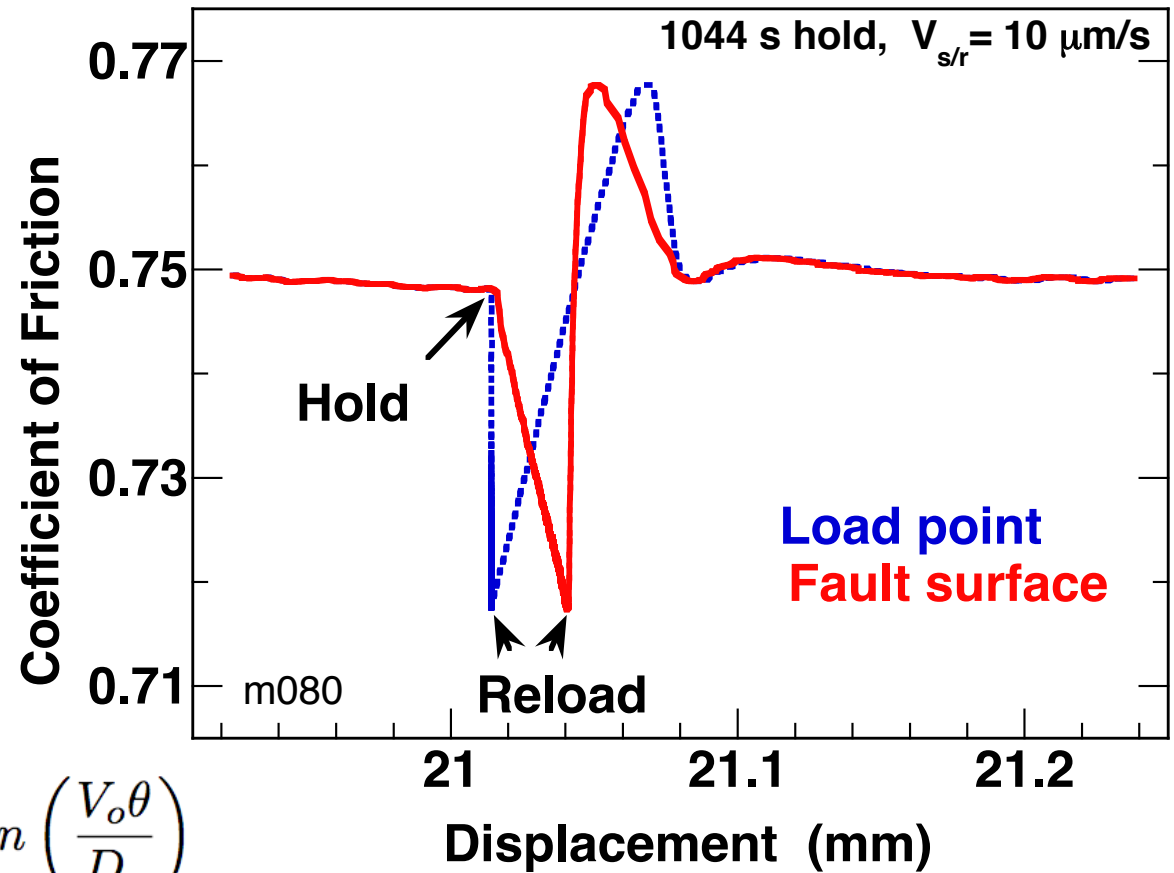
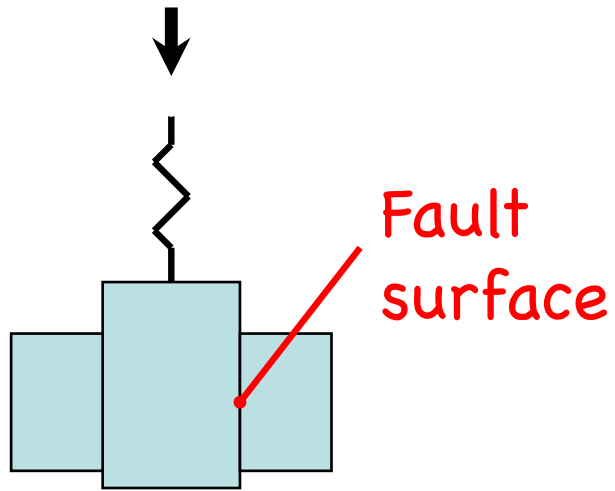


Vidale et al., 1994; Peng, Vidale, Marone & Rubin, *GRL* 2005



Sheared layer of quartz particles  
(100-150  $\mu\text{m}$ ), 25 MPa normal stress .  
Marone, 1998





$$1) \quad \mu(\theta, V) = \mu_o + a \ln \left( \frac{V}{V_o} \right) + b \ln \left( \frac{V_o \theta}{D_c} \right)$$

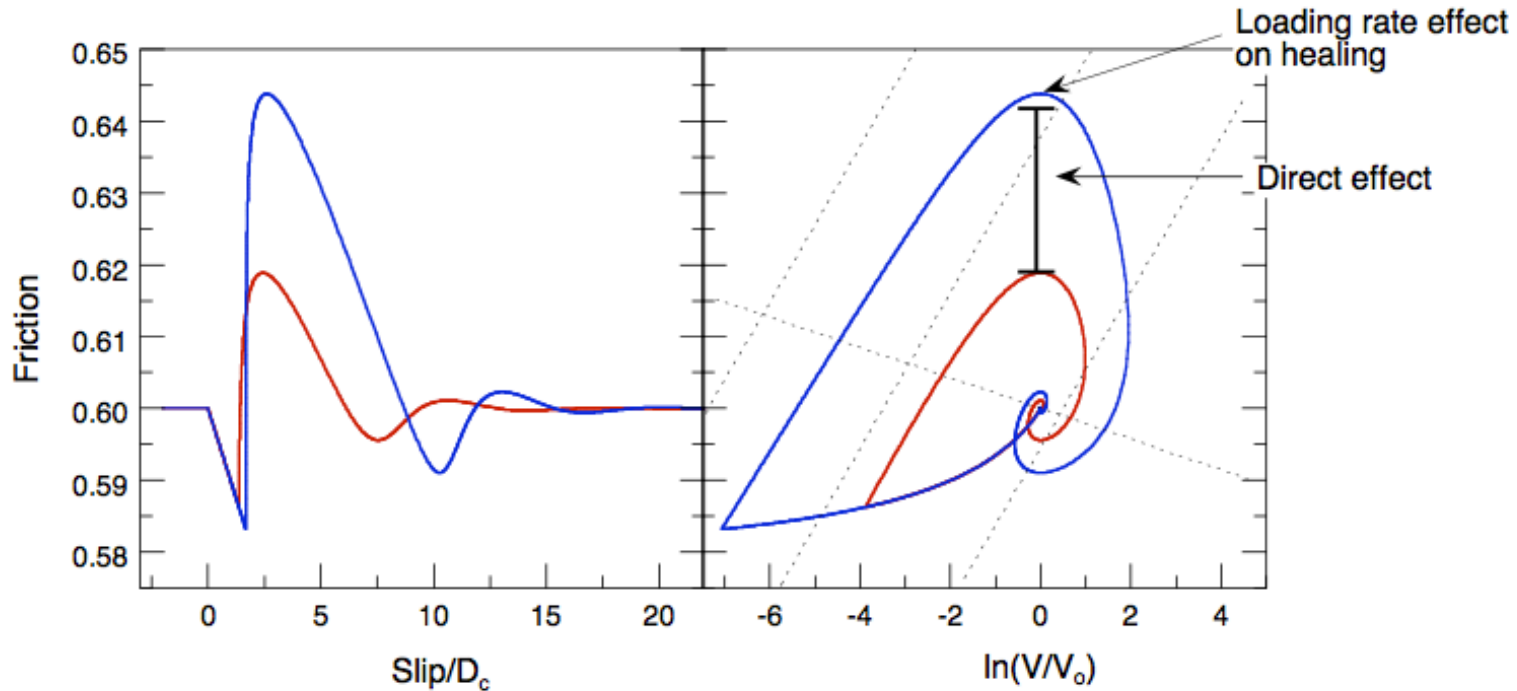
$$2) \quad \frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$3) \quad \frac{d\mu}{dt} = k(V_{lp} - V) \quad \text{Elastic Coupling}$$

$$\frac{d\mu}{dt} = k \left( V_{lp} - V_o \exp \left[ \frac{\mu - \mu_o - b \ln \left( \frac{V_o \theta}{D_c} \right)}{a} \right] \right)$$

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

# Derivation of the healing rate

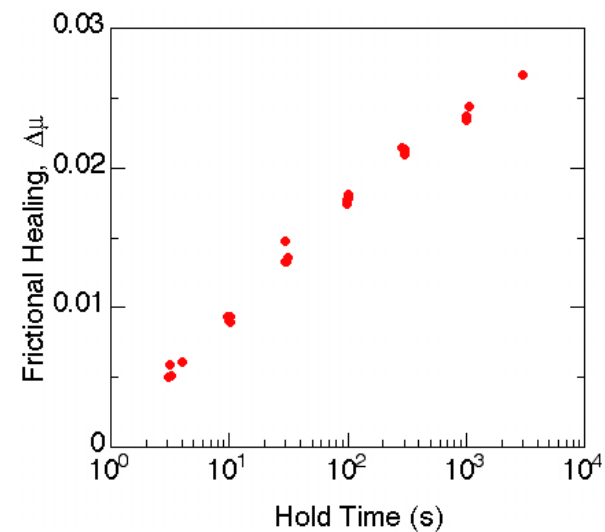


Loading rate effect on frictional healing is due to a combination of the friction direct effect and state evolution

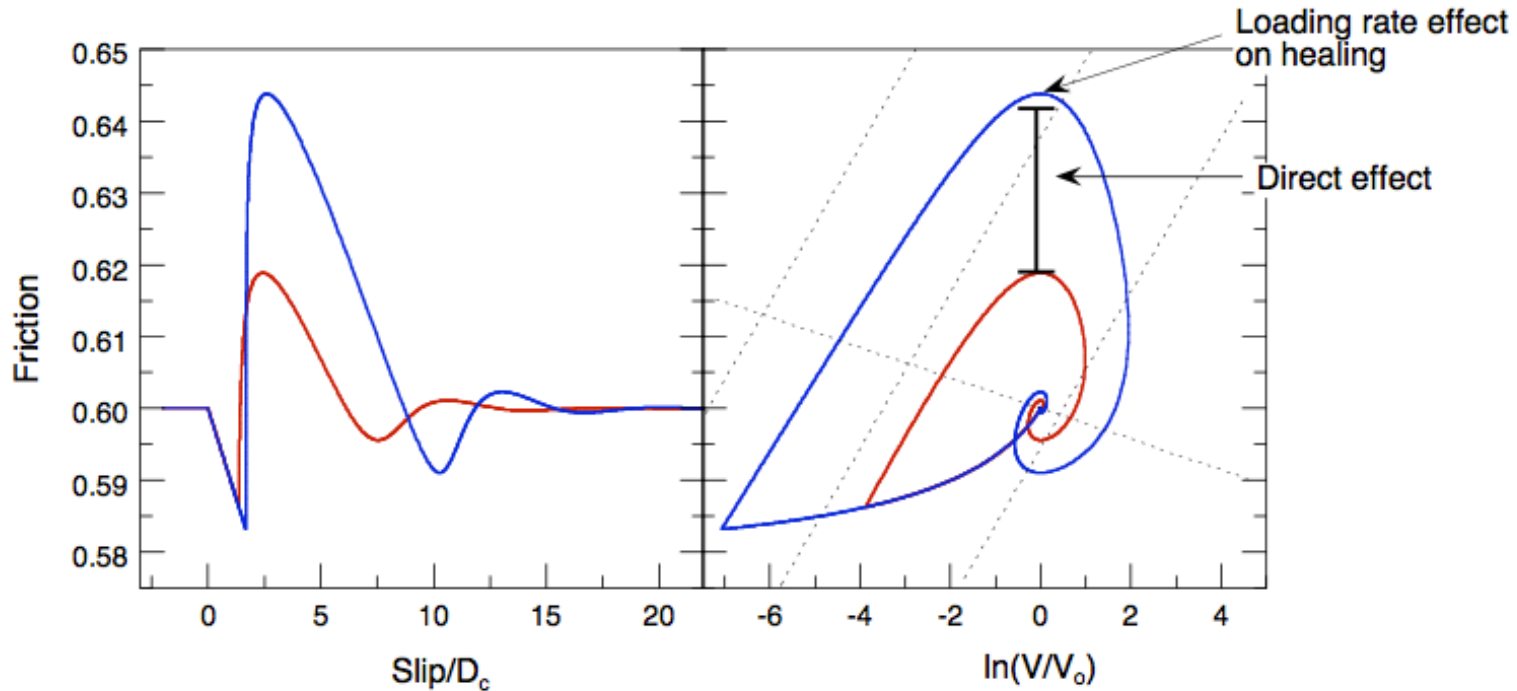
$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$\frac{d\mu}{dt} = k(V_{lp} - V)$$



# Derivation of the healing rate

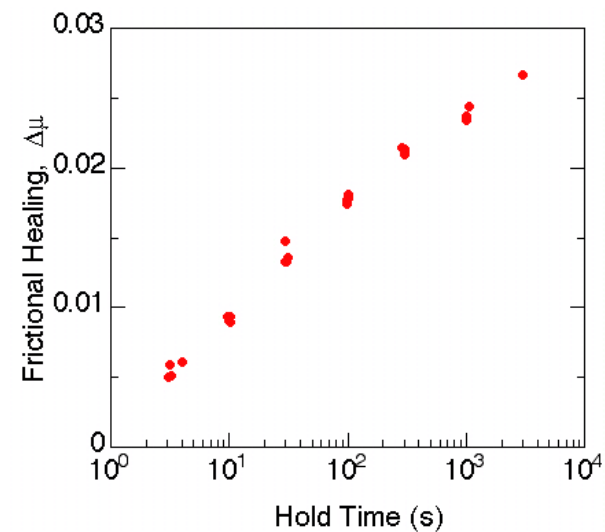


But let's focus on just one velocity, so that we can see what the healing rate tells us.

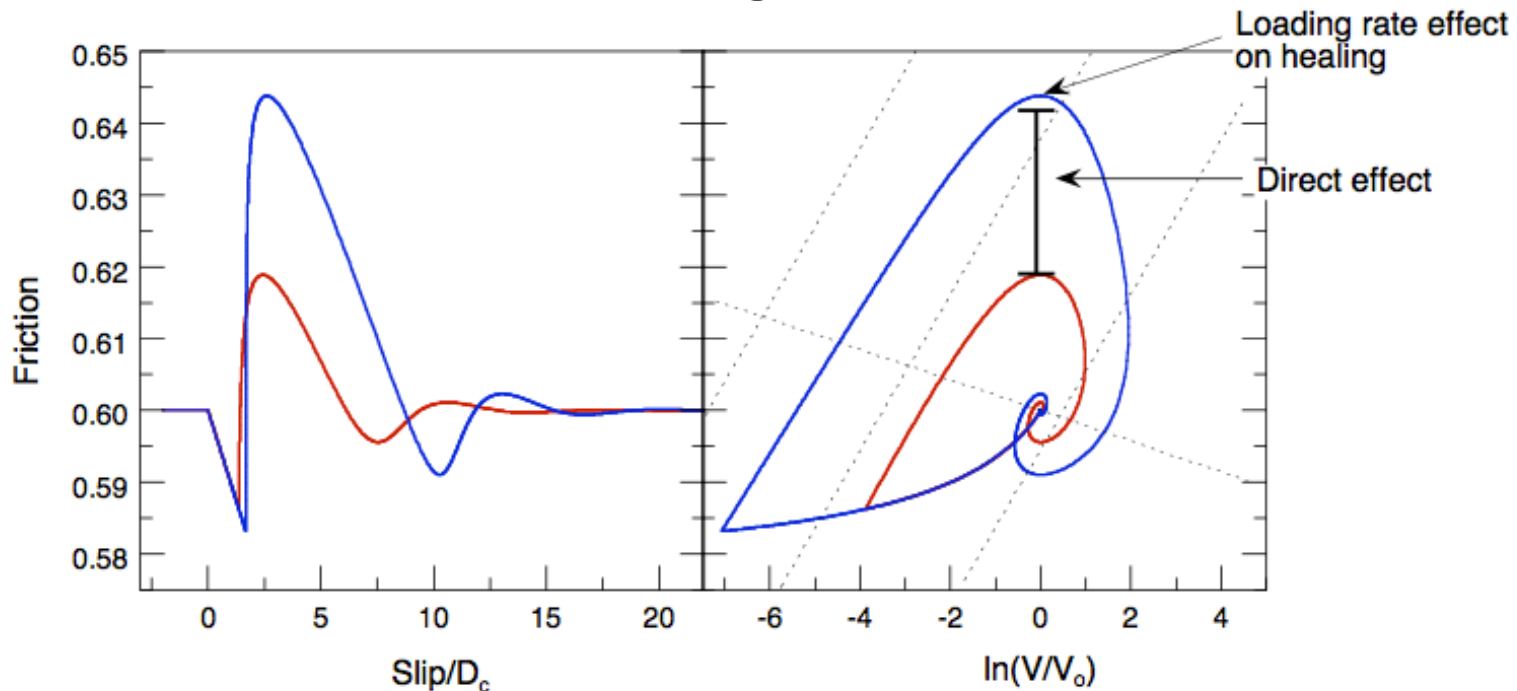
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$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$\frac{d\mu}{dt} = k(V_{lp} - V)$$



# Derivation of the healing rate



But let's focus on just one velocity, so that we can see what the healing rate tells us.

$$(1) \mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

$$(2) \frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$(3) \frac{d\mu}{dt} = k(V_{lp} - V)$$

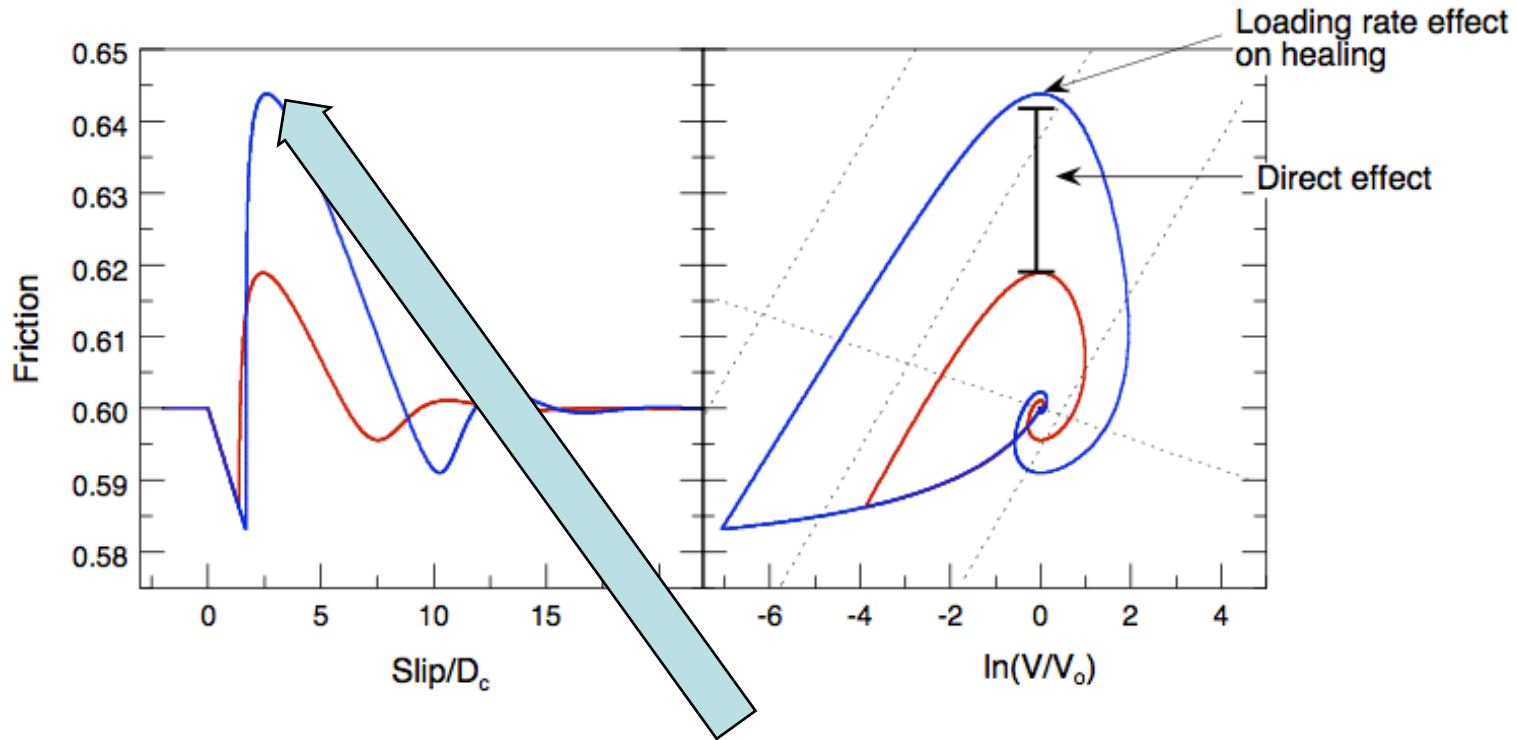
The healing rate involves measurements of  $\Delta\mu$   
-look back at notes from Lec. 7

At steady state, we have

$$\mu_1 - \mu_o = (a - b) \ln\left(\frac{V_1}{V_o}\right)$$

but for the peak ("static") friction, we need to think about both terms on the RSH of eq'n 1

# Derivation of the healing rate



What do we know about  $V$  at the peak?

$V = V_o$  at the peak --right, that's what we see in the phase plane.

Hmmm, that means that the  $\ln(V/V_o)$  term is zero

So  $\Delta\mu = b \ln(V_o \theta / D_c)$

and how does  $\theta$  vary with SHS time?

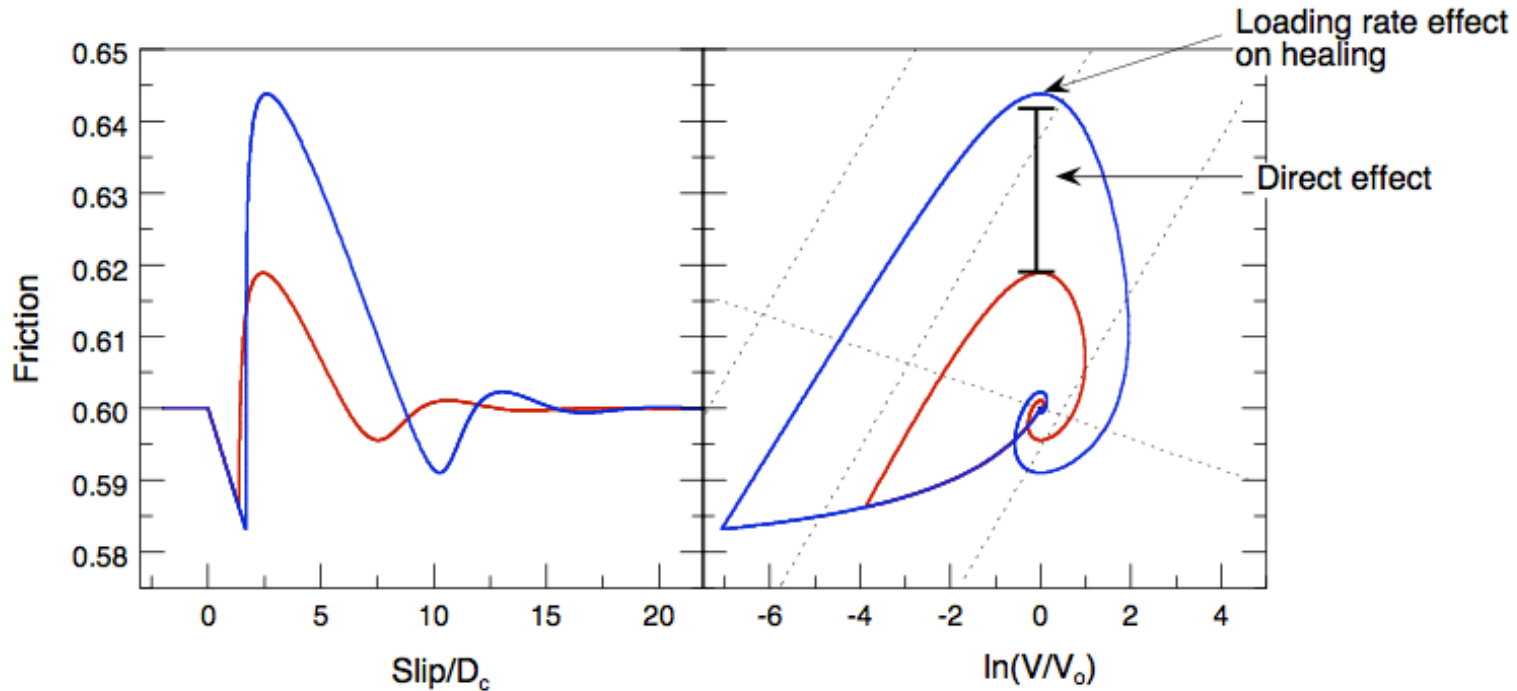
according to (2), as  $V$  goes to 0,  $d\theta = dt$

$$(1) \mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

$$(2) \frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$(3) \frac{d\mu}{dt} = k(V_{lp} - V)$$

# Derivation of the healing rate

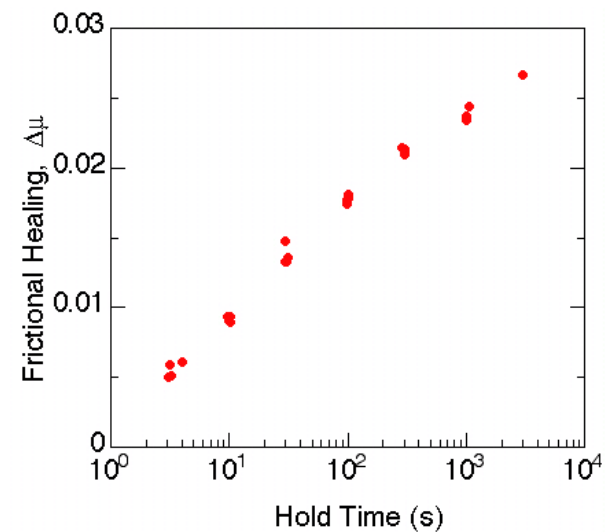


so, if we plot  $\Delta\mu$  vs  $\ln(t)$  we should have the healing rate

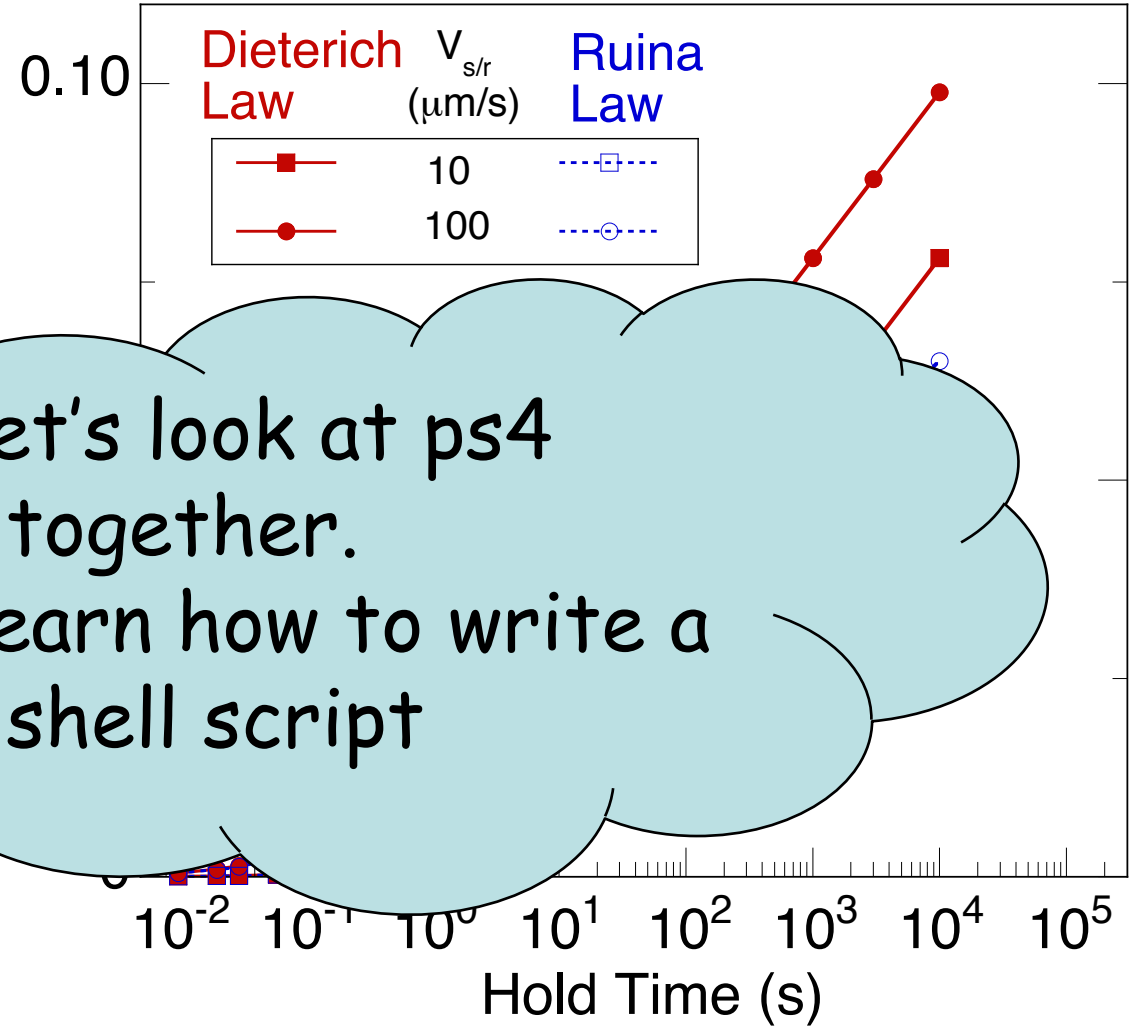
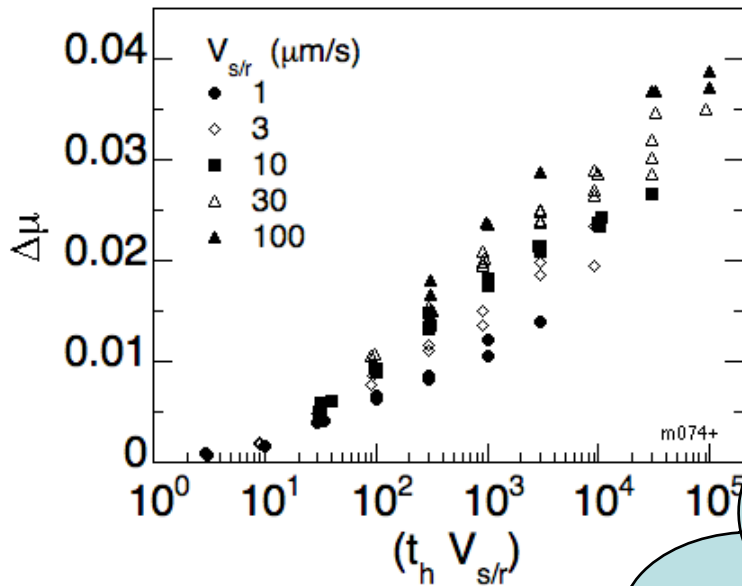
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$$(3) \frac{d\mu}{dt} = k(V_{lp} - V)$$







Let's look at ps4 together.  
Learn how to write a shell script

**Friction Law**

$$\mu = \mu_o + a \ln(V/V_o) + b \ln(V_o \theta/D_c)$$

**State Evolution**

$$d\theta/dt = 1 - V \theta/D_c$$

$$d\theta/dt = -V \theta/D_c \ln(V \theta/D_c)$$

**Elastic Coupling**

$$d\mu/dt = k(V_{ip} - V)$$

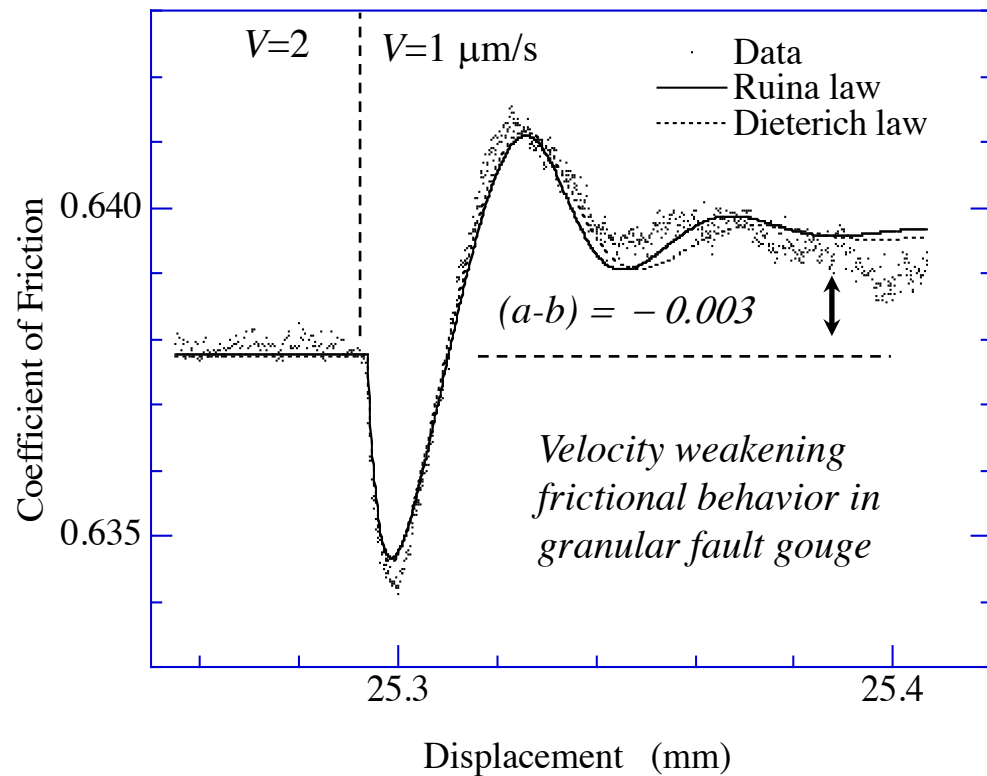
The rate of frictional healing depends on the rate of shearing (Marone, 1998, *Nature*)

Rate State Friction Laws predict this behavior

# Rate/State Friction

## Measuring the friction constitutive parameters

*Empirical laws, based on laboratory friction data*



### Constitutive Modelling

#### Rate and State Friction Law

#### Elastic Interaction, Testing Apparatus

$$\mu(\theta, v) = \mu_0 + a \ln\left(\frac{v}{v_0}\right) + b \ln\left(\frac{v_0 \theta}{D_c}\right)$$

$$\frac{d\theta}{dt} = 1 - \frac{v\theta}{D_c}$$

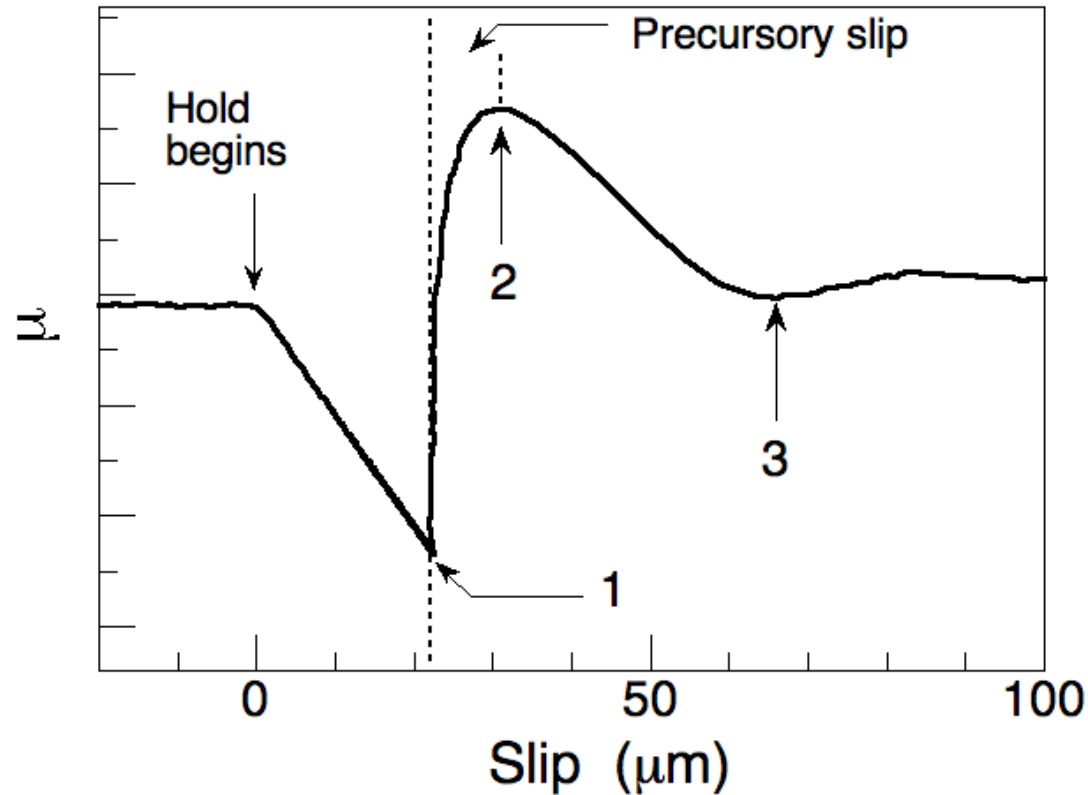
$$\theta_{ss} = \frac{D_c}{v}$$

$$\Delta\mu_{ss} = (a-b) \ln\left(\frac{v}{v_0}\right)$$

$$\frac{d\mu}{dt} = k' \left( v_{lp} - v \right)$$

# Rate/State Friction

## Measuring the friction constitutive parameters



### Constitutive Modelling

#### Rate and State Friction Law

#### Elastic Interaction, Testing Apparatus

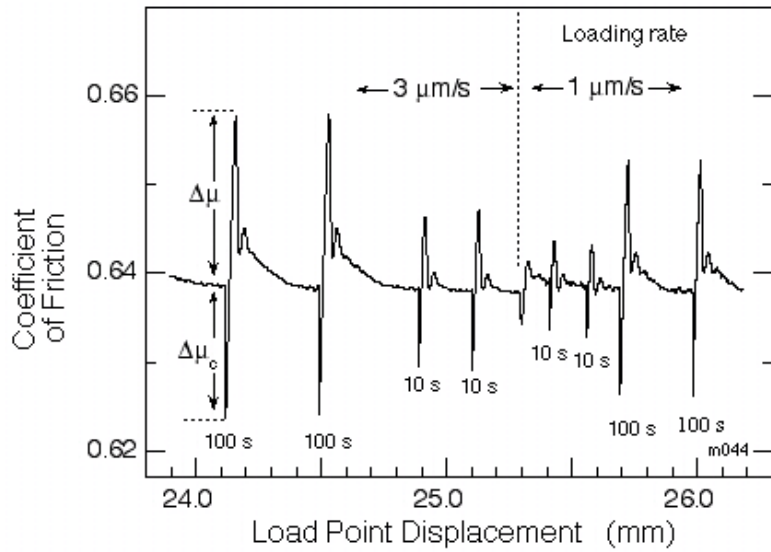
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$$\frac{d\theta}{dt} = 1 - \frac{v\theta}{D_c}$$

$$\theta_{ss} = \frac{D_c}{v}$$

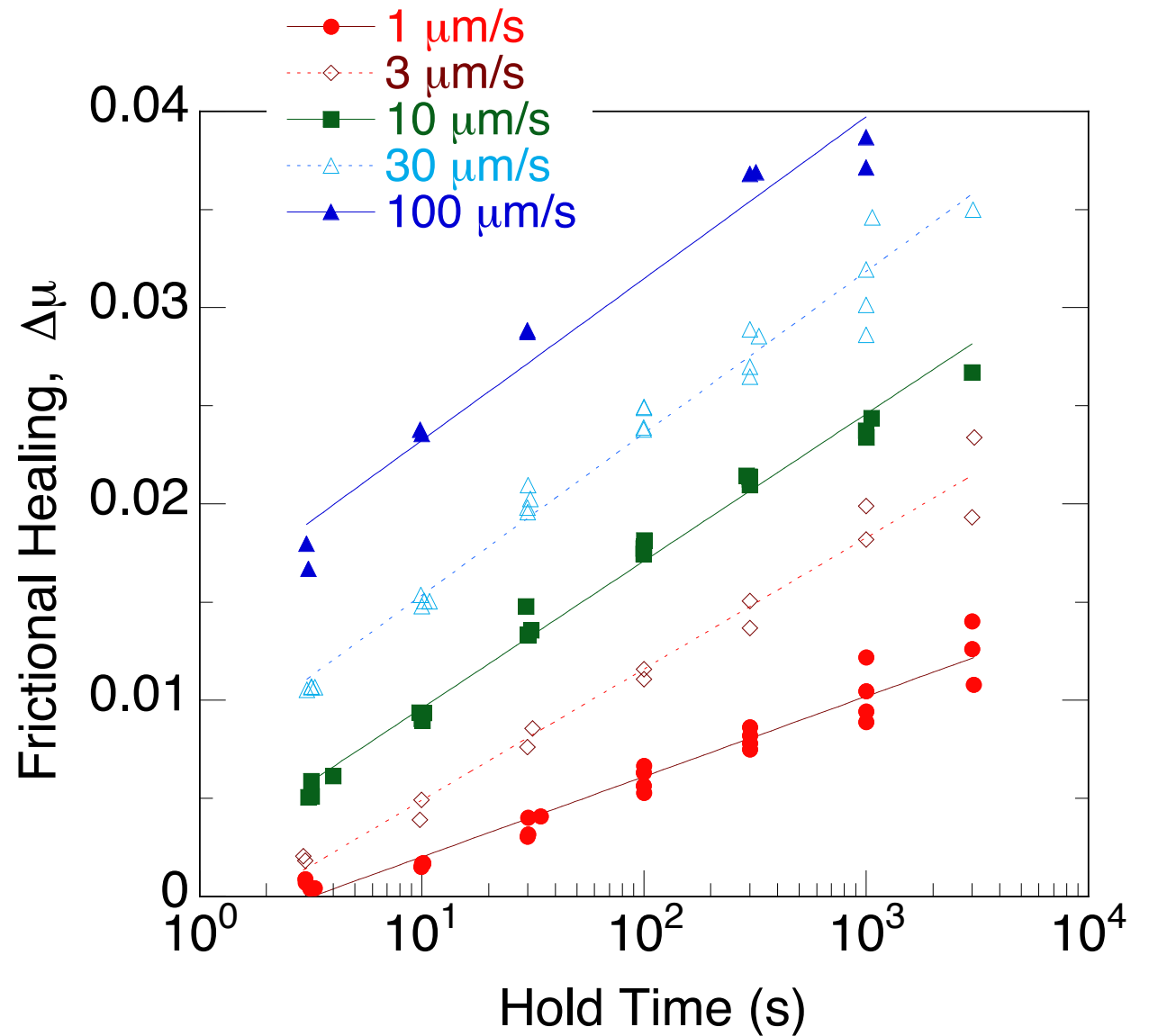
$$\Delta\mu_{ss} = (a - b) \ln\left(\frac{v}{v_0}\right)$$

$$\frac{d\mu}{dt} = k' (v_{lp} - v)$$



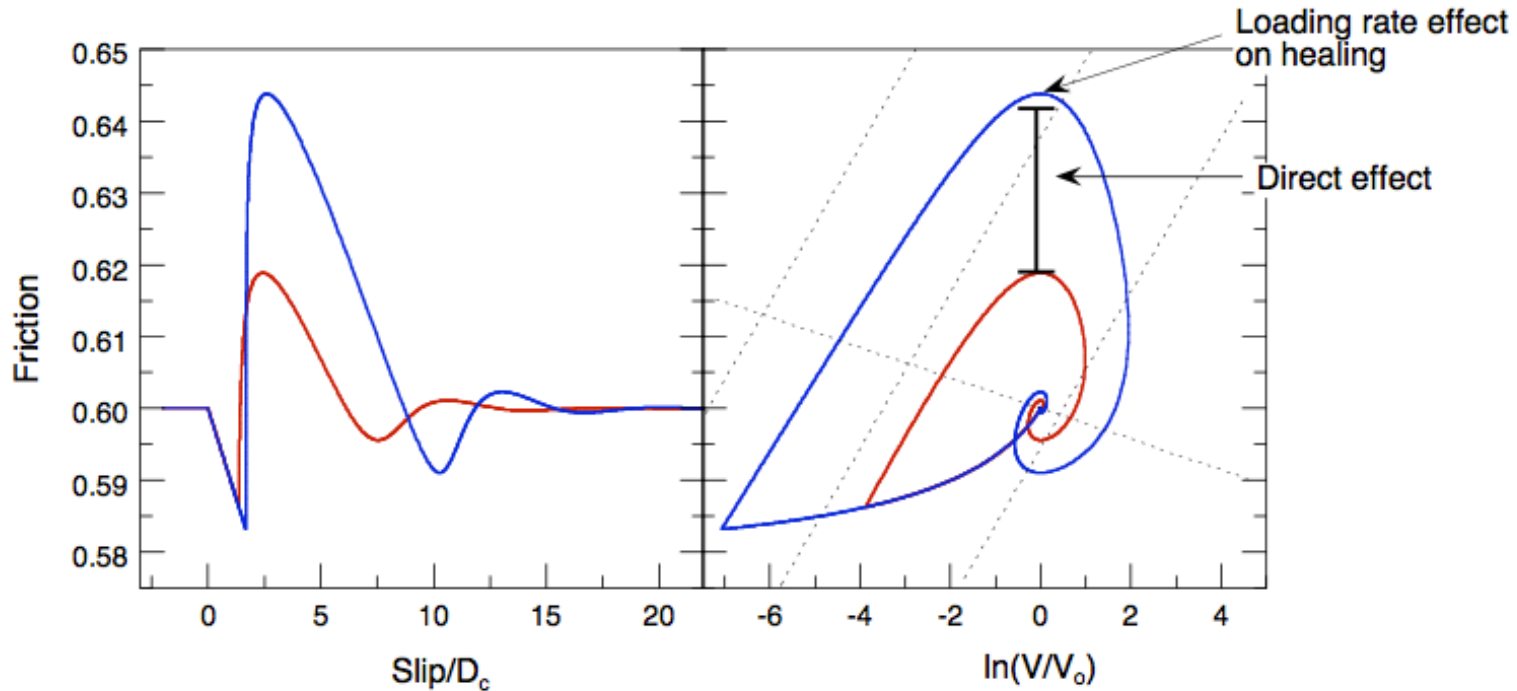
## Stressed Aging

Aging rate depends on the rate of shearing



(Marone, 1998, *Nature*)

# Phase Plane Plots



Loading rate effect on frictional healing is due to a combination of the friction direct effect and state evolution

$$\mu(\theta, V) = \mu_o + a \ln\left(\frac{V}{V_o}\right) + b \ln\left(\frac{V_o \theta}{D_c}\right)$$

$$\frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$\frac{d\mu}{dt} = k(V_{lp} - V)$$

shs test:

1  $\mu\text{m/s}$

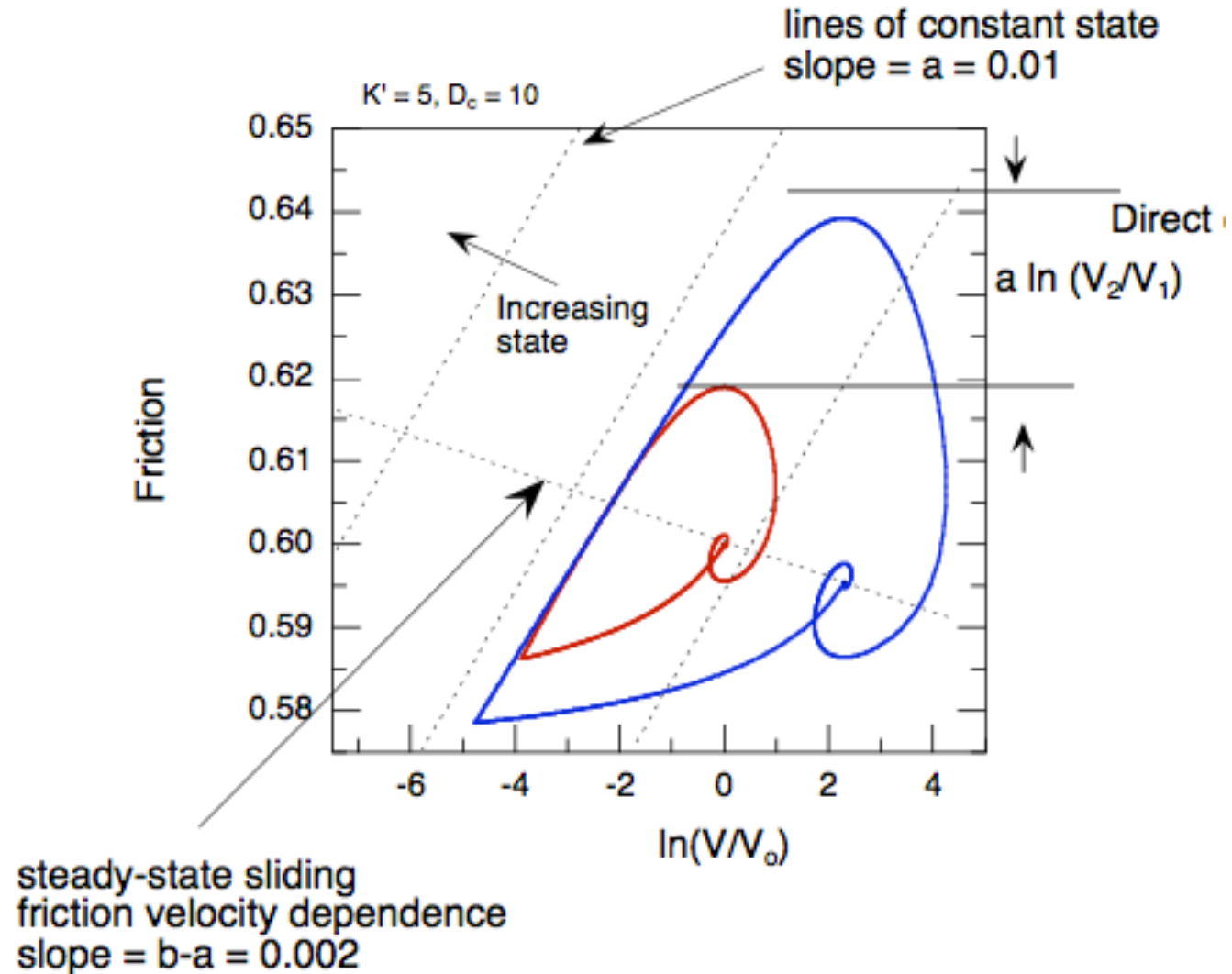
10  $\mu\text{m/s}$

# Phase Plane Plots

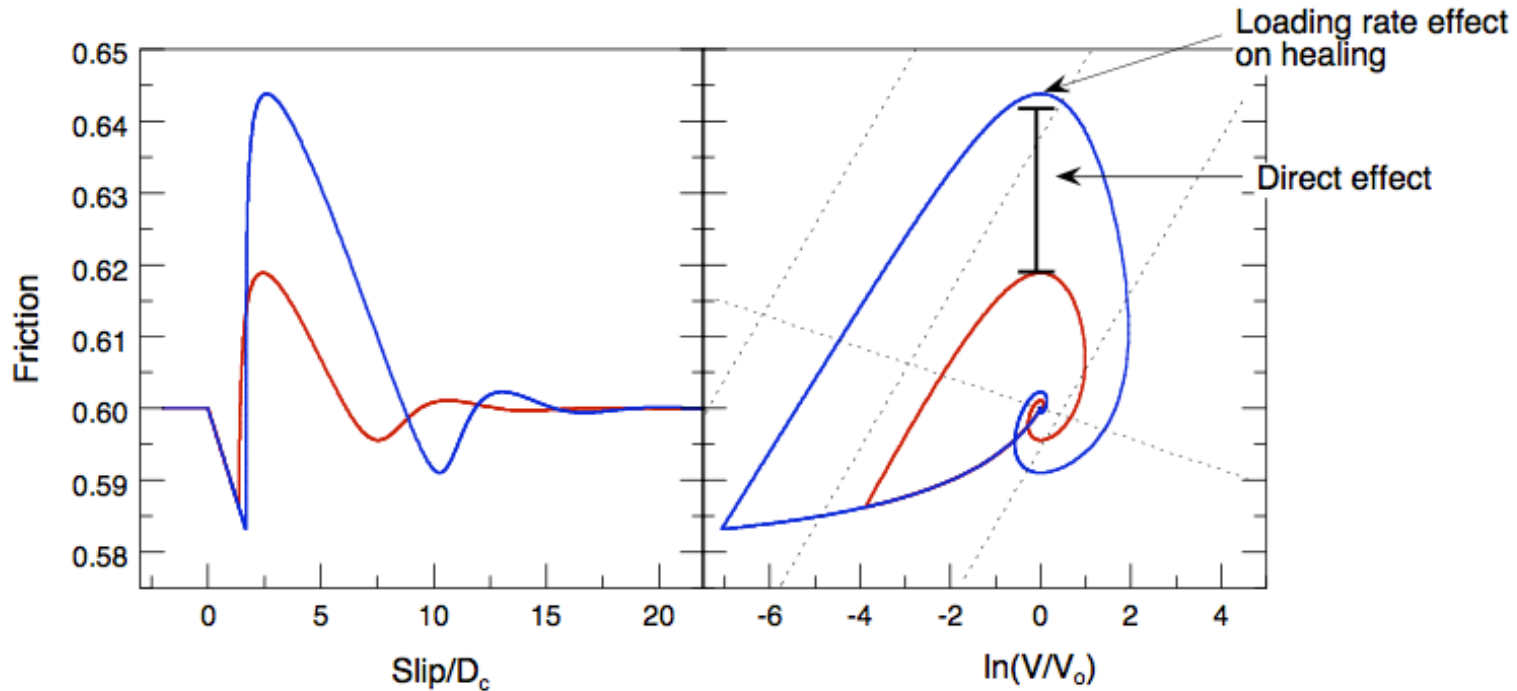
shs test:

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10  $\mu\text{m/s}$



# Derivation of the healing rate

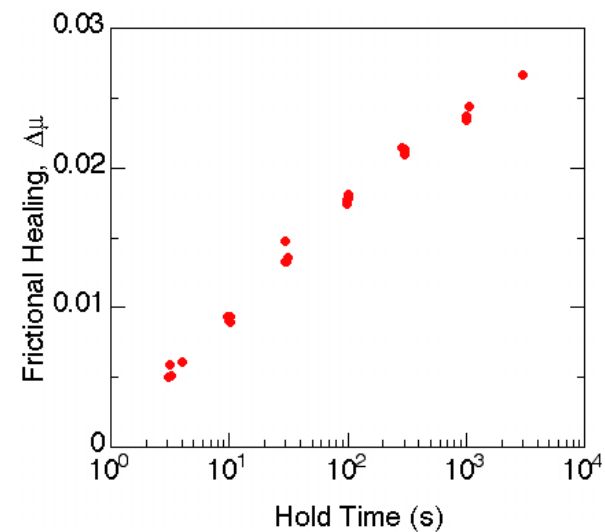


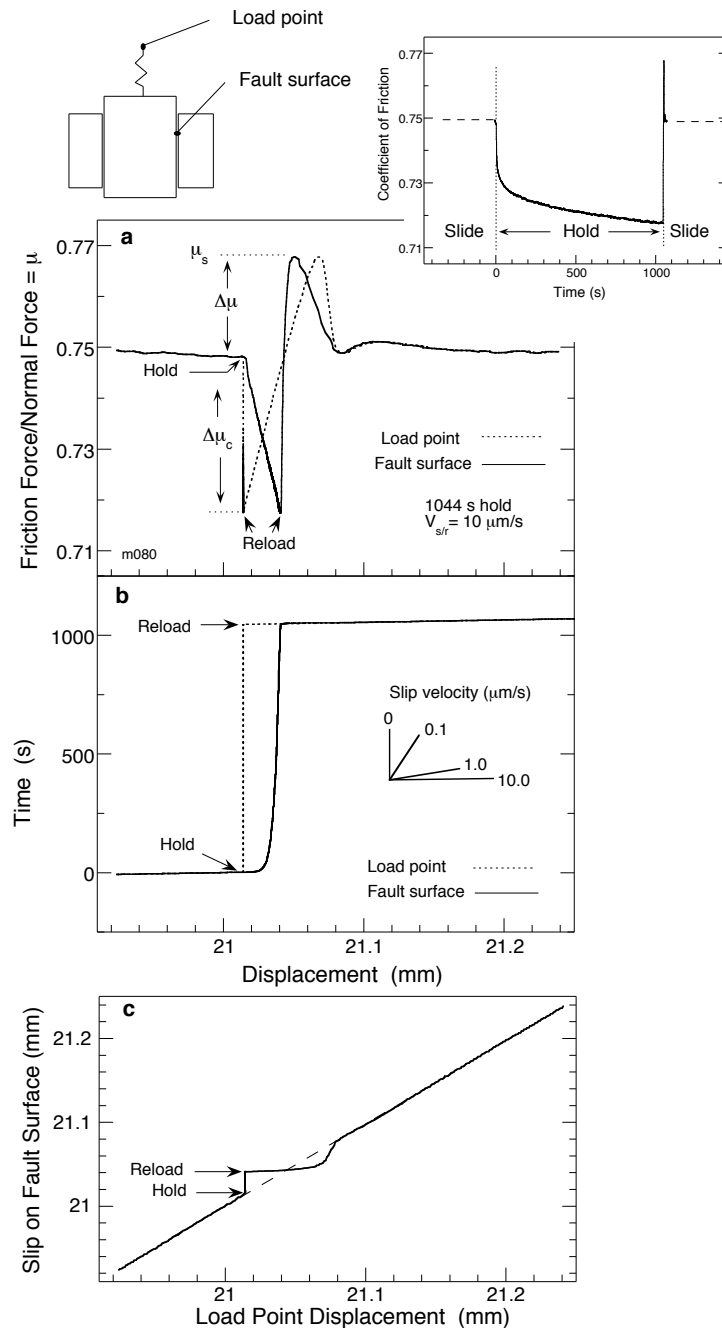
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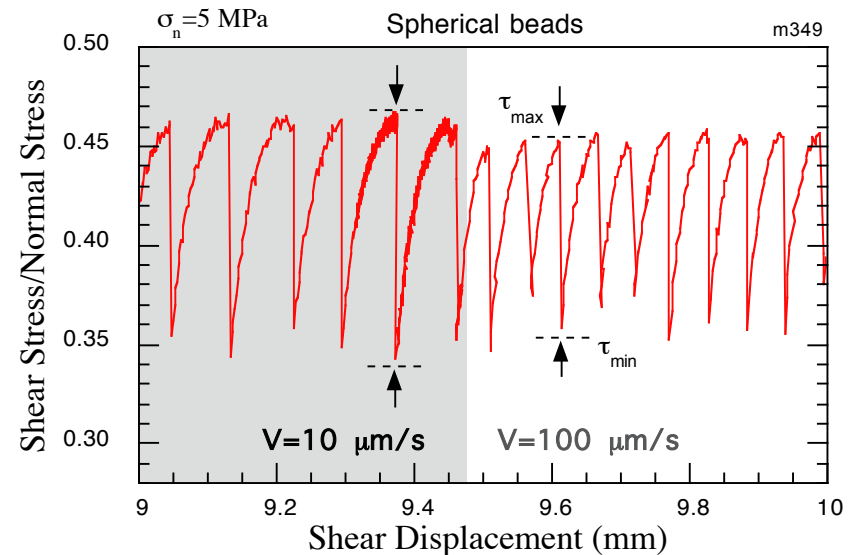


Slide-hold-slide

Slip-reload-slip

Earthquake-interseismic healing and reloading-earthquake

The full seismic cycle of stick-slip, frictional restrengthening, and interseismic reloading



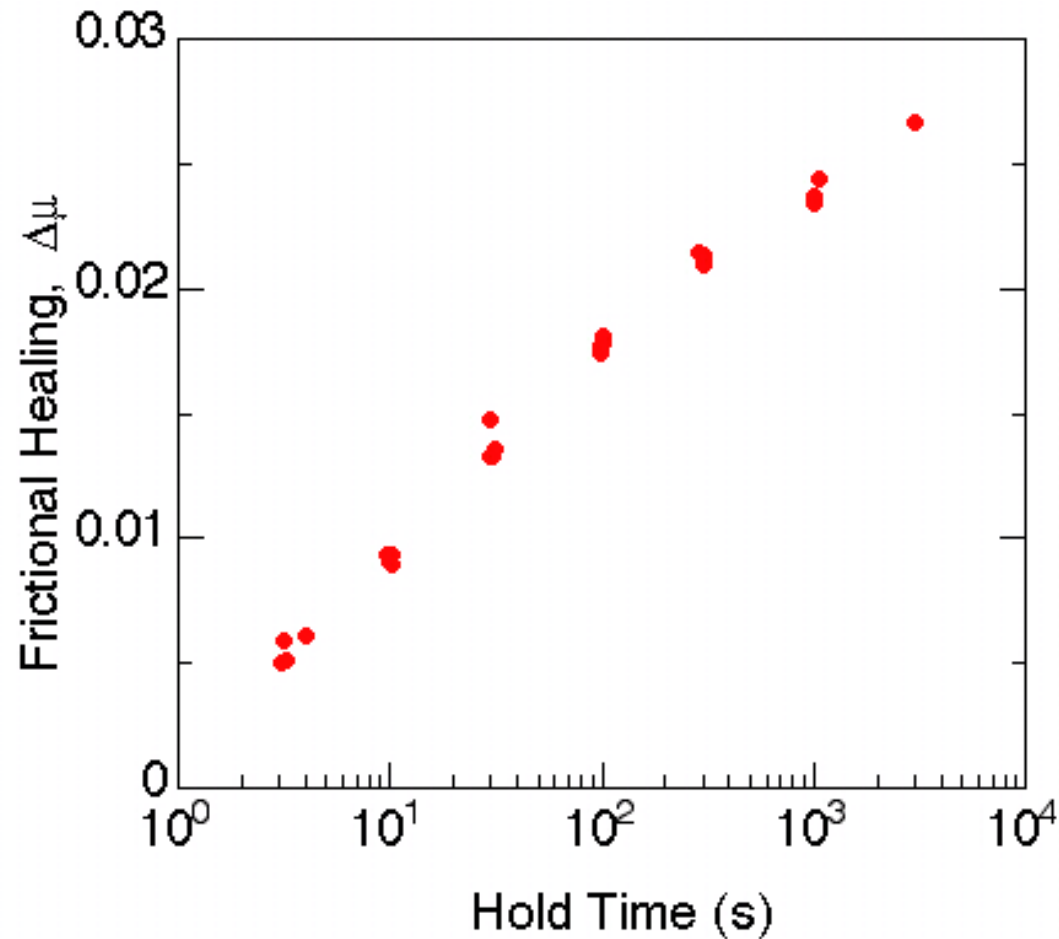


## Rate ( $v$ ) and State ( $\theta$ ) Friction Constitutive Laws

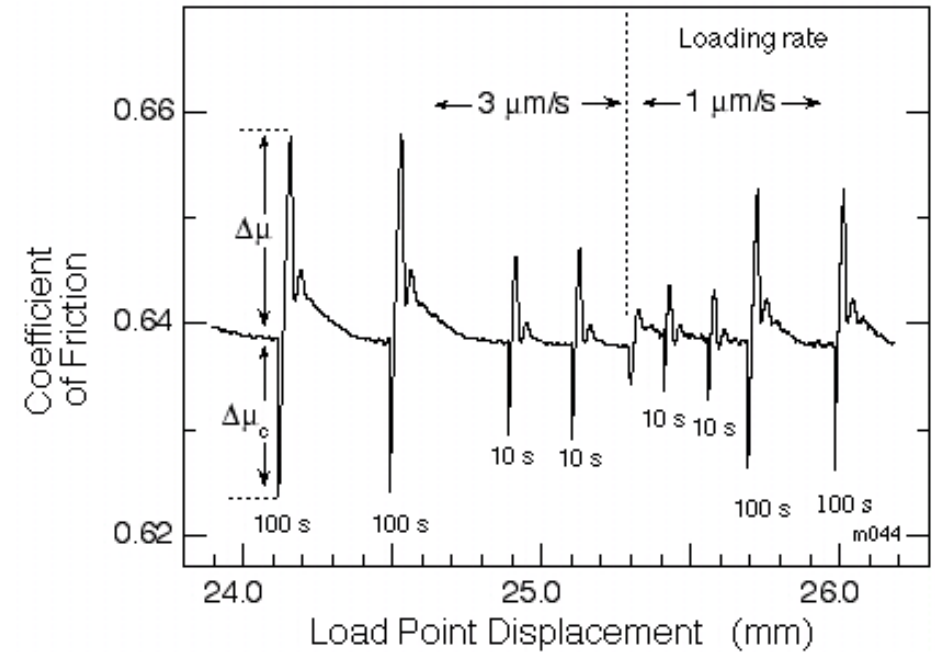
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$$2) \quad \frac{d\theta}{dt} = 1 - \frac{V\theta}{D_c}$$

$$3) \quad \frac{d\mu}{dt} = k(V_{lp} - V) \quad \text{Elastic Coupling}$$



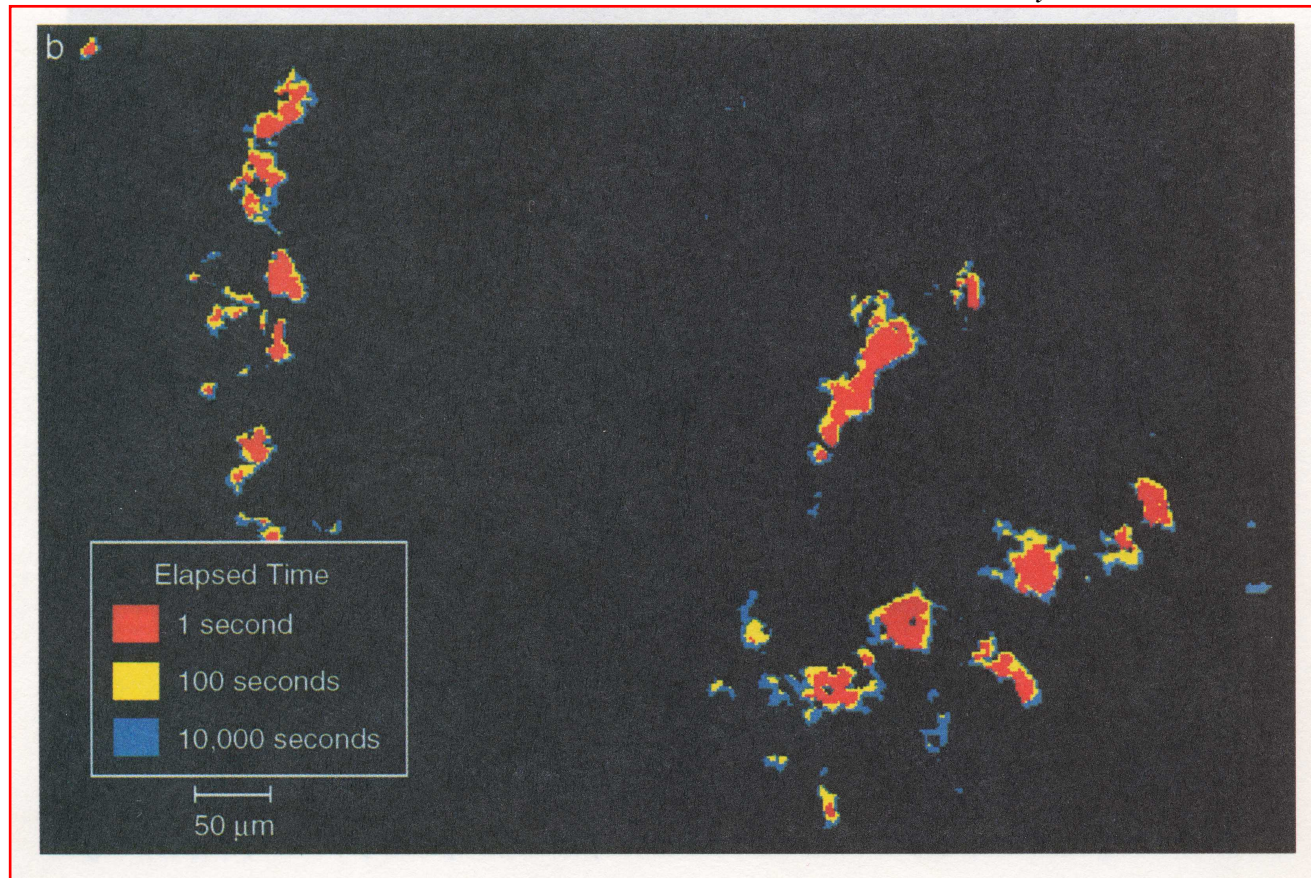
## Modeling experimental data



$$\frac{d\mu}{dt} = k(V_{lp} - V)$$

Time dependent yield strength:

$$\mu = \frac{\tau}{\sigma_n} = \frac{S}{\sigma_y}$$



Dieterich and Kilgore [1994]

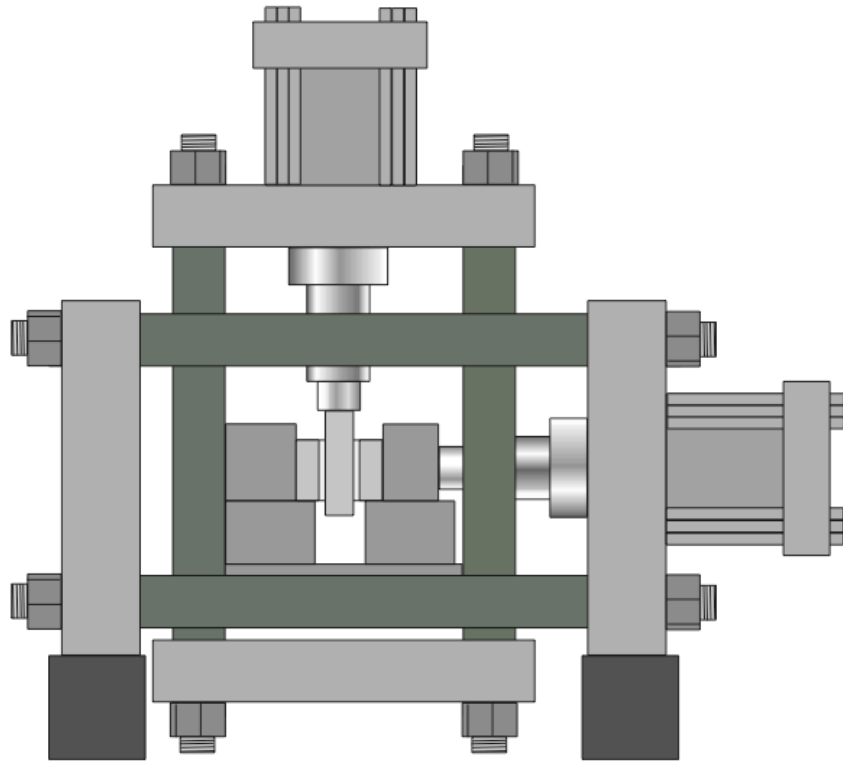
$$\mu = \frac{\tau}{\sigma_n} = \frac{S}{\sigma_y}$$

Time dependent growth of contact (acrylic plastic)- true static contact

$$\sigma_y = \sigma_o + f(t)$$

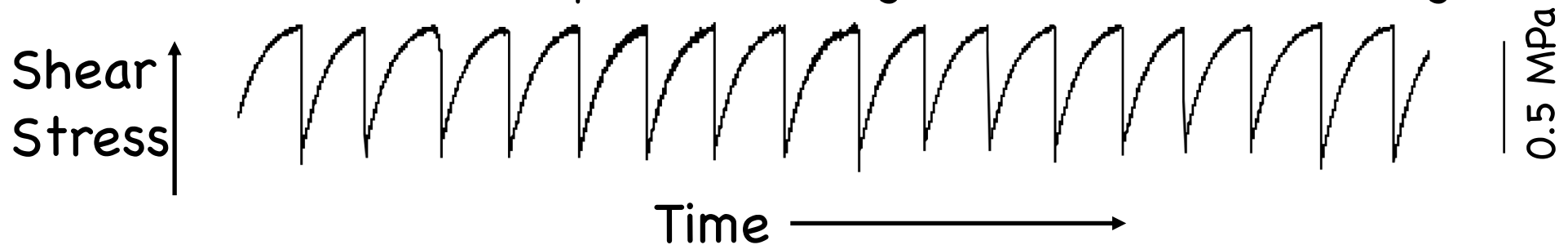
Modified from Beeler, 2003

# How do fault/frictional surfaces heal (regain strength) after failure?



Earthquakes & Fault Mechanics:  
seismic cycle, fault reactivation.  
*(friction and stick slip: doors,  
windows, machines, ships in dry  
dock, dancers...)*

Stick-slip failure during shear at constant loading rate



# Time dependence of “static” friction

## Aging of frictional contacts

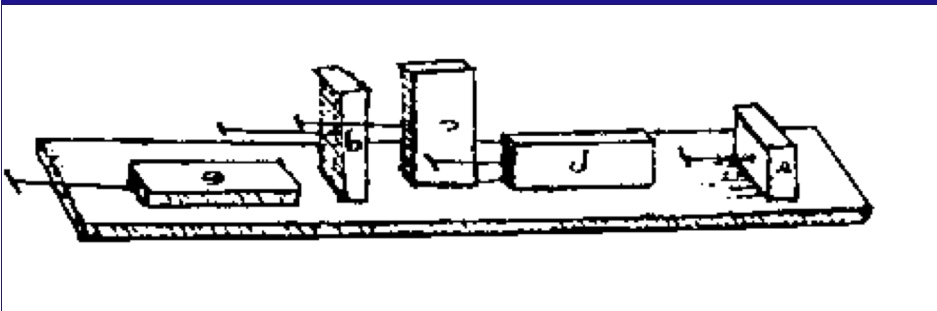


C. A. Coulomb (1736-1806)

Table 9.1

	$T$ (time of repose, min)	$A + mT^m$ (static friction force, lbf)
I <sup>st</sup> observation	0	$A = 502$
II <sup>c</sup>	2	790
III <sup>c</sup>	4	866
IV <sup>c</sup>	9	925
V <sup>c</sup>	26	1,036
VI <sup>c</sup>	60	1,186
VII <sup>c</sup>	960	1,535

static friction of two pieces of well-worn oak lubricated with tallow.



# Time dependence of “static” friction

## Aging of frictional contacts

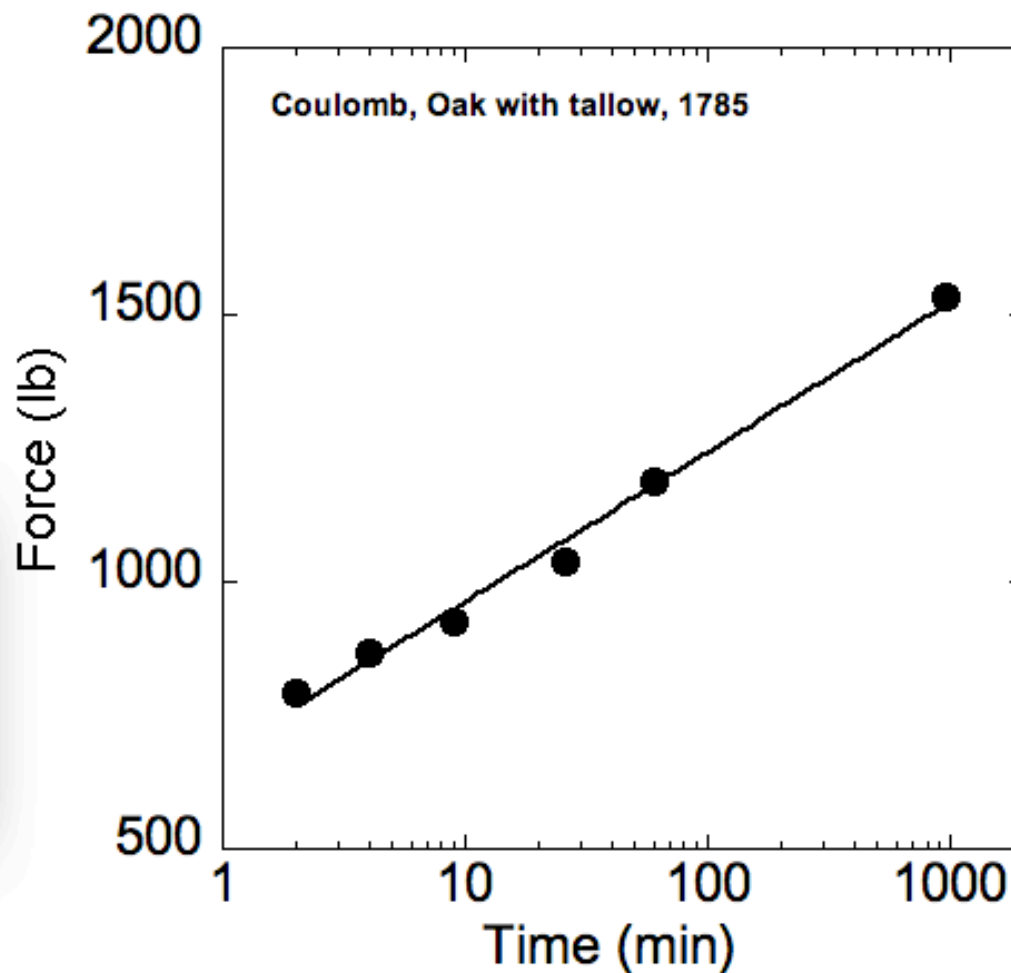


C. A. Coulomb (1736-1806)

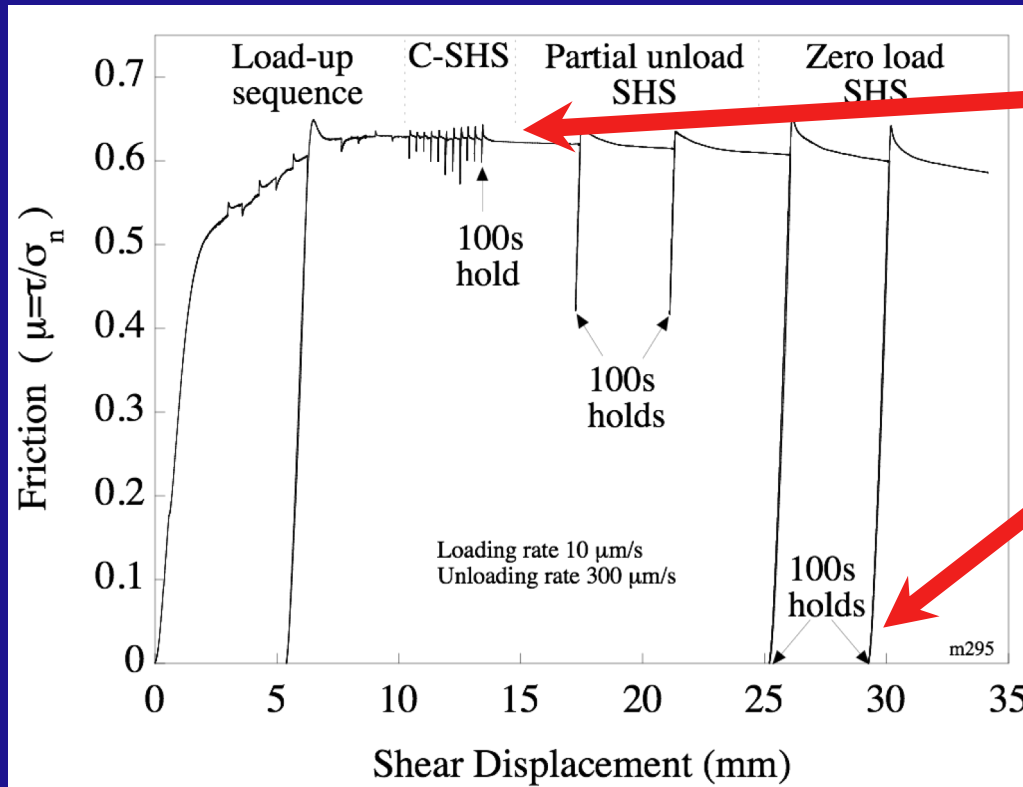
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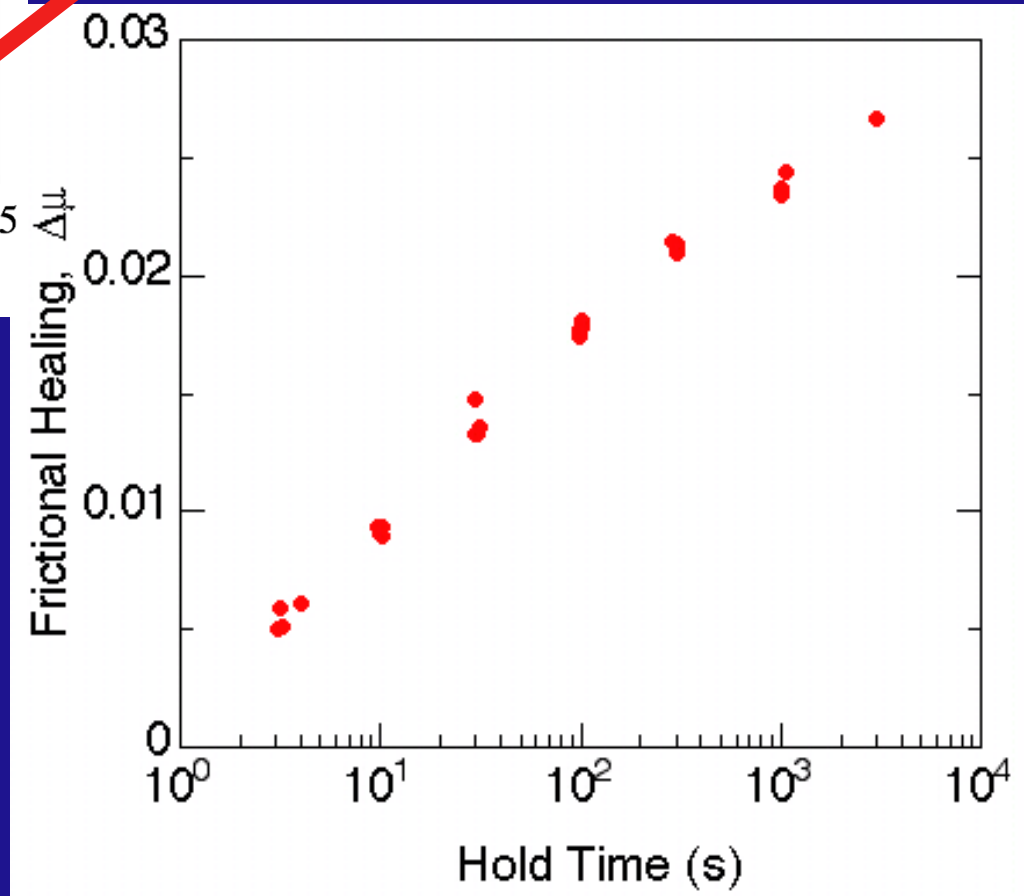


# Time dependence of friction in rocks; Macroscopic frictional aging



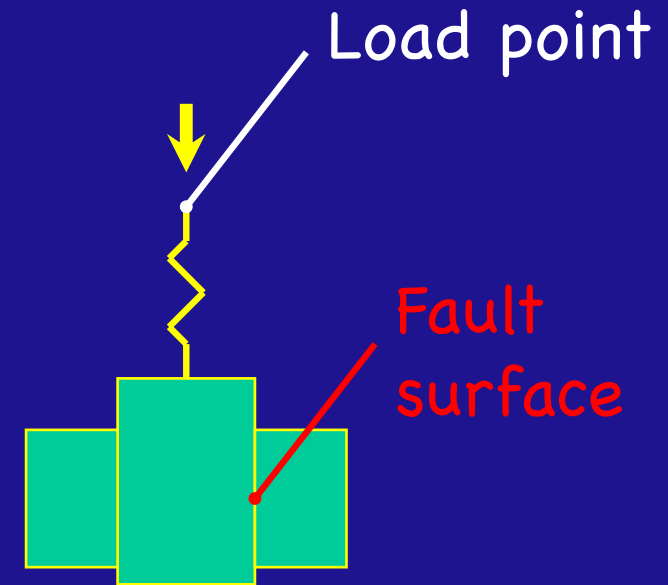
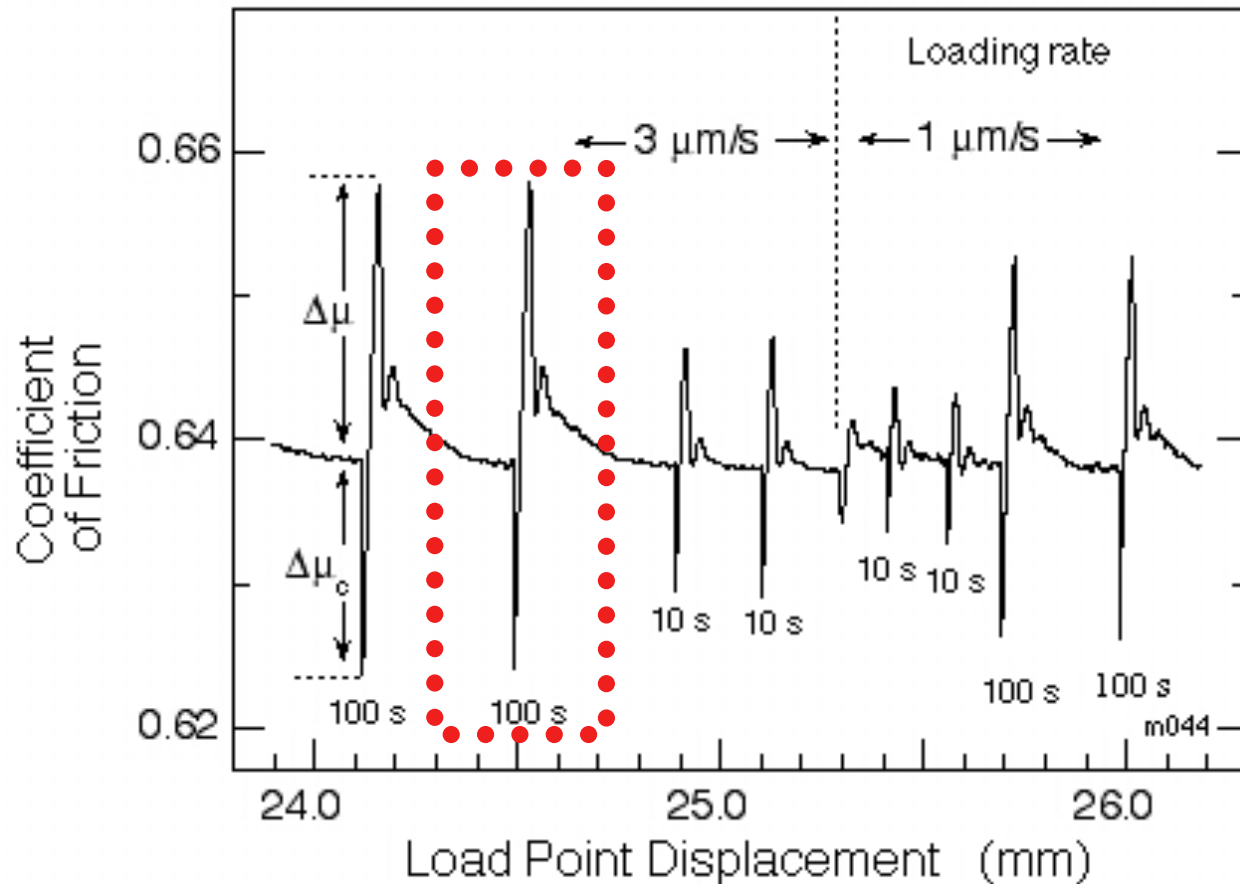
Stressed vs. unstressed aging

Karner and Marone, *J. Geophys. Res.*, 2001.



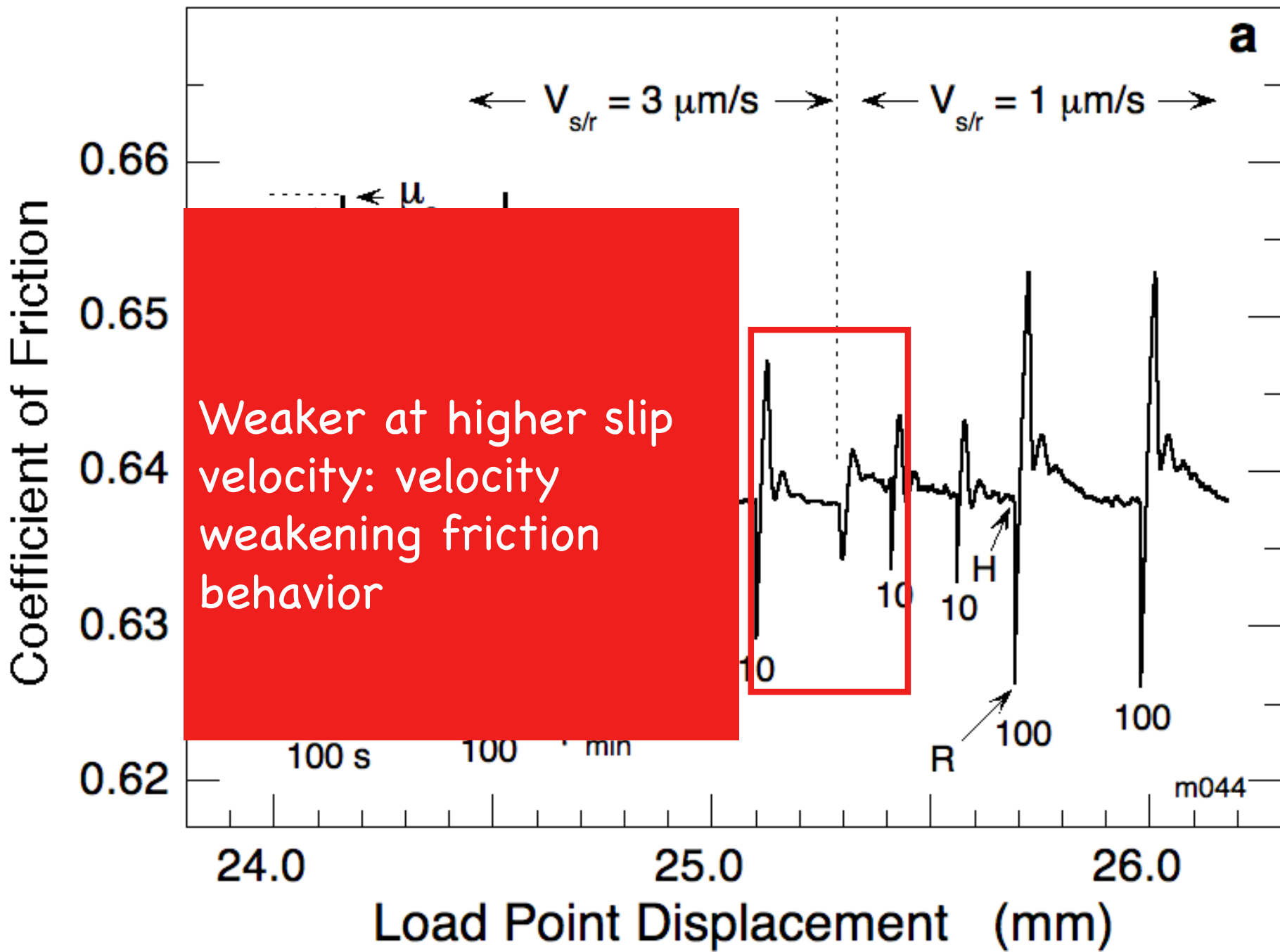
# Frictional Healing

## Stressed aging



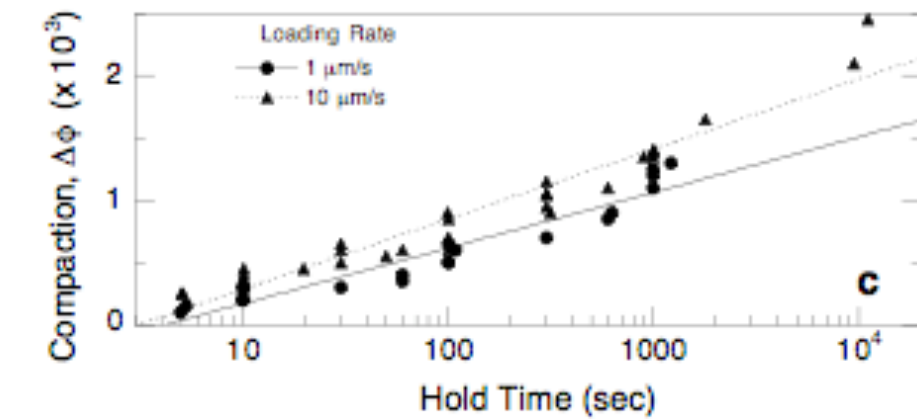
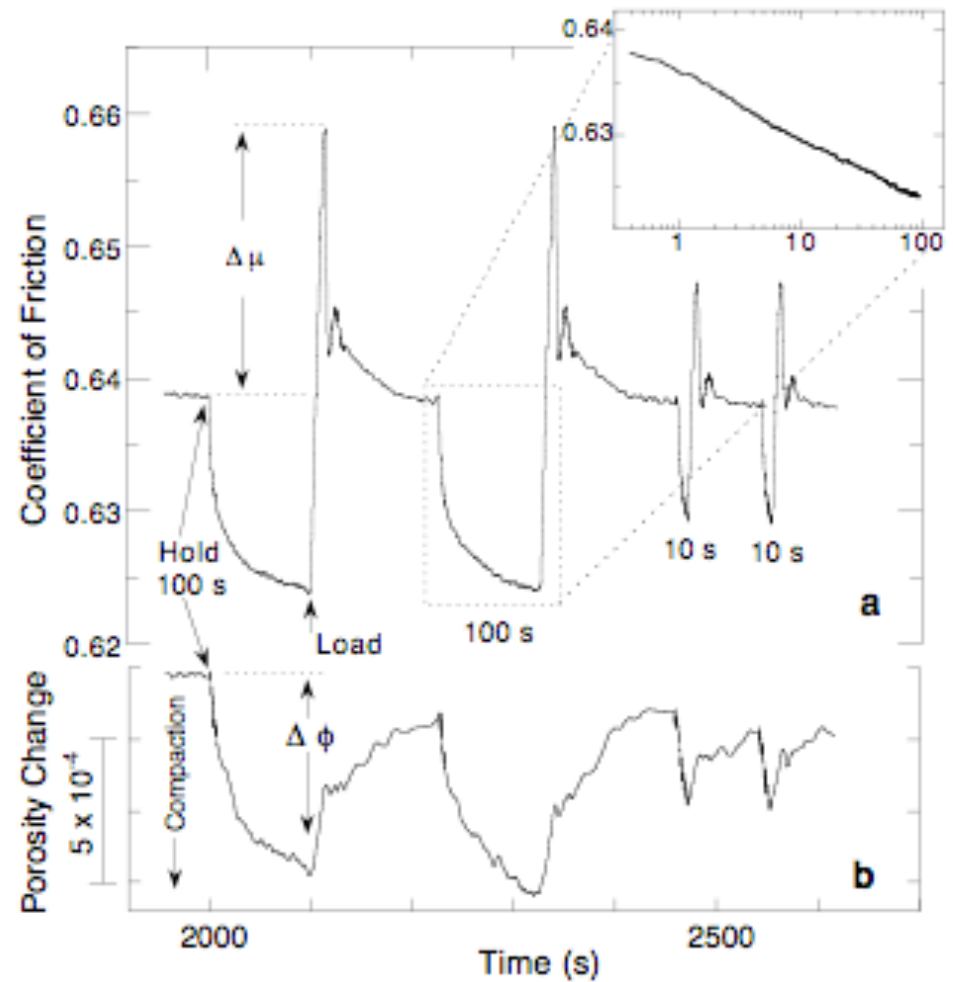
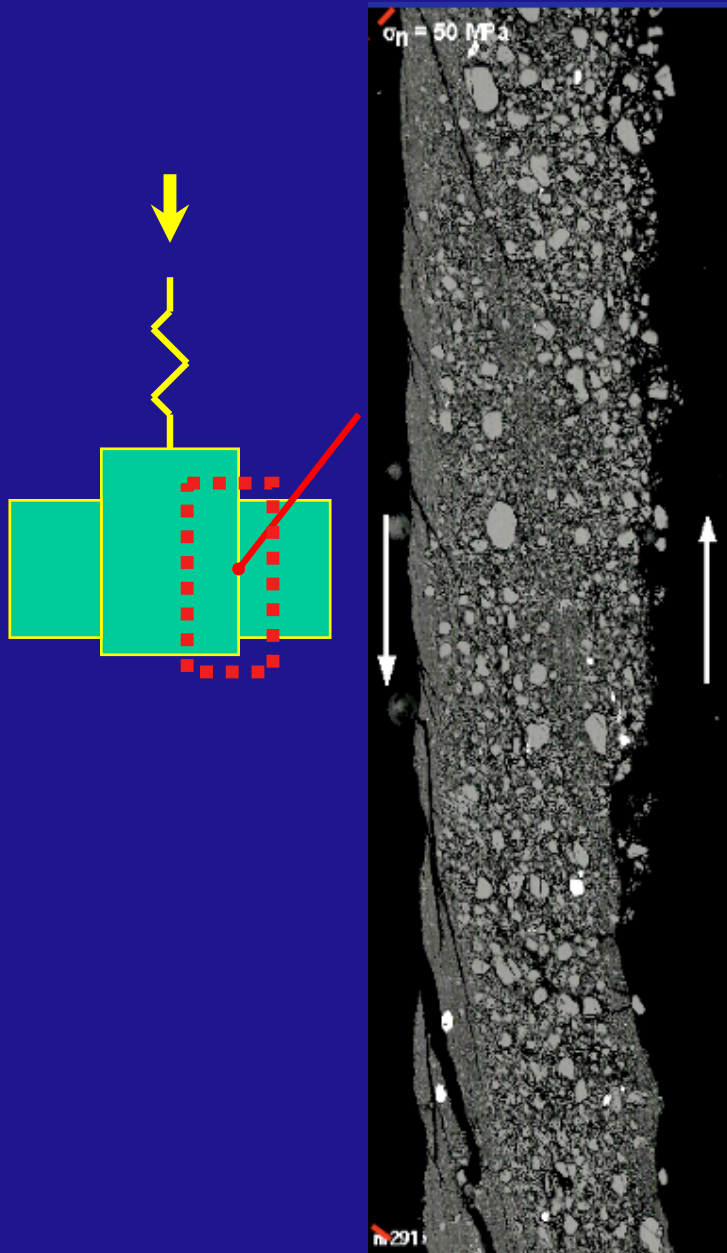
Steady state friction & the rate of healing vary with sliding velocity

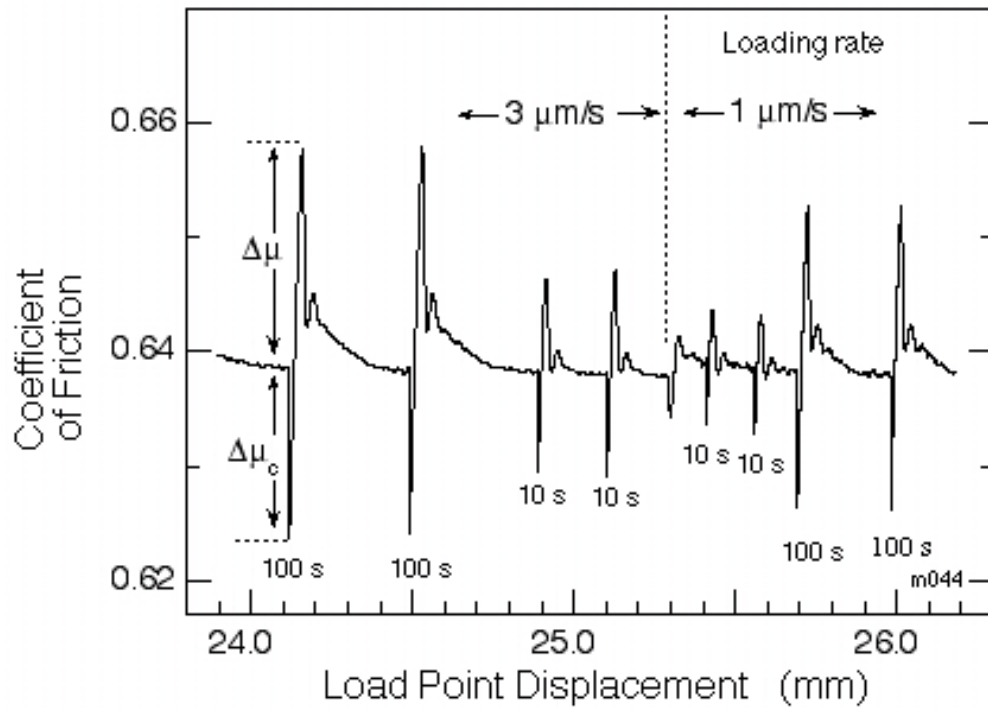
Angular quartz particles (100-150  $\mu\text{m}$ ), 3 mm thick, 25 MPa normal stress. Marone, Nature, 1998



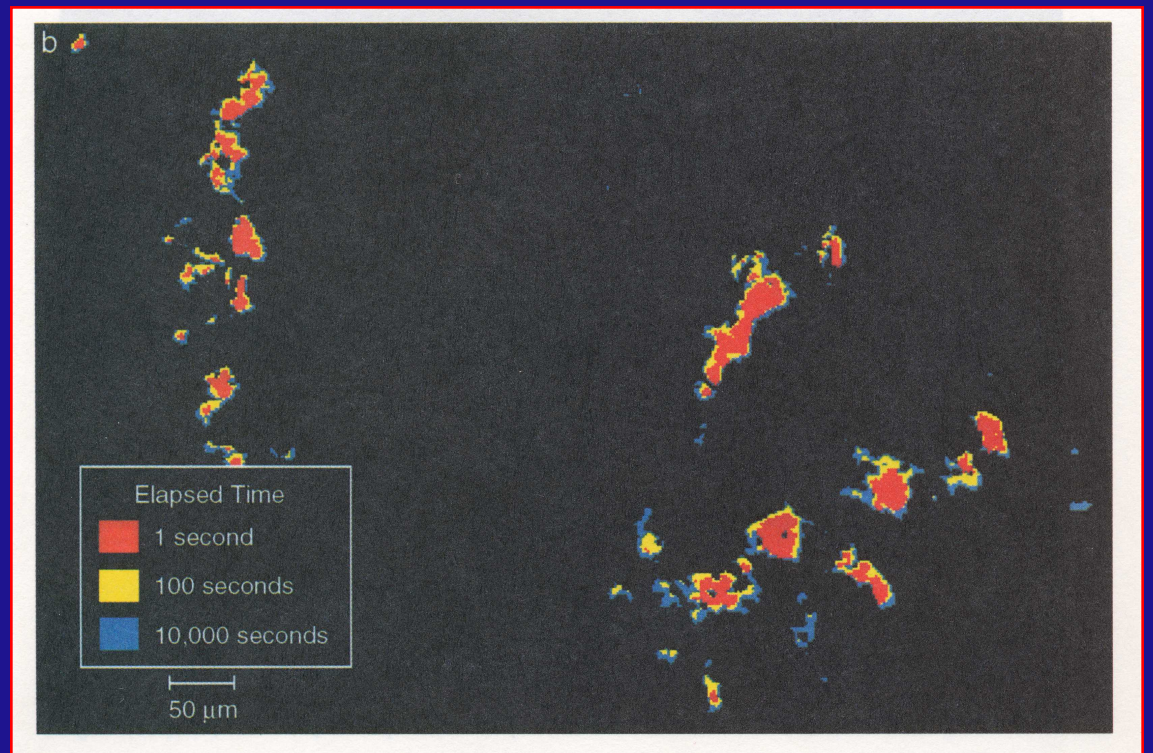


# Consider the role of compaction and densification

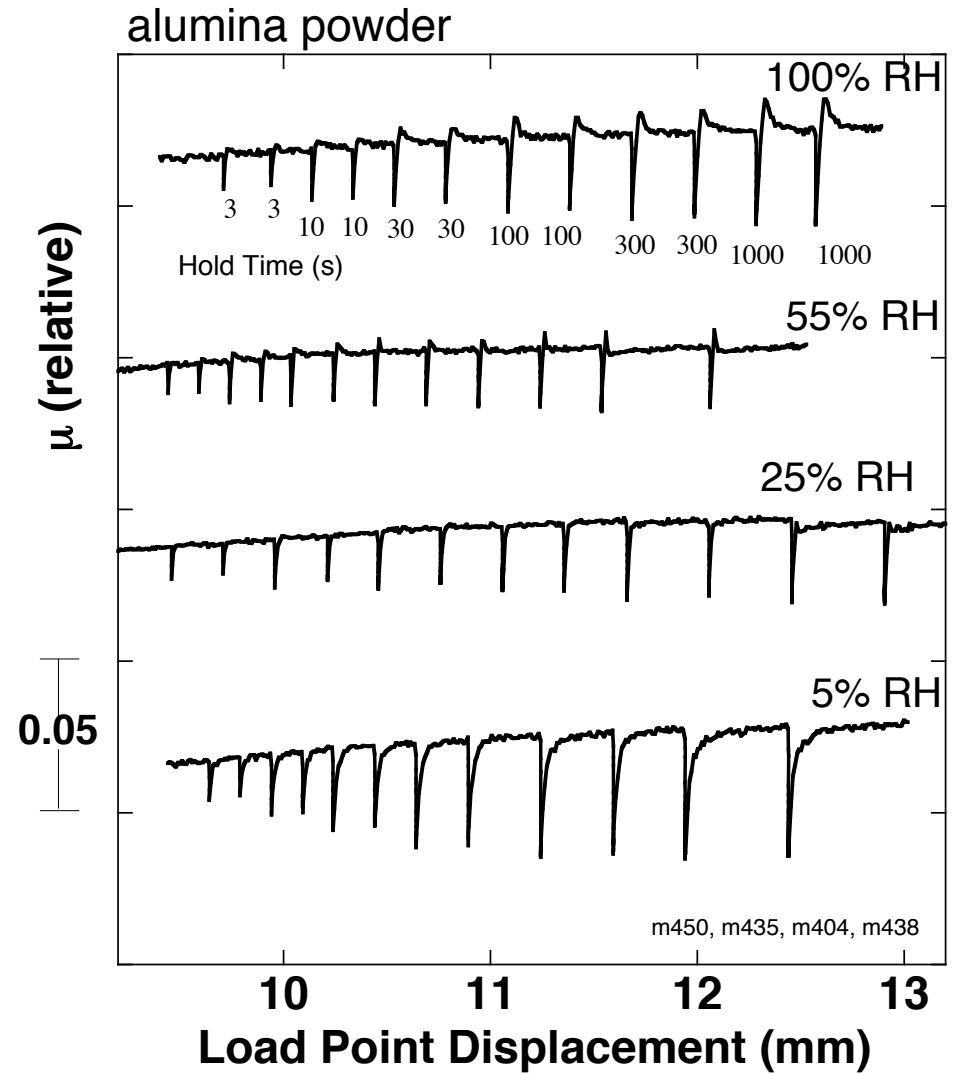
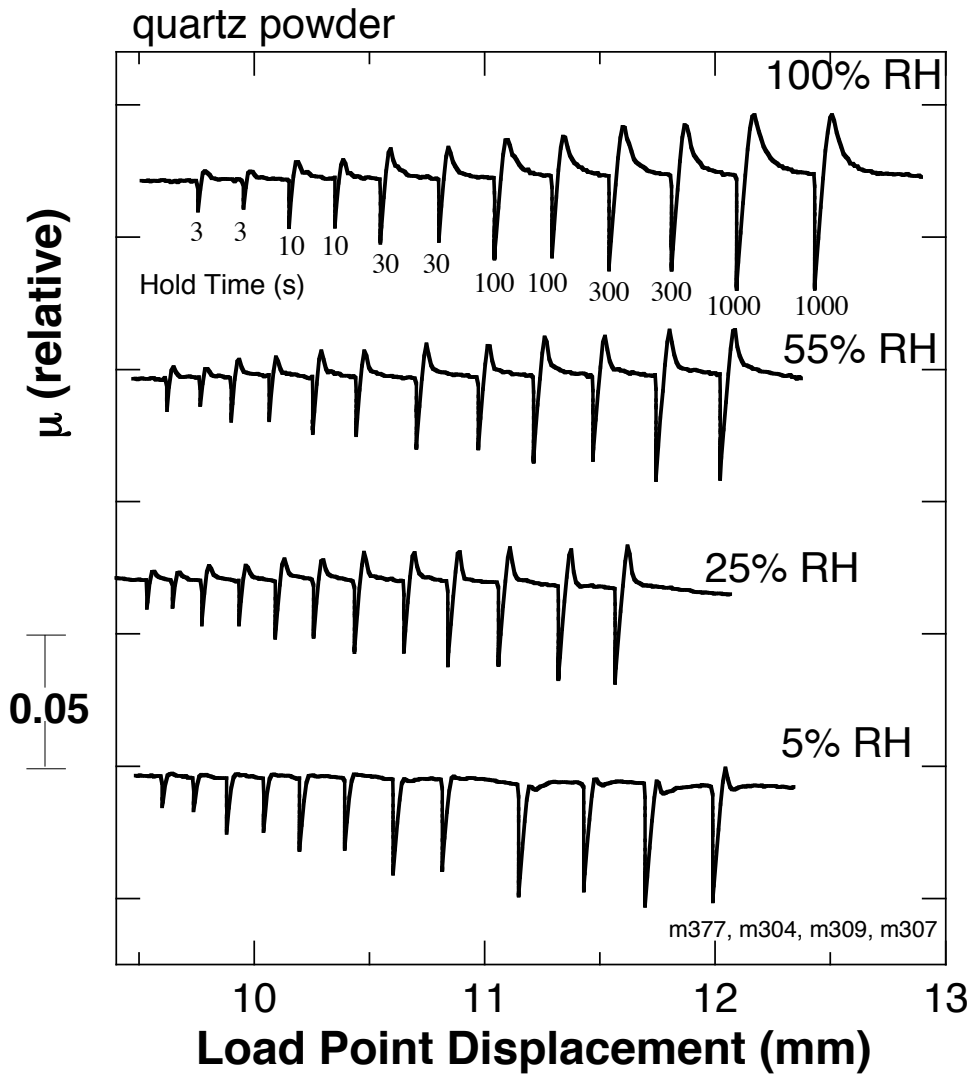




Frictional aging.  
How does it work?



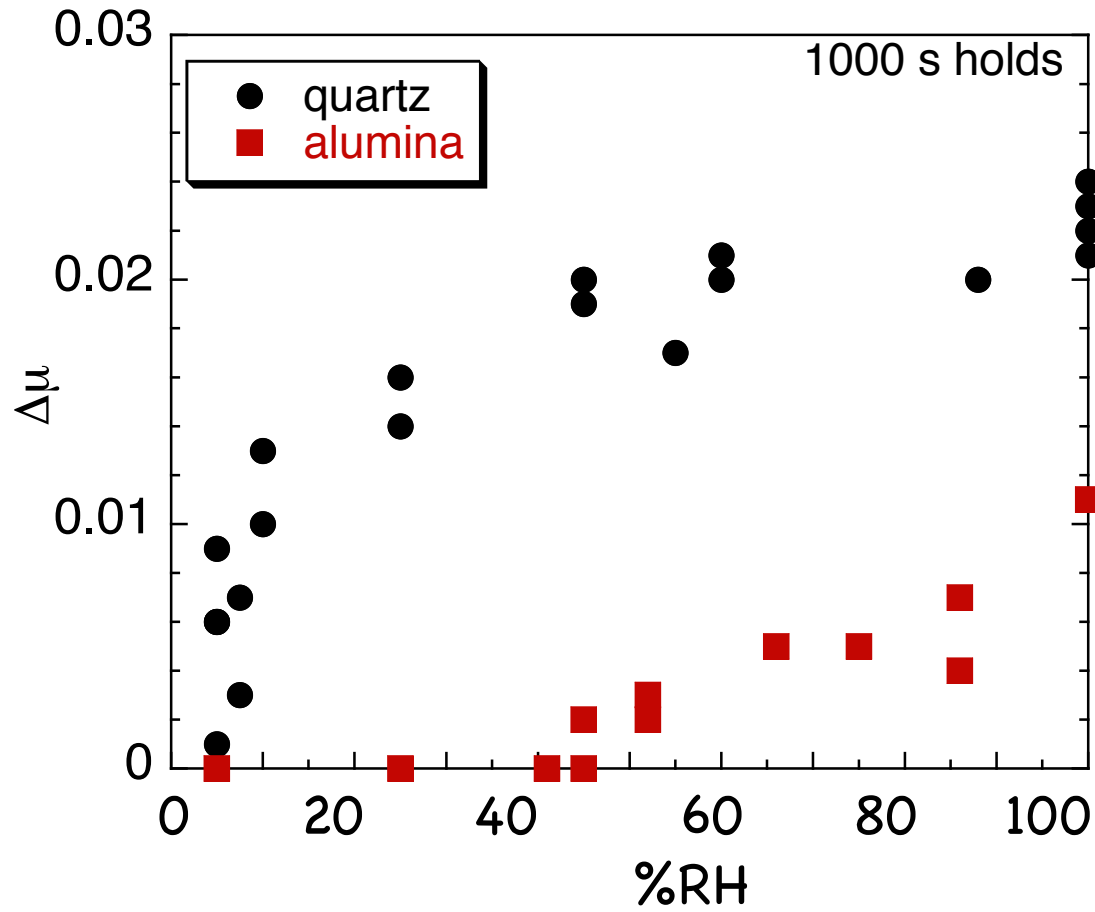
# Chemically-Assisted Frictional Aging; Creep at Adhesive Contact Junctions



**In-situ Particle Comminution; Production of Fresh Surface Area**

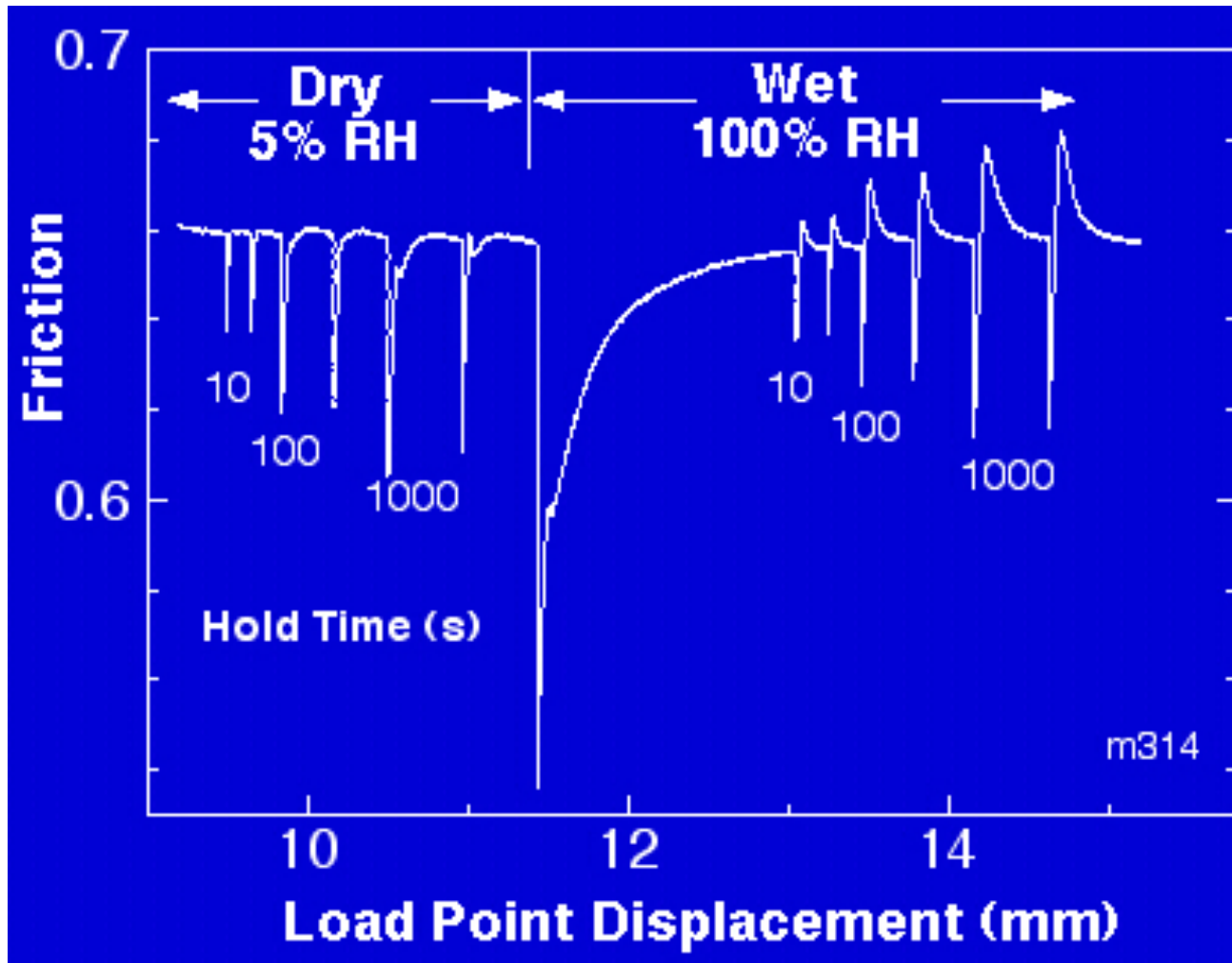
**Frye and Marone, JGR 2002**

## Granular quartz



**Hydrolytic Weakening  
causes enhanced rate of  
strengthening**

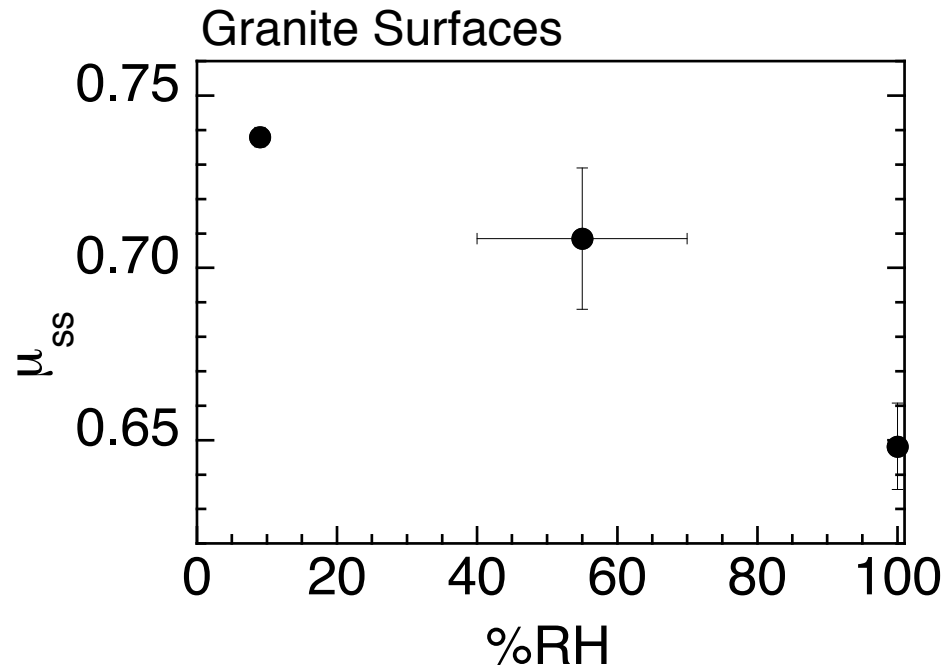
# Chemically-Assisted Frictional Aging; Creep at Adhesive Contact Junctions



Hydrolytic Weakening causes enhanced rate of strengthening, but base level frictional strength is unchanged

Frye and Marone, JGR 2002

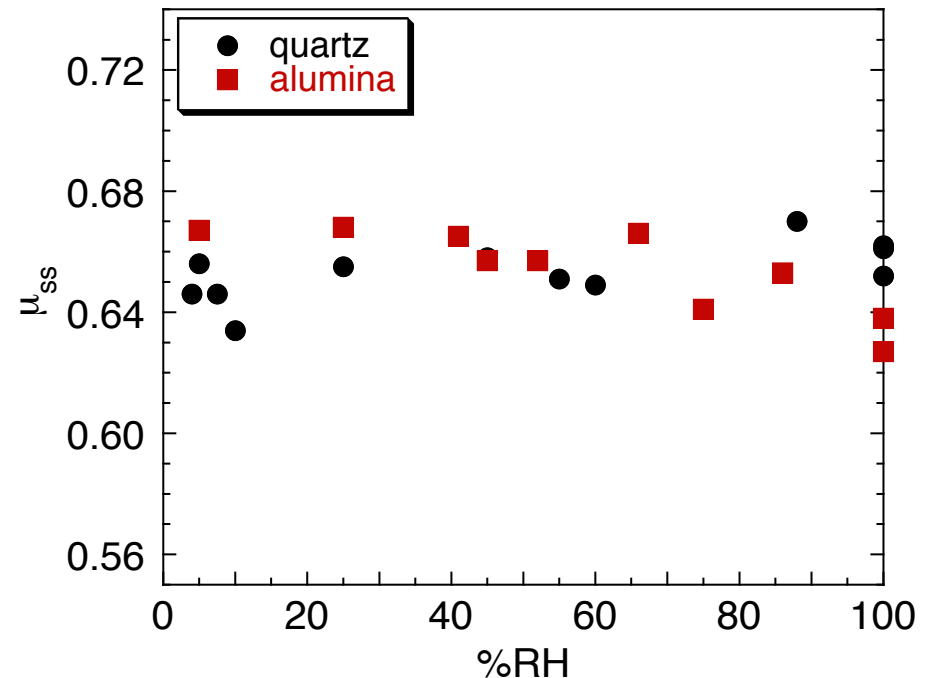
## Granite Surfaces



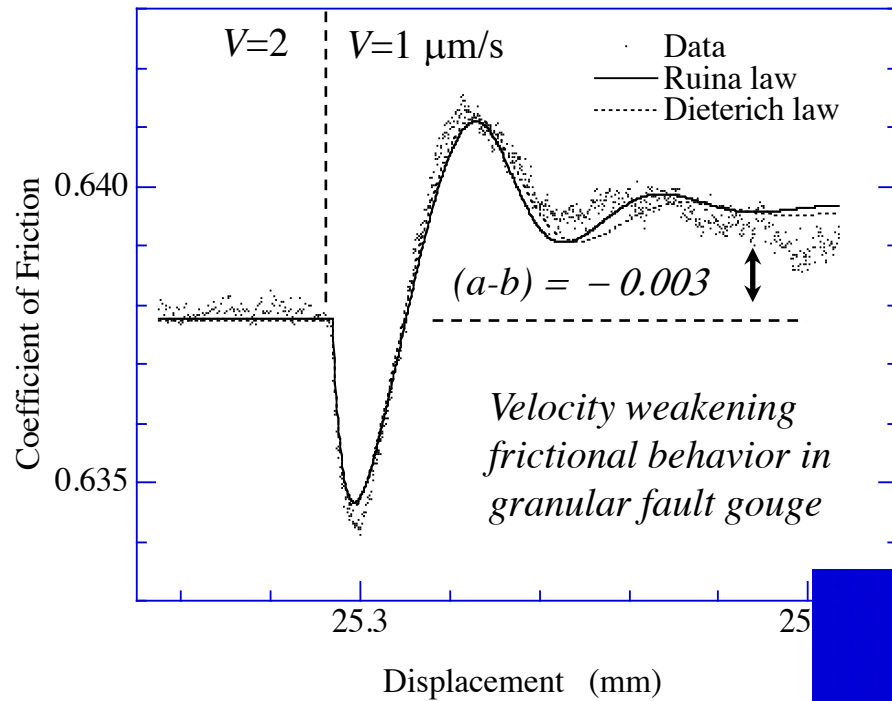
**Solid Surfaces: Base level of frictional strength decreases with increasing water content (cf. Dieterich & Conrad, 1984)**

**Interpretation: Contact junctions subject to time dependent strengthening or growth, which inhibits sliding, but particle rolling is not affected by these factors.**

## Granular Materials

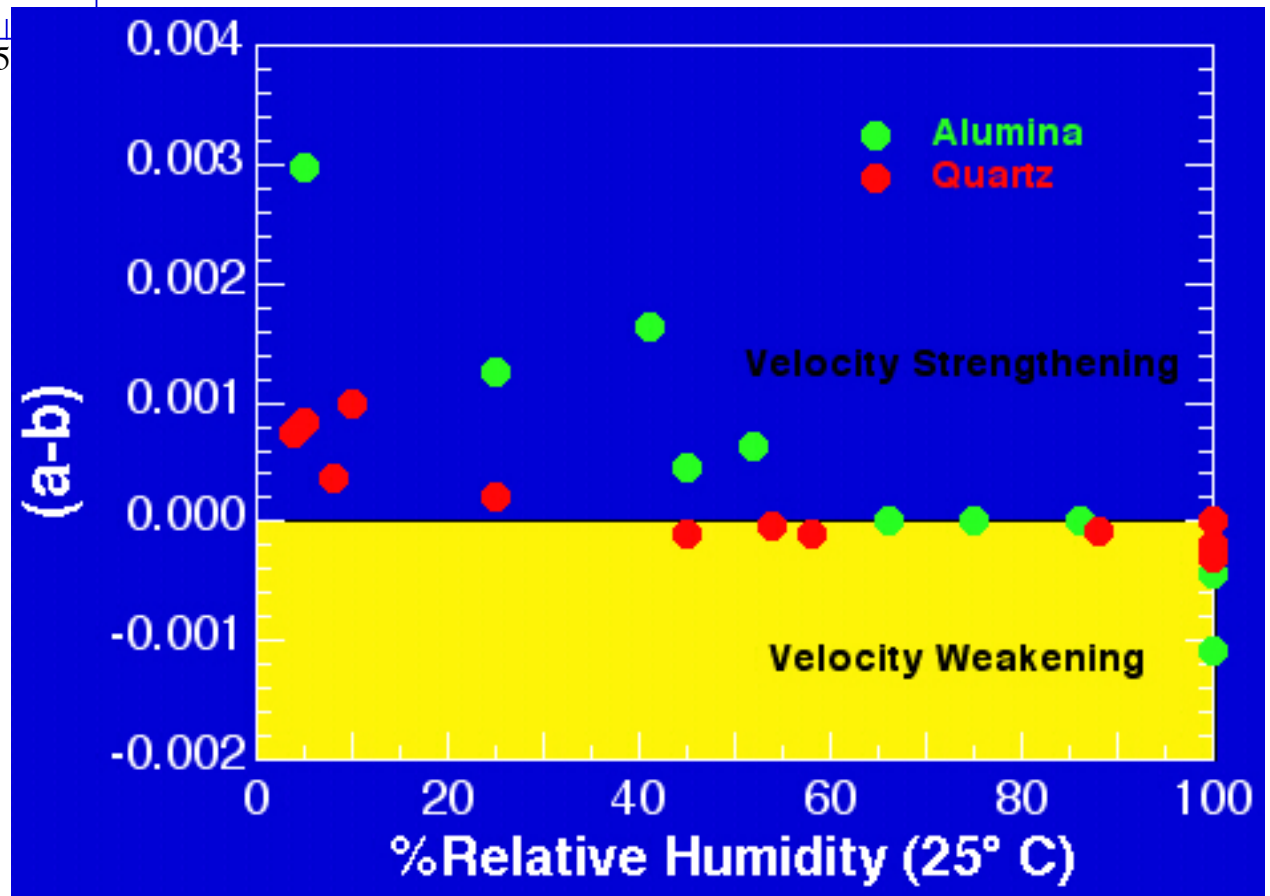


**Granular Materials: Frictional strength is independent of water content**



Velocity dependence of steady state friction varies changes from positive to negative. (cf. Tullis and co-workers)

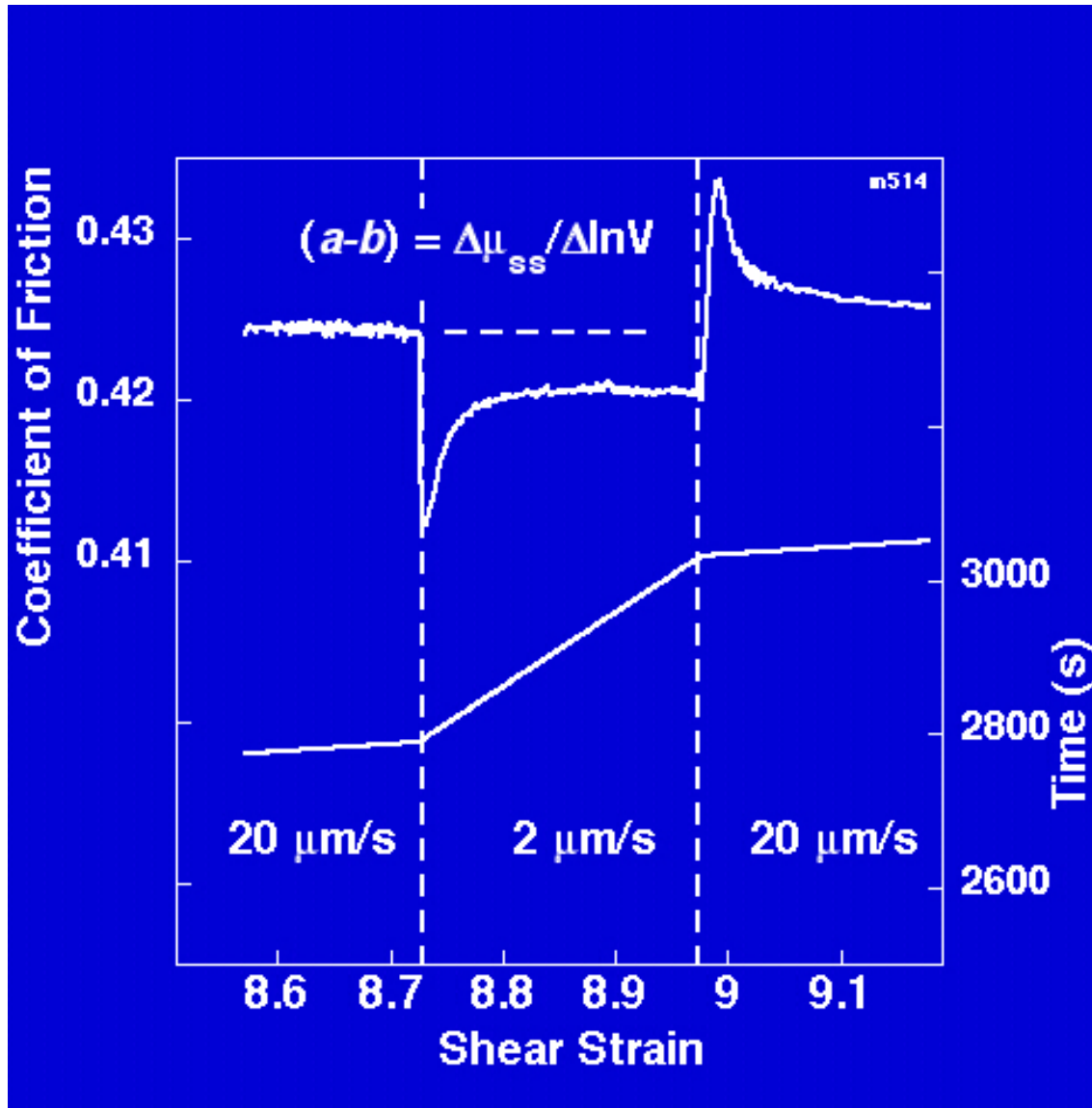
Chemically-assisted creep at adhesive contact junctions



# Measuring the velocity dependence of friction

## Frictional Instability

Requires  $(a-b) < 0$



### Constitutive Modelling

Rate and State Friction Law

Elastic Interaction, Testing Apparatus

$$\mu(\theta, v) = \mu_0 + a \ln\left(\frac{v}{v_0}\right) + b \ln\left(\frac{v_0 \theta}{D_c}\right)$$

$$\frac{d\theta}{dt} = 1 - \frac{v\theta}{D_c}$$

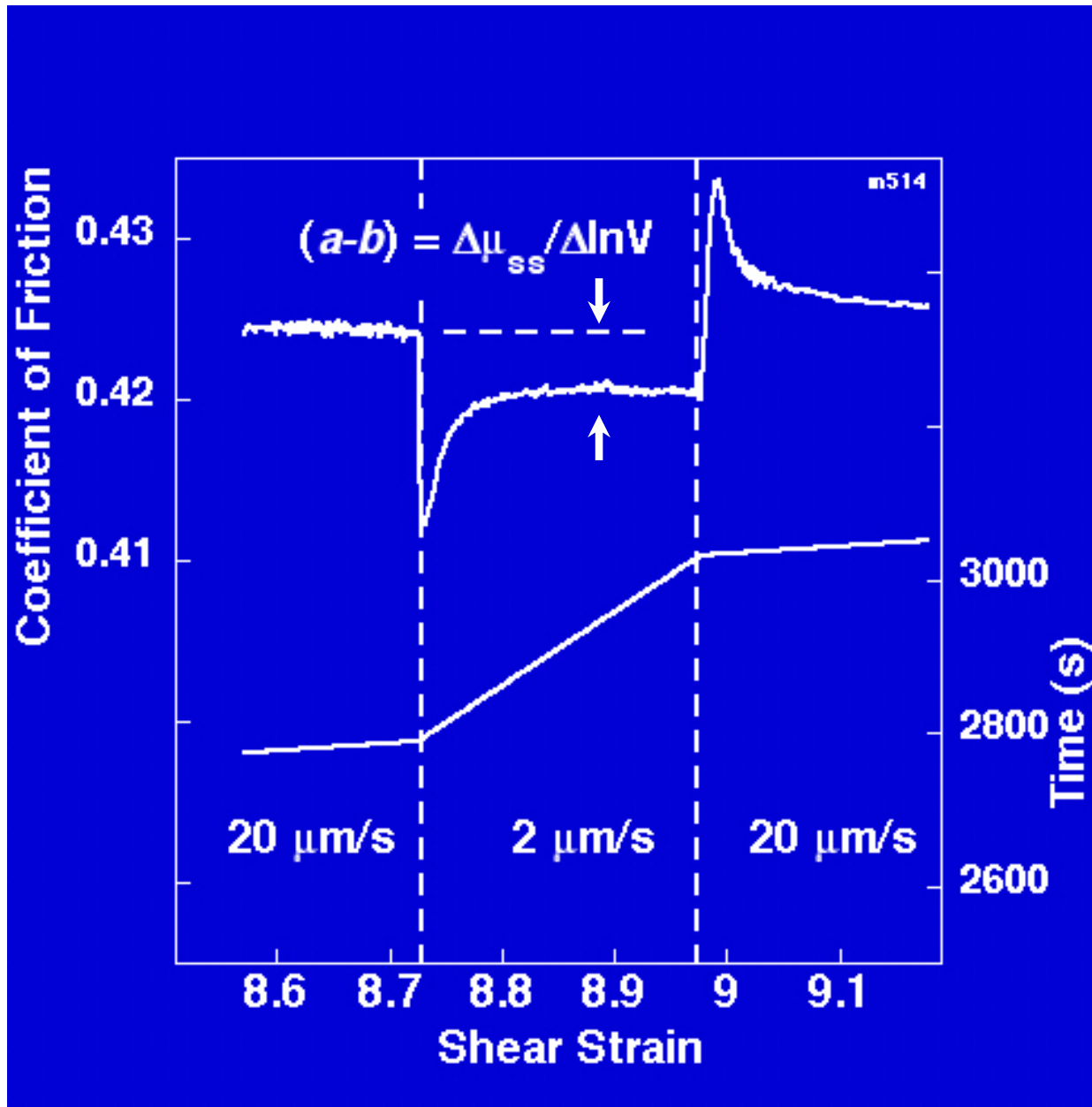
$$\theta_{ss} = \frac{D_c}{v}$$

$$\Delta\mu_{ss} = (a-b) \ln\left(\frac{v}{v_0}\right)$$

$$\frac{d\mu}{dt} = k' (v_{lp} - v)$$



Results: Velocity stepping  
 Measuring the velocity dependence of friction



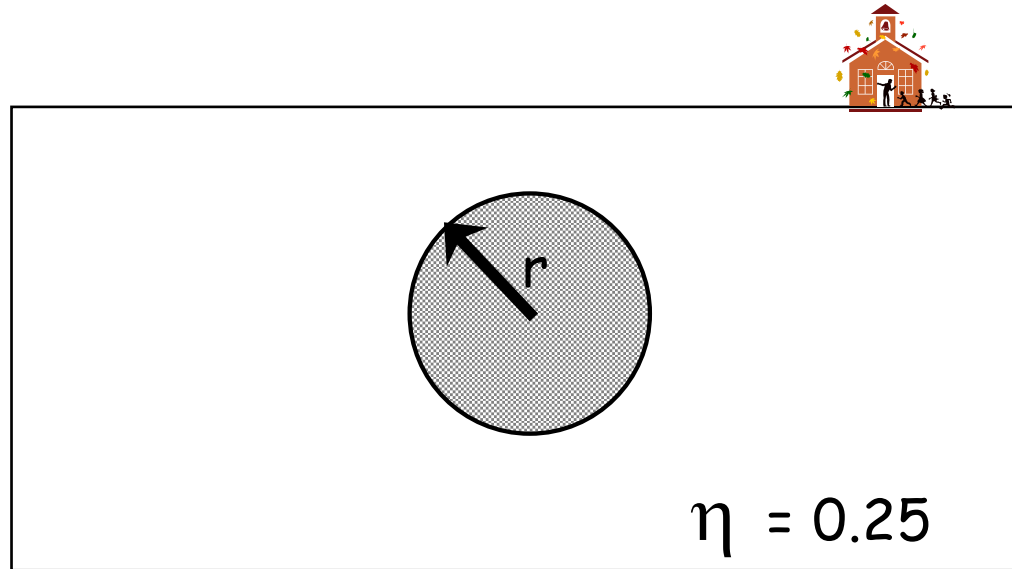
Frictional Instability

Requires  $K < K_c$

$$K_c = \frac{\sigma_n(b - a)}{D_c}$$

This example shows  
 steady-state velocity  
 strengthening:  
 $(a-b) > 0$

# Rupture Patch Size for Earthquake Nucleation



$$K = \frac{\Delta\sigma}{\bar{u}} = \frac{7\pi G}{16 r}$$

$$K_c \approx \frac{\sigma_n(b-a)}{D_c}$$

earthquake nucleation when

$$\frac{K}{K_c} \approx 1.0$$

See also: Earthquake nucleation on rate and state faults  
- Aging and slip laws, Ampuero & Rubin, JGR 2008

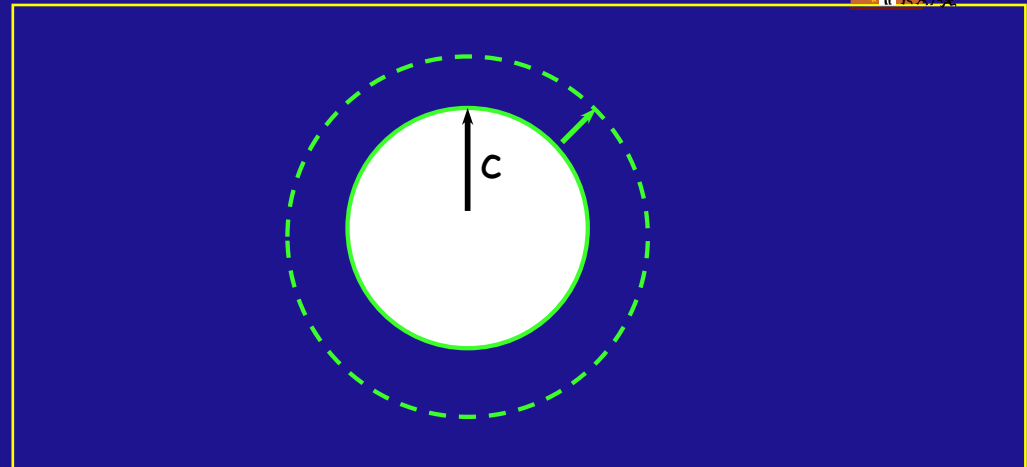
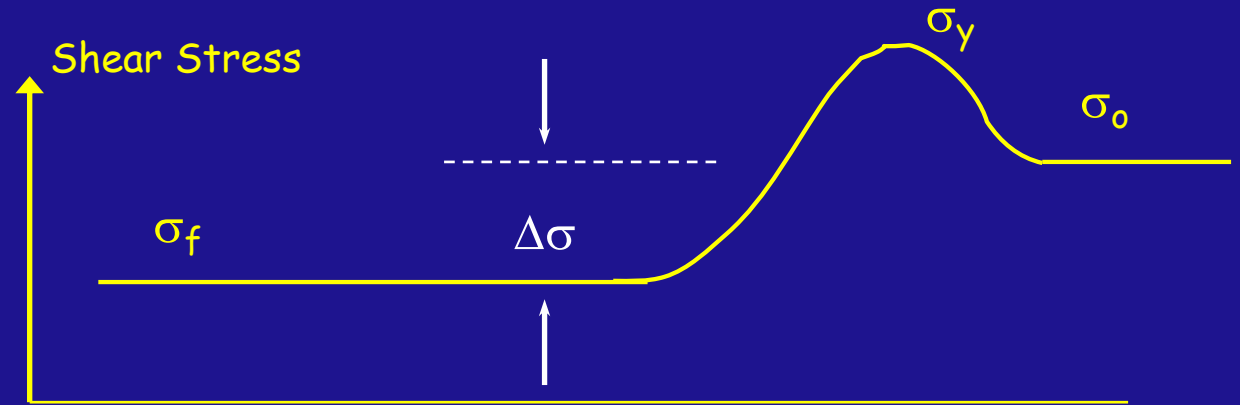
$$h^* = r_c \approx C \frac{GD_c}{\sigma_n(b-a)}$$

# Dislocation model for fracture and earthquake rupture

Dislocation model, circular crack

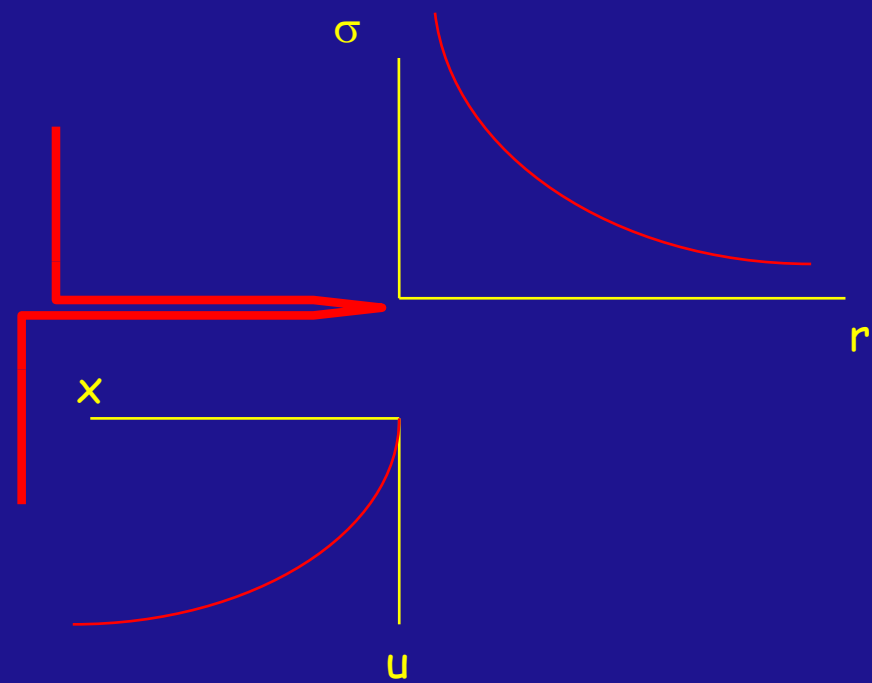
$$\Delta\sigma = (\sigma_0 - \sigma_f)$$

$$\Delta u(x,y) = \frac{24}{7\pi} \frac{\Delta\sigma}{\mu} \sqrt{c^2 - (x^2 + y^2)}$$



# Crack tip stress field, real materials

- Singular crack (Eshelby)



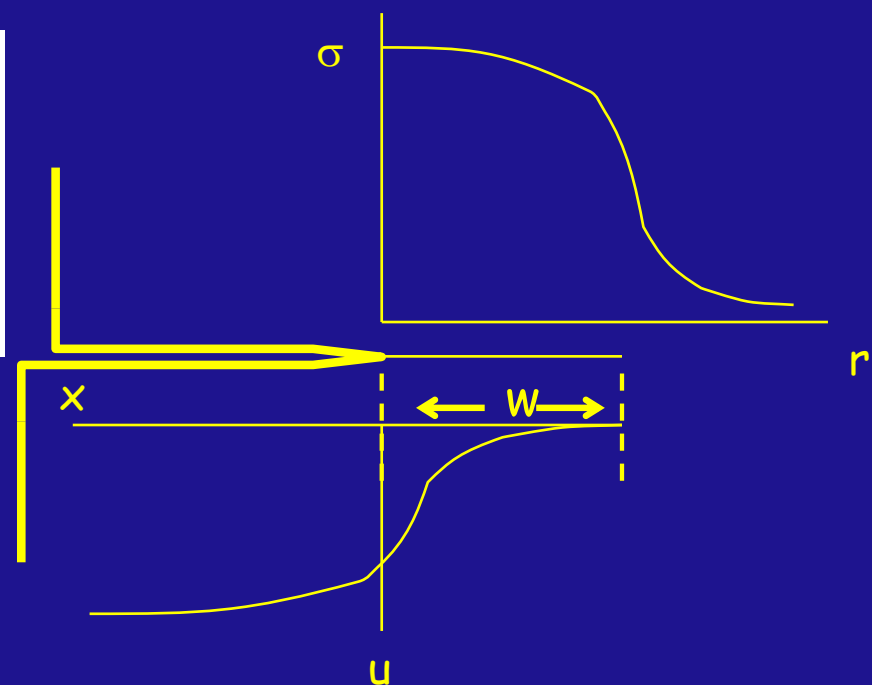
- Dugdale (Barenblatt)

• Assume a yield criterion,  $\sigma_y$ , within a crack tip region

$s = w$

$$\Delta u = \frac{(1-\eta)(\sigma_o - \sigma_f) c}{2\pi\mu} f(\theta, \theta_2), \text{ see 1.30}$$

$$\theta = \cos^{-1}\left(\frac{2x}{c}\right) \text{ for } |x| < \frac{c}{2} \text{ and } \theta_2 = \cos^{-1}\left(\frac{c-2s}{c}\right)$$

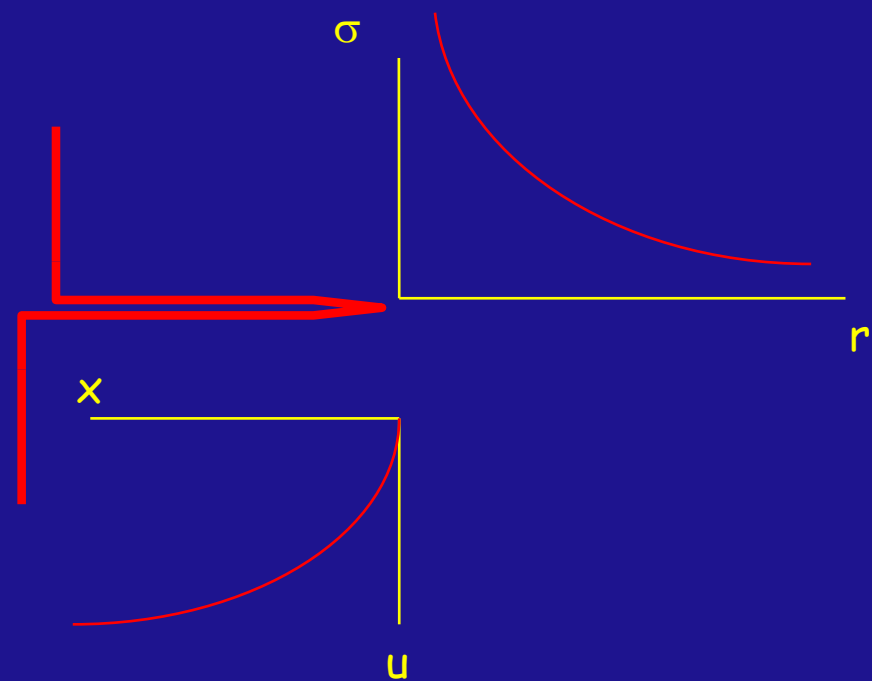


# Crack tip stress field, real materials

- Singular crack (Eshelby)

- Dugdale (Barenblatt)

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- e.g., can we 'read' the state of stress in the crust from earthquake (fault) data

Barenblatt, G. I., 1959, The formation of brittle cracks during brittle fracture. General ideas and hypotheses. Axially-symmetric cracks. *Appl. Math. Mech.* 23, 1273 - 1282.

Barenblatt, G. (1962). The mathematical theory of equilibrium cracks in brittle fracture. *Advances in Applied Mechanics*, 7, 55-129.

Dugdale, D. (1960). Yielding of steel sheets containing slits. *Journal of the Mechanics and Physics of Solids*, 8, 100-104.

# Cohesive zone crack model, applies to fracture and/or friction

## • Dugdale (Barenblatt)

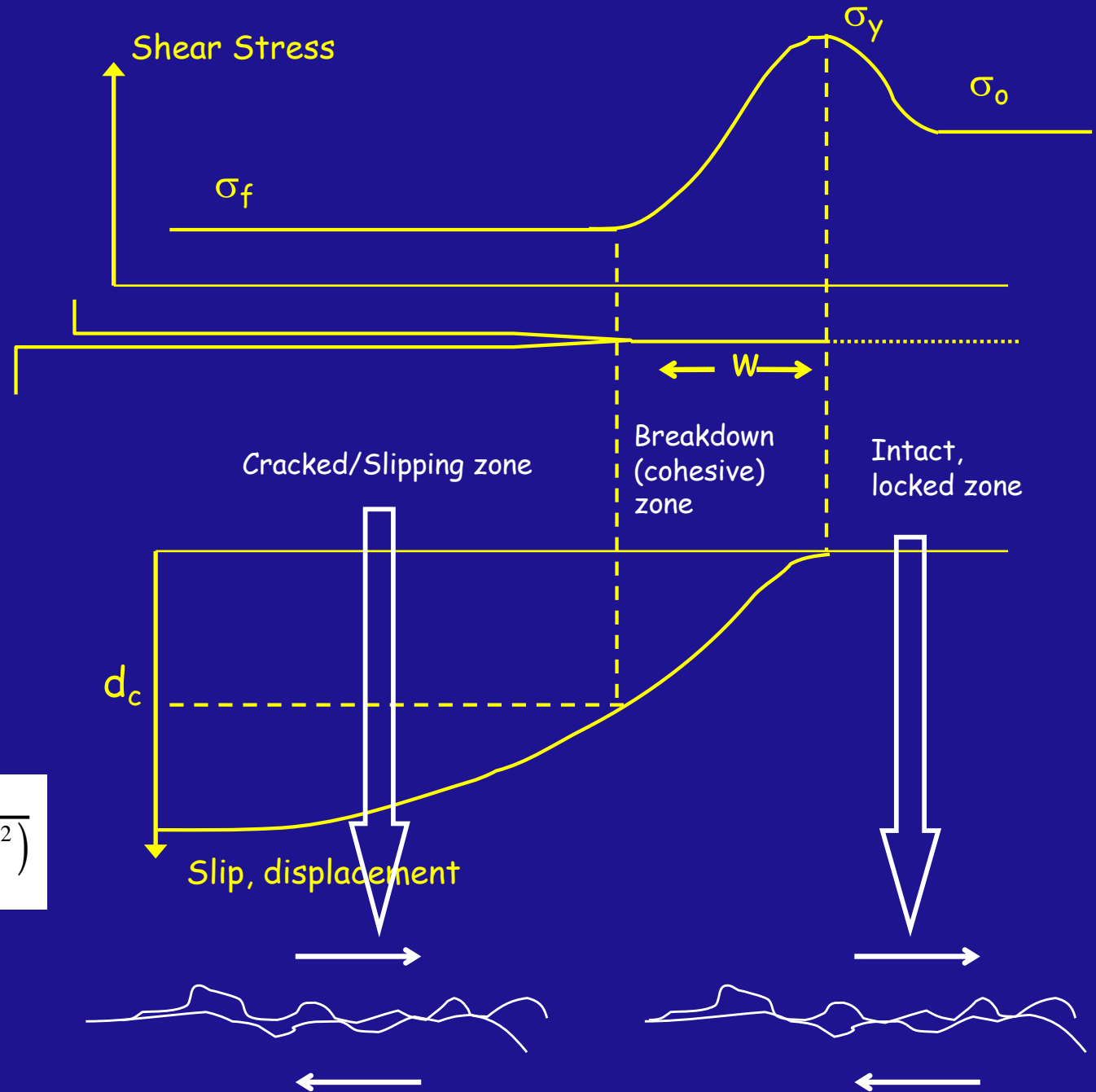
$$\Delta u = \frac{(1-\eta)(\sigma_o - \sigma_f) c}{2\pi\mu} f(\theta, \theta_2), \text{ see 1.30}$$

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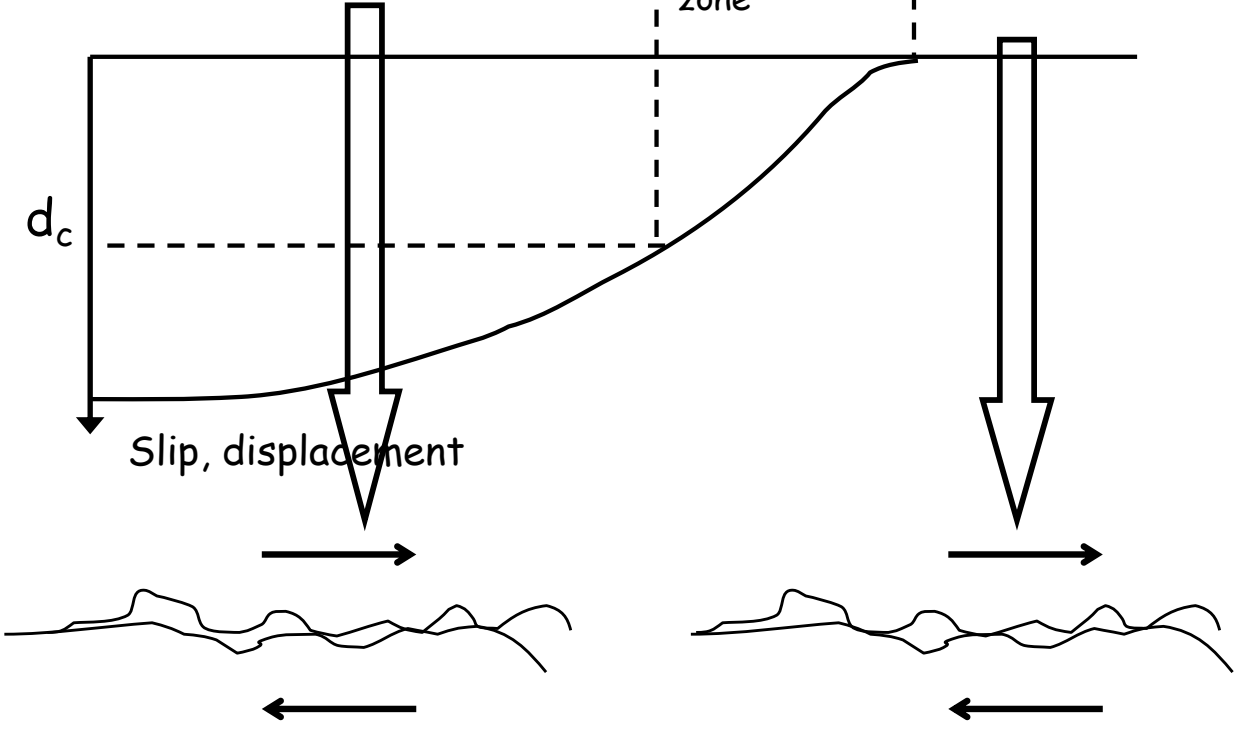
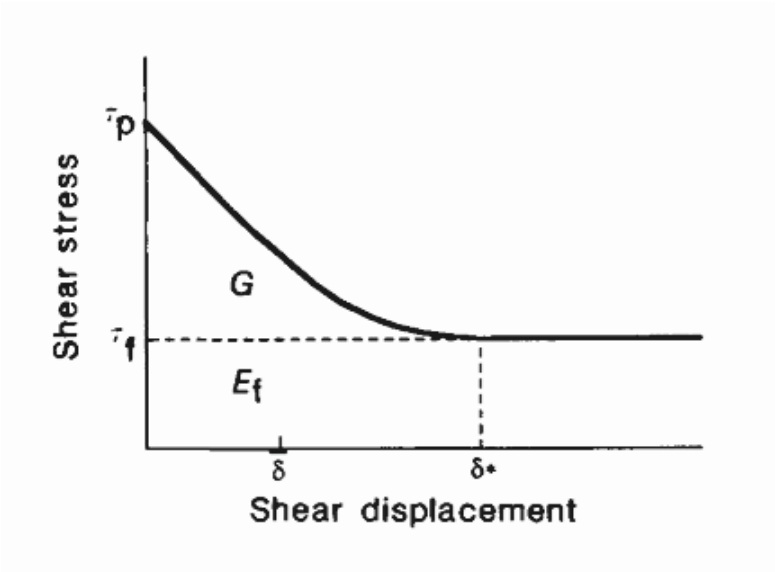
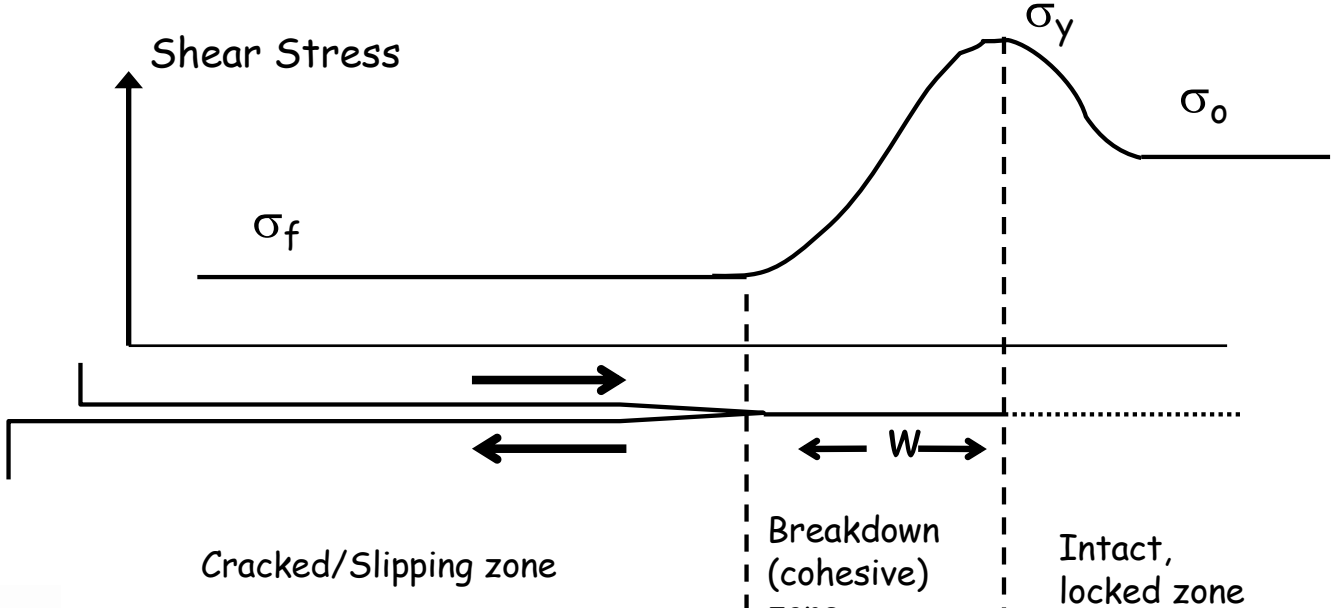
## Dislocation model, circular crack

$$\Delta\sigma = (\sigma_o - \sigma_f)$$

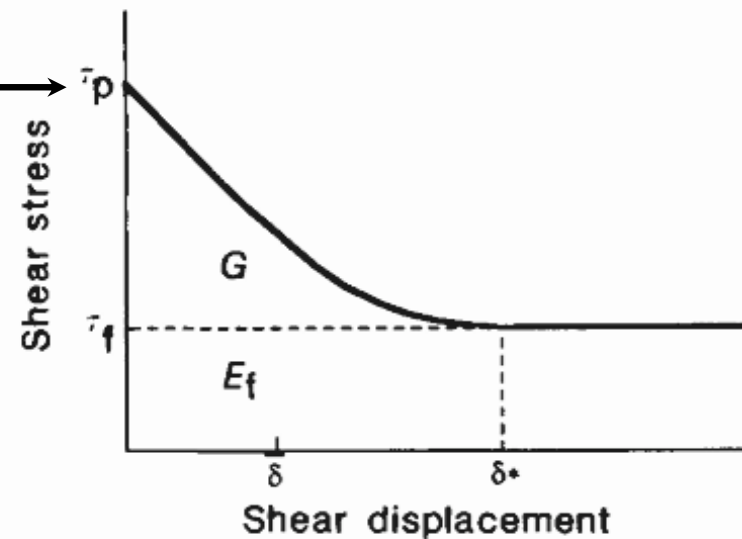
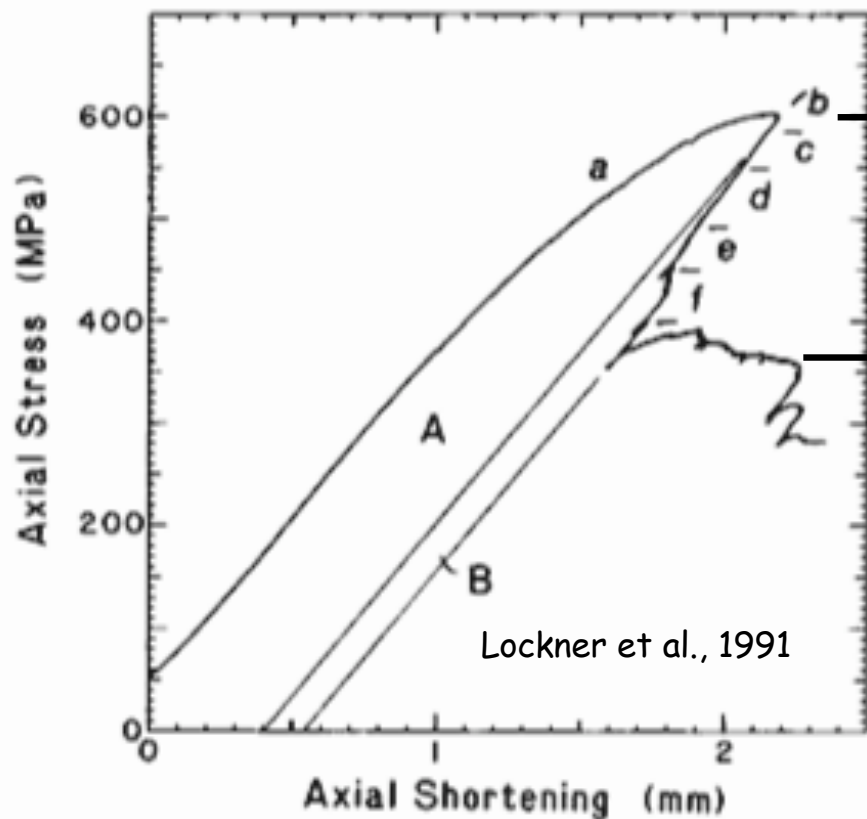
$$\Delta u(x,y) = \frac{24}{7\pi} \frac{\Delta\sigma}{\mu} \sqrt{c^2 - (x^2 + y^2)}$$



# Cohesive zone, slip weakening crack model for friction



# Shear Fracture Energy from Postfailure Behavior



Inferred shear stress vs. slip relation for slip-weakening model. (based on Wong, 1982)

Wong, 1982, found that shear stress dropped  $\sim 0.2$  GPa over a slip distance of  $\sim 50$  microns.

Exercise: Estimate  $G$  from this data and compare it to the values reported in Scholz (Table 1.1) and Wong, 1982.

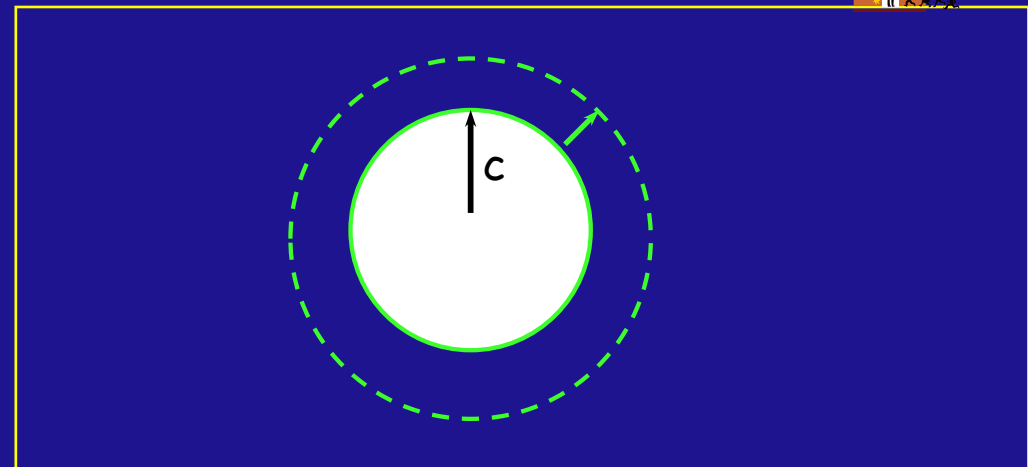
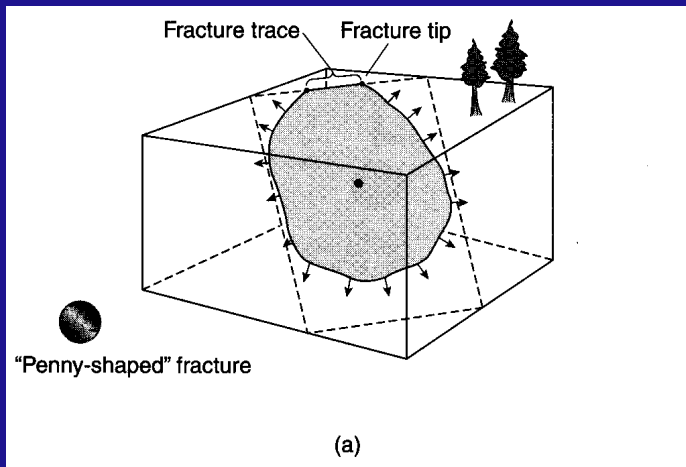
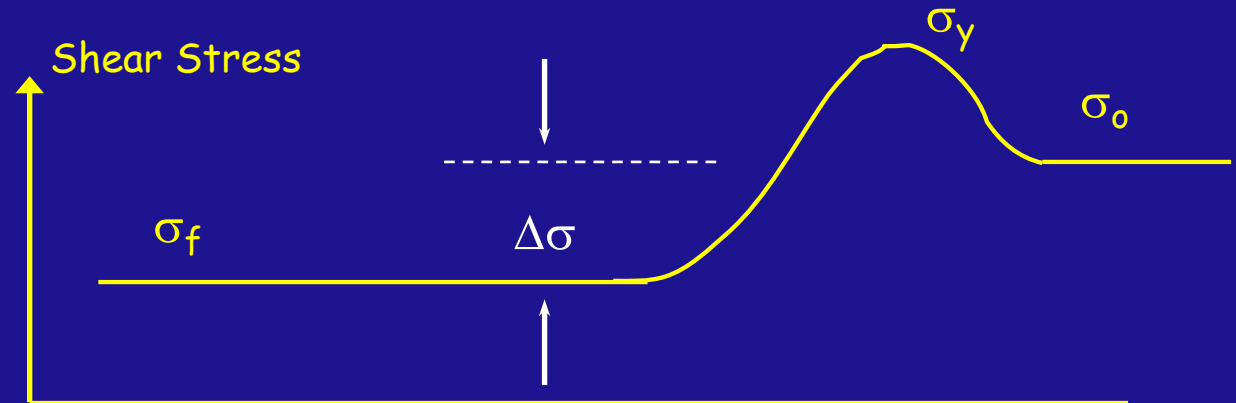


# Dislocation model for fracture and earthquake rupture

Dislocation model, circular crack

$$\Delta\sigma = (\sigma_0 - \sigma_f)$$

$$\Delta u(x,y) = \frac{24}{7\pi} \frac{\Delta\sigma}{\mu} \sqrt{c^2 - (x^2 + y^2)}$$



$$\Delta\sigma = \frac{7\pi}{16} G \frac{\Delta\bar{u}}{r}$$

Relation between stress drop and slip for a circular dislocation (crack) with radius  $r$   
For  $\eta=0.25$ , Chinnery (1969)

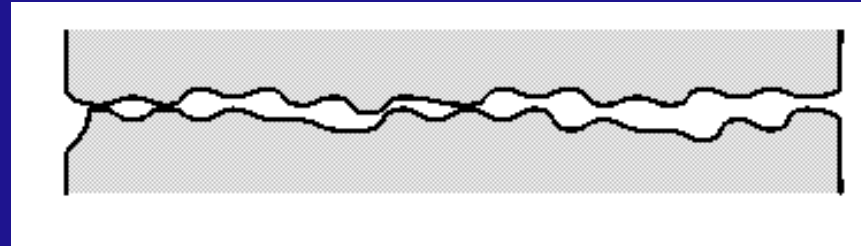
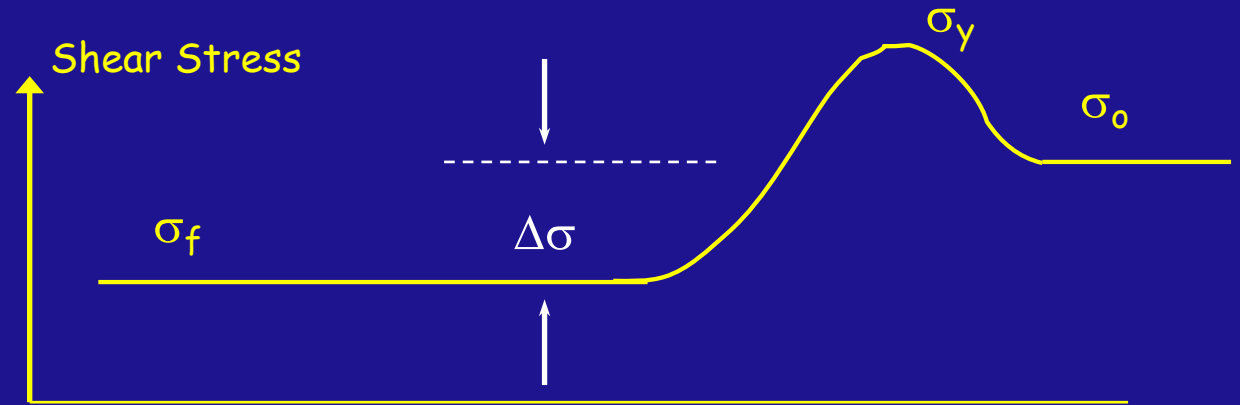
•Importance of slip: e.g.,  $M_0 = \mu A u$

# Dislocation model for fracture and earthquake rupture

Dislocation model, circular crack

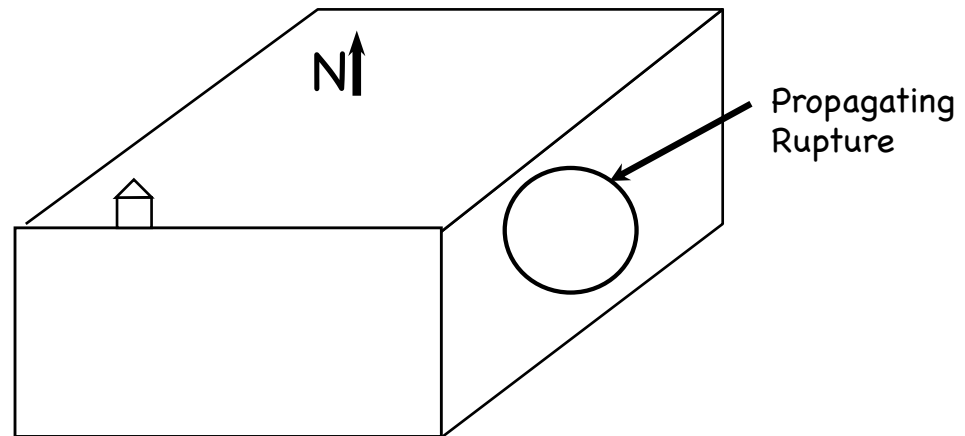
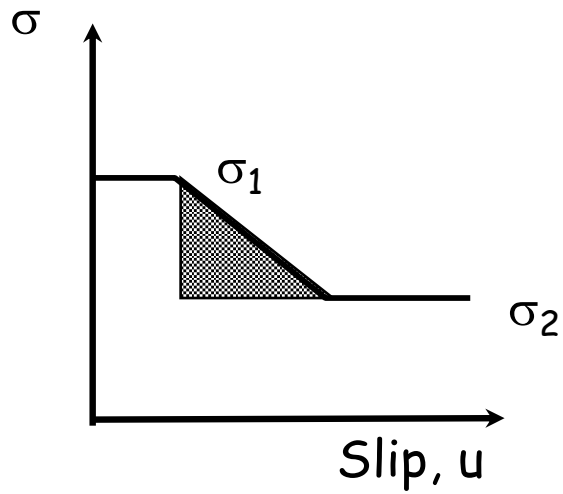
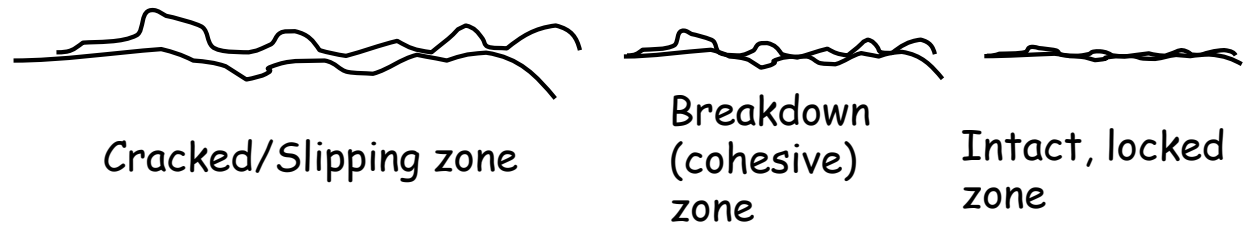
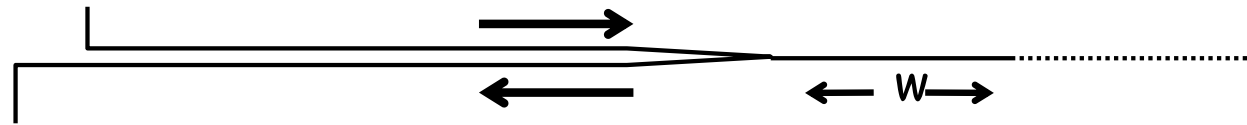
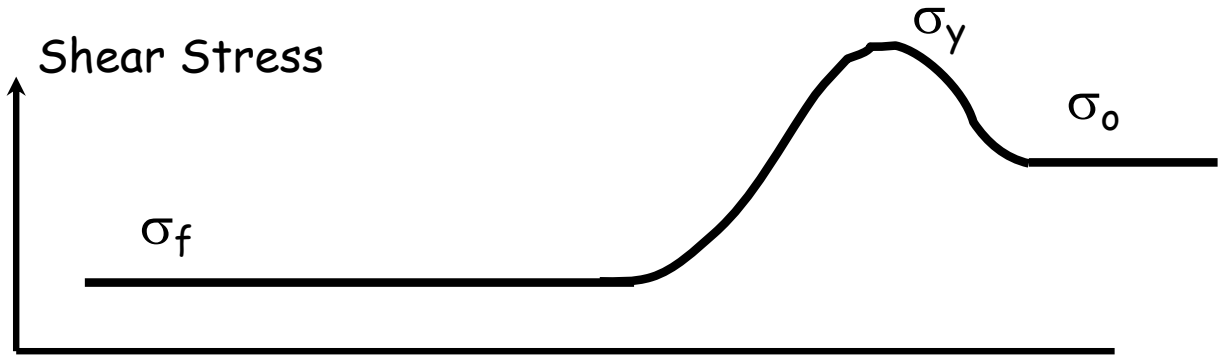
$$\Delta\sigma = (\sigma_0 - \sigma_f)$$

$$\Delta u(x,y) = \frac{24}{7\pi} \frac{\Delta\sigma}{\mu} \sqrt{c^2 - (x^2 - y^2)}$$



Mode II crack propagation at speed  $V_r$   $\longrightarrow$

cohesive zone/slip  
weakening crack model  
for friction



Mode II crack propagation at speed  $Vr$   $\rightarrow$

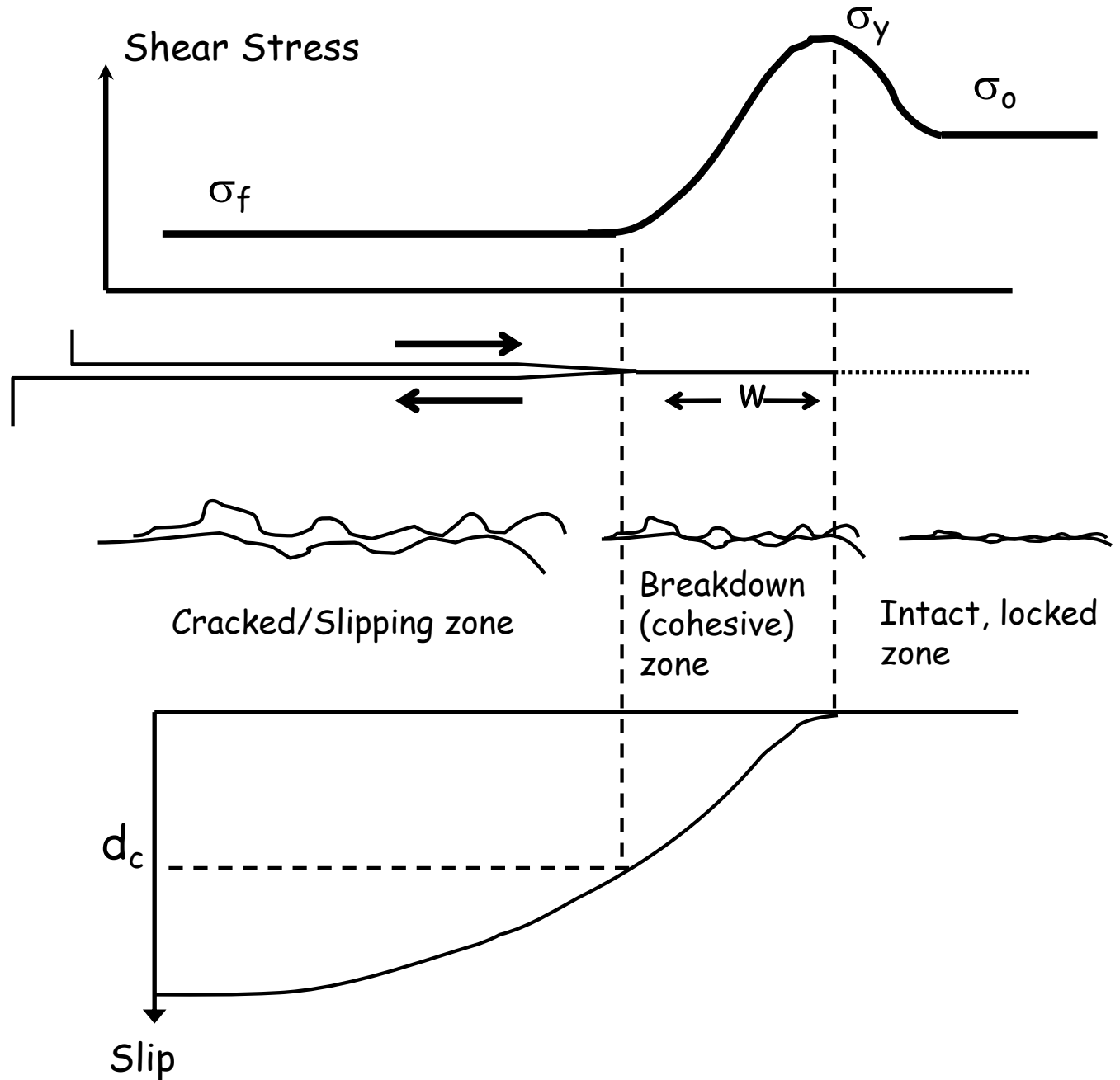
cohesive zone/slip  
weakening crack model  
for friction

$f_{max}$  scales as:  
 $f_{max} = w/Vr$

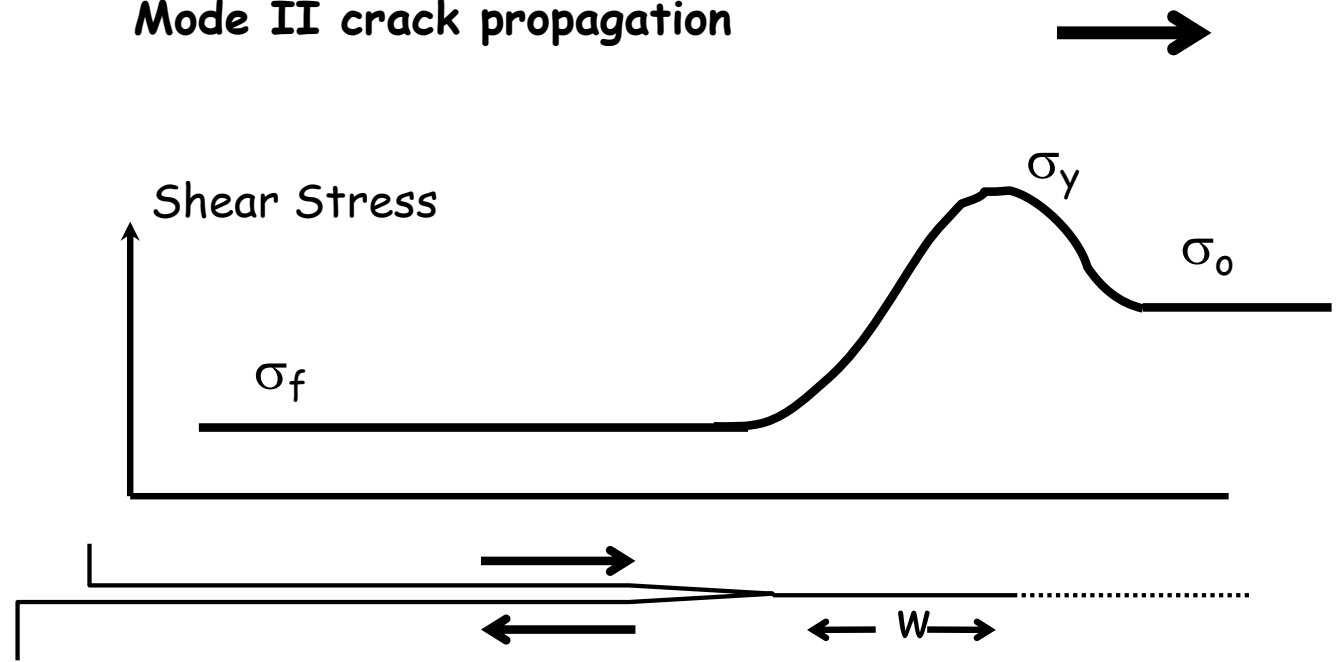
$$\Delta\sigma = \frac{7\pi}{24} G \frac{u_{max}}{r}$$

$$\Delta\sigma = \frac{7\pi}{24} G \frac{D_c}{w}$$

$$D_c = \frac{24}{7\pi} \frac{\Delta\sigma}{G} w$$



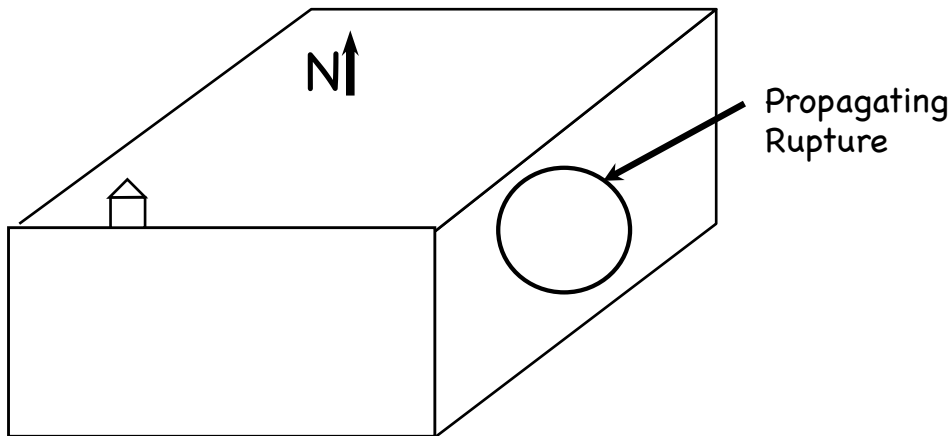
# Mode II crack propagation



Cohesive zone length  $W$  scales as:

$$w \approx C \frac{D_c G}{\Delta\sigma}$$

$$\frac{w}{D_c} \approx C \frac{G}{\Delta\sigma}$$



## Slip weakening model (Ida, 1972, 1973)

$$R_c = \frac{CGD_c}{\sigma_n(b-a)} \text{ or } L_c = \frac{E D_c}{2(1-\nu^2)\sigma_n \Delta\mu}$$

Recall the derivation of this result:

### Frictional Instability

Requires  $K < K_c$

$$K_c = \frac{\sigma_n(b-a)}{D_c}$$

$$\Delta\sigma = \frac{7\pi}{24} G \frac{u_{max}}{r}$$

Relation between stress drop and slip for a circular dislocation (crack) with radius  $r$   
For  $\nu = 0.25$ , Chinnery (1969)

$$\Delta\sigma = \frac{7\pi}{24} G \frac{D_c}{w}$$

$$D_c = \frac{24}{7\pi} \frac{\Delta\sigma}{G} w$$